

### 18.435/2.111 Homework 3 Solutions

1a:

$$\sum_{i=1}^k E_i = \sum_{i=1}^k T \Pi_i T^\dagger \quad (1)$$

$$= T \left( \sum_{i=1}^k \Pi_i \right) T^\dagger \quad (2)$$

$$= T T^\dagger \quad (3)$$

$$= I_n \quad (4)$$

**1b:** I have received feedback that this problem was confusing in terms of dimensions, so I will try to be precise. The confusion arises when we try to interpret a *ket*  $|v\rangle$  in terms of the more familiar *vector* or *n-tuple*  $[v_1 \ v_2 \ \cdots \ v_n]^T$  or when we interpret an *operator* as a familiar *matrix*. Both the matrix and the vector representations for operators and kets arise when we define a basis for the space of interest. If we define a basis  $\{|i\rangle\}_{i=1\dots d}$  we can say that  $\mathcal{H}_d$  is the span of this basis. We can further define the space  $\mathcal{H}_n$  as the span of the first  $n$  kets of the basis. If we want to represent  $|v\rangle$  as a vector, we can do it in this basis by defining  $v_i = \langle i|v\rangle$ .

Similarly, for any operator  $A : \mathcal{H}_d \mapsto \mathcal{H}_d$ , we can fully describe it by its action on the basis kets  $\{|i\rangle\}_{i=1}^d$ . We define the *matrix elements*  $A_{ij} = \langle i|A|j\rangle$  and we can always write the operator as  $A = \sum_{ij} A_{ij} |i\rangle \langle j|$ . Let's interpret  $T$  as an operator. We can see from the matrix elements that  $T = \sum_{i=1}^n |i\rangle \langle i|$ .

Now to the problem of interest. We have been told that  $|v\rangle$  is in the subspace generated by the first  $n$  basis vectors, *i.e.*  $v_i = 0$  for  $i > n$ . Viewed in terms of its basis elements, it should be clear that  $T|v\rangle = |v\rangle$ . In that case, we insert  $T$  on either side of  $\Pi_i$  and see the desired equality:

$$\langle v| \Pi_i |v\rangle = \langle v| T \Pi_i T |v\rangle = \langle v| E_i |v\rangle \quad (5)$$

You might complain that I've left off the adjoint  $^\dagger$ . When we interpret  $T$  in bra-ket notation, as above, we can see that  $T = T^\dagger$ .

I can see why this may be confusing. It may be easier to see if we were to rewrite the desired equation as

$$\mathcal{H}_d \langle v| \Pi_i |v\rangle_{\mathcal{H}_d} = \mathcal{H}_n \langle v| E_i |v\rangle_{\mathcal{H}_n} \quad (6)$$

where I labelled with the bras and kets with subscripts indicating which space they live in. In that case,  $T$  is defined as  $T = \sum_{i=1}^n |i\rangle_{\mathcal{H}_n} \langle i|_{\mathcal{H}_d}$ . We

can see that  $T \neq T^\dagger$ . In this form, we should derive the desired result by noting that when  $|v\rangle_{\mathcal{H}_d}$  is spanned by the first  $n$  basis vectors, we have  $|v\rangle_{\mathcal{H}_n} = T^\dagger|v\rangle_{\mathcal{H}_d} = T^\dagger T|v\rangle_{\mathcal{H}_d}$ . Using this, we can get

$$\mathcal{H}_d\langle v|\Pi_i|v\rangle_{\mathcal{H}_d} = \mathcal{H}_d\langle v|T^\dagger T\Pi_i T^\dagger T|v\rangle_{\mathcal{H}_d} = \mathcal{H}_n\langle v|E_i|v\rangle_{\mathcal{H}_n} \quad (7)$$

**1c:** First, let's understand the probabilities associated with the  $k+1$  outcomes of the consecutive projective measurement. The first measurement has two outcomes, namely  $N$  and NOT  $N$ . We can state  $p(N) = \langle\psi|T^\dagger T|\psi\rangle$  and  $p(\text{NOT } N) = \langle\psi|(I - T^\dagger T)|\psi\rangle = 1 - p(N)$ . Now, to derive the probabilities of the subsequent measurement (for if we observed outcome  $N$ ), first let's write down the resultant state after outcome  $N$ :  $(T^\dagger T)/\sqrt{p(N)}|\psi\rangle$ . From here, we can write the conditional probabilities for the  $k$  outcomes associated with  $\Pi_i$ :

$$p(i|N) = \frac{\langle\psi|T^\dagger T\Pi_i T^\dagger T|\psi\rangle}{p(N)}. \quad (8)$$

We use the definition of conditional probability to conclude

$$p(i) = p(i|N)p(N) = \langle\psi|T^\dagger T\Pi_i T^\dagger T|\psi\rangle. \quad (9)$$

Combining these, we can define a POVM by the  $k+1$  operators  $\{I - T^\dagger T, T^\dagger T\Pi_i T^\dagger T\}$ . These form a valid POVM which we can see by summing them to the identity:

$$I - T^\dagger T + \sum_{i=1}^k T^\dagger T\Pi_i T^\dagger T = I - T^\dagger T + T^\dagger T\left(\sum_{i=1}^k \Pi_i\right)T^\dagger T \quad (10)$$

$$= I - T^\dagger T + T^\dagger T \quad (11)$$

$$= I. \quad (12)$$

**2:** Let's first define what we mean by  $E^{1/2}$ . We can see by its construction that  $E$  is a positive matrix (since it is defined as  $E = A^\dagger A$ ), so we know that it can be written as  $E = V^\dagger D V$ , where  $V$  is unitary and  $D$  is a diagonal matrix where the diagonal entries are non-negative real numbers. By convention, we write the diagonal elements of  $D$  in descending order. This is the spectral decomposition of  $E$ . We define  $E^{1/2} \equiv V^\dagger D^{1/2} V$ . (This definition is unique if  $E$  has distinct eigenvalues, but this will not be important for the existence proof we need here.)

Now let's write down the singular value decomposition of  $A = U S W$ , where  $U$  and  $W$  are unitary and  $S$  is a diagonal matrix with non-negative real

diagonal entries. Again, we write the diagonal elements of  $S$  in descending order. By writing  $E = A^\dagger A = W^\dagger S U^\dagger U S W = W^\dagger S^2 W$ , we see that  $S = D^{1/2}$ .

If we write  $A = U D^{1/2} W = U W W^\dagger D^{1/2} W$ , we can define  $U' = W^\dagger U^\dagger$  and we see that  $U' A = W^\dagger U^\dagger U W W^\dagger D^{1/2} W = W^\dagger D^{1/2} W$ . This is what we wanted to show, as the left hand side is  $E^{1/2}$ . (Note, we haven't claimed  $W = V$ . This is only true the eigenvalues of  $E$  are distinct and  $E^{1/2}$  is uniquely defined. This is a technicality that is not too important to our purposes but is certainly worth understanding.)

**3:** Let's start by noting that  $Q^2 = R^2 = S^2 = T^2 = I$ . This follows immediately from the fact that the eigenvalues are  $\pm 1$ . With this, the first result is straight algebra:

$$\begin{aligned}
& (Q \otimes S + R \otimes S + R \otimes T - Q \otimes T)^2 \\
&= ((R + Q) \otimes S + (R - Q) \otimes T)^2 \\
&= (Q + R)^2 \otimes I + (R - Q)^2 \otimes I \\
&\quad + (R + Q)(R - Q) \otimes ST + (R - Q)(R + Q) \otimes TS \\
&= 4I \otimes I + QR \otimes I + RQ \otimes I - QR \otimes I - RQ \otimes I \\
&\quad + (R^2 - Q^2 - RQ + QR) \otimes ST + (R^2 - Q^2 - QR + RQ) \otimes TS \\
&= 4I + QR \otimes ST - RQ \otimes ST - QR \otimes TS + RQ \otimes TS \\
&= 4I + [Q, R] \otimes ST - [Q, R] \otimes TS \\
&= 4I + [Q, R] \otimes [S, T].
\end{aligned}$$

The bound is easiest to see if we define

$$X = Q \otimes S + R \otimes S + R \otimes T - Q \otimes T.$$

We know that the variance of a random variable is non-negative, so  $\langle X^2 \rangle - \langle X \rangle^2 \geq 0$ . We rearrange this and take the square root to get  $\langle X \rangle \leq \langle X^2 \rangle^{1/2}$ . So to finish the proof of the bound, we need to show that  $\langle X^2 \rangle \leq 8$ .

To continue, let's remind ourselves about the meaning of the notation  $\langle A \rangle$ . Above, I used it as the expected value. This is a shorthand referring to the expected value of an observable (i.e. Hermitian operator)  $A$  when the quantum state is  $|\psi\rangle$ :  $\langle A \rangle = \langle \psi | A | \psi \rangle$ . It should be clear from linear algebra that  $\langle A \rangle \leq \lambda_{\max}(A)$  where  $\lambda_{\max}(A)$  is the largest eigenvalue of  $A$ .

Now let's look at  $\lambda_{\max}([Q, R])$ . Writing out the commutator, we can see that  $\lambda_{\max}(QR - RQ) \leq \lambda_{\max}(RQ) + \lambda_{\min}(RQ) \leq 2$ . A similar argument shows that  $\lambda_{\max}([S, T]) \leq 2$ . Finally, since  $\lambda_{\max}(A \otimes B) = \lambda_{\max}(A) \lambda_{\max}(B)$ , we have  $\lambda_{\max}([Q, R] \otimes [S, T]) \leq 4$ .

Combining all of this together,

$$\begin{aligned}\langle X^2 \rangle &= \langle 4I + [Q, R] \otimes [S, T] \rangle \\ &\leq 4 + \lambda_{\max}([Q, R] \otimes [S, T]) \\ &\leq 8.\end{aligned}$$

. This is the desired result.

**4:** See attached diagram.

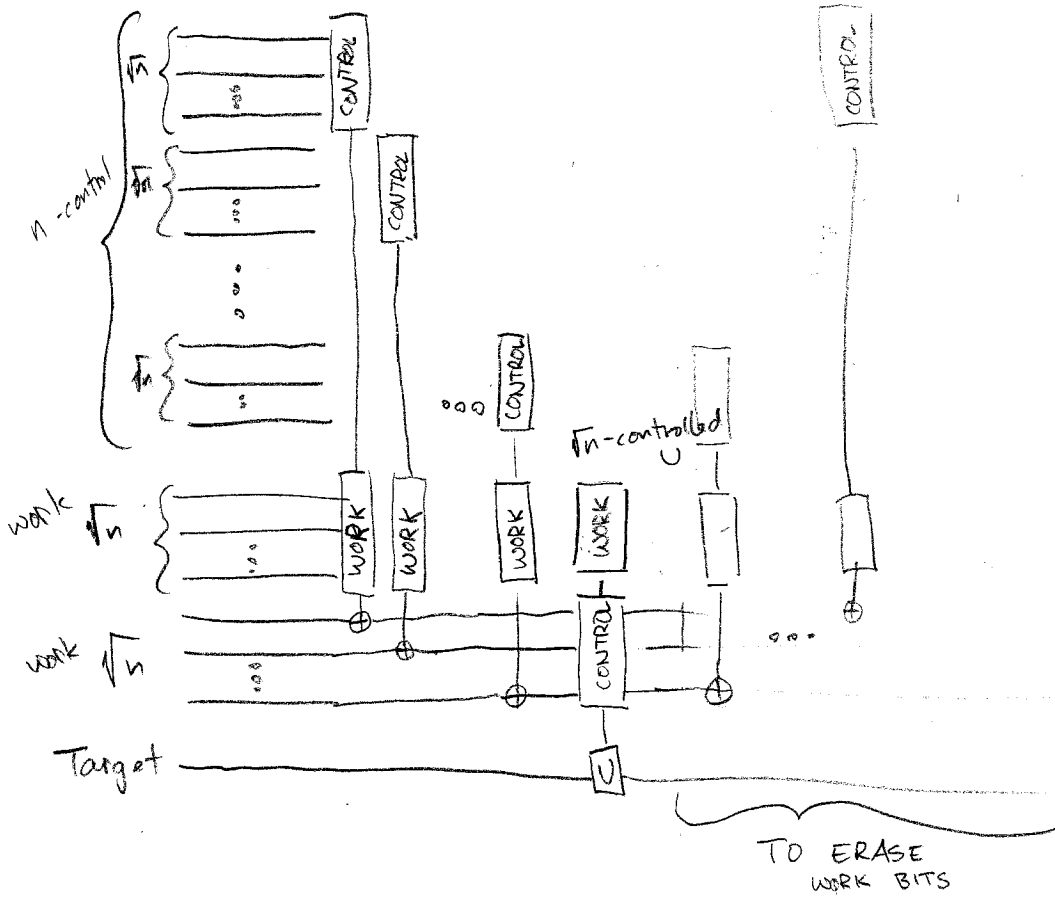
**5:** As mentioned in the hint, we can do this in  $O(\sqrt{n})$  or  $O(\log n)$  work bits, each with  $O(n)$  gates. We'll describe the construction for  $O(\sqrt{n})$ .

We know from class that we can design a  $\sqrt{n}$ -controlled  $U$  gate using  $\sqrt{n}$  work bits and  $O(\sqrt{n})$  gates. We'll need  $2\sqrt{n}$  work bits. The first  $\sqrt{n}$  bits will be used for each  $\sqrt{n}$ -controlled gate. The second  $\sqrt{n}$  bits will be used as follows. Divide the  $n$  control bits into  $\sqrt{n}$  groups of  $\sqrt{n}$  bits each. Using  $\sqrt{n}$ -controlled Nots and the first  $\sqrt{n}$  work bits, set a work bit to one for each set of  $\sqrt{n}$  control bits. When we have done this for each group, we will use the second set of work bits as control for a  $\sqrt{n}$ -controlled  $U$  gate acting on the target bit. Finally, we need to reverse the operations to erase the work bits. Since the control-NOT is its own inverse, we simply repeat the task.

See attached diagram.

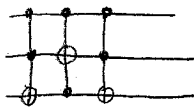
#5

$\sqrt{n}$ -controlled NOTs



#4

(a)



(b)

