## Lecture 25 Prof. Peter Shor

In this lecture, we will give a lower bound that shows that Grover's algorithm is nearly optimal. In particular, we assume that you have a quantum state space with basis $\{|1\rangle,|2\rangle, \ldots|N\rangle\}$ and an oracle function $O_{x}$ such that

$$
O_{x}|y\rangle=\left\{\begin{aligned}
-|y\rangle & \text { if } y=x \\
|y\rangle & \text { if } y \neq x
\end{aligned}\right.
$$

We will show that any algorithm which always returns the output $|x\rangle$ must call the oracle on the order of $\sqrt{N}$ times. We first show that this is true if the algorithm is required to give the answer $|x\rangle$ with probability 1 , and then show that the proof still works if we just require the answer with probability at least $1-\epsilon$.

One amazing thing about this proof is that it was developed independently from Grover's algorithm, starting with the motivation of whether you can show that quantum computers can't solve NP-complete problems (although it was not published until after Grover's paper, and the original paper mentions Grover).

We model our algorithm as follows. First, by the principle of deferred measurement (see homework), we can assume that we postpone all the measurements in our algorithm until the end. The algorithm will then consist of unitary operations and calls to the oracle. We will define $U_{j}$ so that the algorithm alternates calling the oracle $O_{x}$ and implementing the unitary $U_{k}$. Suppose the algorithm requires $t$ calls to the oracle in the worst case. Then, we define the algorithm to be:

$$
\left|\psi_{t}^{x}\right\rangle=U_{t} O_{x} U_{t-1} O_{x} U_{t-2} \ldots U_{1} O_{x}\left|\psi_{0}\right\rangle
$$

If not all runs of the algorithm use the same number of calls to the oracle, we can always add more dummy calls to the oracle that in essence do nothing.

Now, suppose we are able to identify $x$ with probability 1 after this algorithm. Then it must be the case that all the $\left|\psi_{t}^{x}\right\rangle$ are orthonormal. We will bound how quickly they are able to become distinct.

Define the state of the quantum computer after $k$ oracle calls to be

$$
\left|\psi_{k}^{x}\right\rangle=U_{k} O_{x} U_{k-1} O_{x} U_{k-2} \ldots U_{1} O_{x}\left|\psi_{0}\right\rangle
$$

We will define a measure of how quickly the states $\left|\psi_{k}^{x}\right\rangle$ are diverging from each other. In order to do this, let

$$
\left|\psi_{t}\right\rangle=U_{t} U_{t-1} U_{t-2} \ldots U_{1}\left|\psi_{0}\right\rangle
$$

so $\left|\psi_{t}\right\rangle$ is what we get when we run the algorithm with no oracle calls (or equivalently, when there are no solution states). Now, define

$$
\left.D_{k}=\sum_{x}| | \psi_{k}^{x}\right\rangle-\left.\left|\psi_{k}\right\rangle\right|^{2}
$$

We will use $D_{k}$ as a measure of how far the states $\left|\psi_{k}^{x}\right\rangle$ are from each other. In other to get a bound on the number of oracle calls we need, we will need to prove two things:

1. an upper bound on $D_{k}$,
2. a proof that if $D_{t}$ is too small, then we cannot distinguish the $\left|\psi_{x}\right\rangle$ from each other well.

We first prove the upper bound for $D_{k}$.
We have

$$
\begin{aligned}
D_{k+1} & \left.\left.=\sum_{x}\left|U_{k+1} O_{x}\right| \psi_{k}^{x}\right\rangle-U_{k+1}\left|\psi_{k}\right\rangle\right)\left.\right|^{2} \\
& \left.=\sum_{x}\left|O_{x}\left(\left|\psi_{k}^{x}\right\rangle-\left|\psi_{k}\right\rangle\right)+\left(O_{x}-I\right)\right| \psi_{k}\right\rangle\left.\right|^{2} \\
& \left.\left.\leq \sum_{x}\left|O_{x}\left(\left|\psi_{k}^{x}\right\rangle-\left|\psi_{k}\right\rangle\right)\right|^{2}+\left|\left(O_{x}-I\right)\right| \psi_{k}\right\rangle\left.\right|^{2}+2\left|\left(O_{x}-I\right)\right| \psi_{k}\right\rangle|\cdot| O_{x}\left(\left|\psi_{k}^{x}\right\rangle-\left|\psi_{k}\right\rangle\right) \mid,
\end{aligned}
$$

where we have added and subtracted the term $O_{x}\left|\psi_{k}\right\rangle$, and used the inequality $|a+b|^{2} \leq|a|^{2}+|b|^{2}+2|b||a|$.

We will deal with the three terms in this inequality in order. For the first term, because unitary transformations preserve length,

$$
\left.\sum_{x}\left|O_{x}\left(\left|\psi_{k}^{x}\right\rangle-\left|\psi_{k}\right\rangle\right)\right|^{2}=\sum_{x}| | \psi_{k}^{x}\right\rangle-\left.\left|\psi_{k}\right\rangle\right|^{2}=D_{k}
$$

For the second term, we have $\left(O_{x}-I\right)\left|\psi_{k}\right\rangle=-2\left\langle x \mid \psi_{k}\right\rangle|x\rangle$ because $O_{x}|y\rangle=|y\rangle$ if $y \neq x$ and $O_{x}|x\rangle=-2|x\rangle$. Thus,

$$
\left.\sum_{x}\left|\left(O_{x}-I\right)\right| \psi_{k}\right\rangle\left.\right|^{2}=\sum_{x}\left|2\left\langle x \mid \psi_{k}\right\rangle\right|^{2}=4
$$

because $\{|x\rangle\}$ is a basis. Finally, we have

$$
\left.\left.2 \sum_{x}\left|\left(O_{x}-I\right)\right| \psi_{k}\right\rangle|\cdot| O_{x}\left(\left|\psi_{k}^{x}\right\rangle-\left|\psi_{k}\right\rangle\right)\left|=2 \sum_{x}\right|\left\langle x \mid \psi_{k}\right\rangle|2 \cdot|\left|\psi_{k}^{x}\right\rangle-\left|\psi_{k}\right\rangle\right) \mid .
$$

Now, we use the Cauchy-Schwarz inequality, $v \cdot w \leq|v||w|$, on this term. Thus,

$$
\begin{aligned}
\left.2 \sum_{x}\left|2\left\langle x \mid \psi_{k}\right\rangle\right| \cdot| | \psi_{k}^{x}\right\rangle-\left|\psi_{k}\right\rangle \mid & \leq 4 \sqrt{\sum_{x}\left|\left\langle x \mid \psi_{k}\right\rangle\right|^{2}} \sqrt{\left.\sum_{x}| | \psi_{k}^{x}\right\rangle-\left.\left|\psi_{k}\right\rangle\right|^{2}} \\
& =4 \sqrt{D_{k}}
\end{aligned}
$$

We thus have, substituting into the equation for $D_{k+1}$,

$$
D_{k+1} \leq D_{k}+4 \sqrt{D_{k}}+4
$$

We show by induction that this gives $D_{k} \leq 4 k^{2}$. It's clear that $D_{0}=0$. Now, assume that this equation holds true for $D_{k}$. We have

$$
D_{k+1} \leq D_{k}+4 \sqrt{D_{k}}+4 \leq 4 k^{2}+8 k+4=4(k+1)^{2}
$$

and we are done.
The next thing we need to do is show that $D_{t}$ has to be large in order to identify $x$. We will assume that we find $x$ with certainty; you can modify the proof so that this is not necessary. If we find $x$ with certainty, then the $\left|\psi_{t}^{x}\right\rangle$ are othonormal for all $N$ values of $x$. We will then show that $D_{t} \geq 2 N-2 \sqrt{N}$.

Let's look at the equivalent problem where we have $\left|e_{1}\right\rangle,\left|e_{2}\right\rangle,\left|e_{3}\right\rangle, \ldots,\left|e_{N}\right\rangle$, and we want to find the vector $|v\rangle$ that minimizes $\left.\sum_{i=1}^{N}| | e_{i}\right\rangle-\left.|v\rangle\right|^{2}$. It turns out that $|v\rangle=\frac{1}{\sqrt{N}} \sum_{i=1}^{N}\left|e_{i}\right\rangle$, and one can calculate that for this value of $\left.|v\rangle, \sum_{i=1}^{N}| | e_{i}\right\rangle-$ $\left.|v\rangle\right|^{2}=2 N-2 \sqrt{N}$. Since $D_{t} \leq 2 t^{2}$, this shows that $2 t^{2} \geq 2 N-2 \sqrt{N}$, which implies that $t>\sqrt{N}-1$.

We now show that this is indeed the minimum. We have that

$$
\left.\left|\left|e_{i}\right\rangle-\right| v\right\rangle\left.\right|^{2}=\left\langle e_{i} \mid e_{i}\right\rangle+\langle v \mid v\rangle+\left\langle e_{i} \mid v\right\rangle+\left\langle v \mid e_{i}\right\rangle \geq 2-2\left|\left\langle e_{i} \mid v\right\rangle\right|
$$

so

$$
\left.\sum_{i=1}^{N}| | e_{i}\right\rangle-\left.|v\rangle\right|^{2} \geq 2 N-2 \sum_{i}\left|\left\langle e_{i} \mid v\right\rangle\right|
$$

By Cauchy-Schwarz,

$$
\sum_{i=1}^{N}\left|\left\langle e_{i} \mid v\right\rangle\right| \cdot 1 \leq \sqrt{\sum_{i=1}^{N}\left|\left\langle e_{i} \mid v\right\rangle\right|^{2}} \sqrt{\sum_{i=1}^{N} 1}=\sqrt{N}
$$

Thus, we have

$$
\left.\sum_{i=1}^{N}| | e_{i}\right\rangle-\left.|v\rangle\right|^{2} \geq 2 N-2 \sqrt{N}
$$

and we are done.
In fact, you can show that even if you only require the algorithm to succeed half the time, it still must take order $\sqrt{N}$ steps; you can see the proof of this in the Nielsen and Chuang.

