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Modular Forms & Spectra

①

Let f be a classical modular form (for $\Gamma(1)$, say); recall that this is a special sort of function

$$f: \mathbb{H} (= SO_2(\mathbb{R}) \backslash SL_2(\mathbb{R})) \longrightarrow \mathbb{C}$$

We can define a function

$$F: SL_2(\mathbb{R}) \longrightarrow \mathbb{C}$$

via

$$F(g) = j(g, i)^{-2k} f(g \cdot i)$$

Two things follow immediately from this definition:

(1) For all $\gamma \in \Gamma(1)$, one has

$$F(\gamma g) = F(g);$$

(2) For all $r(\theta) = \begin{pmatrix} \cos \theta & -\sin \theta \\ \sin \theta & \cos \theta \end{pmatrix} \in K_\infty = SO_2(\mathbb{R})$, one has

$$F(g \cdot r(\theta)) = e^{-2ik\theta} F(g)$$

We can use the same technique to get a function on $SL_2(\mathbb{A})$: strong approximation tells us

$$SL_2(\mathbb{A}) = SL_2(\mathbb{Q}) SL_2(\mathbb{R}) SL_2(\hat{\mathbb{Z}}),$$

where $SL_2(\hat{\mathbb{Z}}) = \prod_p SL_2(\mathbb{Z}_p)$, and so $g \in SL_2(\mathbb{A})$ can be written $g = \gamma g_\infty k_0$, and we can define

$$\begin{aligned} \phi_f: SL_2(\mathbb{A}) &\longrightarrow \mathbb{C} \\ \gamma g_\infty k_0 &\longmapsto F(g_\infty) \end{aligned}$$

Recall: Strong approximation actually says

$$SL_2(\mathbb{A}) = SL_2(\mathbb{Q}) SL_2(\mathbb{R}) K_0$$

whenever $K_0 \subset SL_2(\mathbb{Z})$ is open compact, so we can do this for modular forms of arbitrary weight & level.

(2)

Let's see how this works explicitly. Recall

$$E_{2k}(z) = \sum_{\gamma \in B(\mathbb{Q}) \backslash G(\mathbb{Q})} j(\gamma, z)^{-2k}, \quad k > 1,$$

the normalized Eisenstein series of weight k for $\Gamma(1)$. The associated function on $K_{\infty} = SO_2(\mathbb{R})$ is

$$F(g) = j(g, i)^{-2k} \sum_{\gamma \in B(\mathbb{Q}) \backslash G(\mathbb{Q})} j(\gamma, g \cdot i)^{-2k}$$

Using the cocycle condition

$$j(g_1 g_2, z) = j(g_1, g_2 z) j(g_2, z),$$

one has

$$F(g) = j(g, i)^{-2k} \sum_{\gamma \in B(\mathbb{Q}) \backslash G(\mathbb{Q})} j(\gamma g, i)^{-2k} j(g, i)^{2k}$$

$$= \sum_{\gamma \in B(\mathbb{Q}) \backslash G(\mathbb{Q})} j(\gamma g, i)^{2k}, \quad g \in K_{\infty}.$$

Thus we have our function

$$\begin{aligned} \phi_{\mathbb{Z}k}^E: SL_2(\mathbb{A}) &\longrightarrow \mathbb{C} \\ \gamma g_{\infty} k_0 &\longmapsto \sum_{\gamma \in B(\mathbb{Q}) \backslash G(\mathbb{Q})} j(\gamma g_{\infty}, i)^{-2k} \end{aligned}$$

In general, such "automorphic ^{Eisenstein series} ~~forms~~" will take the form

$$Z(g; s, f) = \sum_{\gamma \in B(\mathbb{Q}) \backslash G(\mathbb{Q})} \text{Im}(\gamma g i)^{s/2} f(\gamma g),$$

where $f = \otimes_v f_v$ is a function on $SL_2(\mathbb{A})$ coming from a particular induced representation. In

our example, $f_v(g_v) = \begin{cases} 1 & v < \infty, \text{ lives in} \\ j(g_v, i)^{-2k} & v = \infty \end{cases}$

$\text{Ind}_{B(\mathbb{A})}^{G(\mathbb{A})}(\mathbb{1})$; we will not elaborate on these points here.

As we briefly mentioned last time, the quantity $s \in \mathbb{C}$ generalizes the "weight" $2k$ of the classical setting. What role does it now play?

If H is a locally compact group, we can consider the space $L^2(H)$ after equipping H with some left Haar measure. We have the right regular representation

$$R: H \xrightarrow{\text{End}} L^2(H),$$

where $(R(y)f)(x) = f(xy)$. The Peter-Weyl theorem says that, if H is compact, then

$$L^2(H) = \hat{\bigoplus}_{(\pi, V)} m(\pi) V,$$

the sum taken over all irreducible unitary representations of H . When H is not compact — say, when $H = SL_2(\mathbb{R}) \backslash SL_2(\mathbb{A})$ — things are more complicated.

Motivating examples: (1) $H = \mathbb{R}/\mathbb{Z}$, compact. The Pontryagin dual is $\hat{H} = \mathbb{Z}$; any irreducible unitary representation is one-dimensional, and corresponds to an integer element of \hat{H} via

$$\begin{aligned} \pi_n: H &\longrightarrow \mathbb{C} \\ e^{i\theta} &\longmapsto e^{in\theta} \end{aligned}$$

As such, Peter-Weyl tells us

$$L^2(\mathbb{R}/\mathbb{Z}) = \hat{\bigoplus}_{n \in \mathbb{Z}} V_{\pi_n}$$

(4)

2. $H = \mathbb{R} = \hat{H}$. This time we don't have Peter-Weyl; nonetheless, elements of $L^2(H)$ are still parameterized by elements of $\mathbb{R} = \hat{H}$. Intuitively, the irreducible unitary representations are

$$x \mapsto e^{2\pi i x y}, \quad y \in \mathbb{R};$$

while not themselves square-integrable, integration against the character $e^{-2\pi i x y}$ yields the Fourier transform isomorphism $L^2(H) \xrightarrow{\sim} L^2(\hat{H})$

We see this time that $L^2(H)$ will not decompose as a direct sum. Instead, we have the direct integral

$$L^2(\mathbb{R}) = \int_{\mathbb{R}}^{\oplus} \mathbb{R}_x dx,$$

a family of vector spaces \mathbb{R}_x indexed by $x \in \mathbb{R}$. Here, \mathbb{R}_x is the unitary representation space

$$\begin{aligned} (H, \mathbb{R}) &\longrightarrow \mathbb{R} \\ (y, z) &\longmapsto e^{2\pi i x y} z \end{aligned}$$

In general, we expect the decomposition of $L^2(H)$ to be a mix of these discrete (direct sum) and continuous (direct integral) parts; we will write

$$L^2(H) = L^2_{\text{disc}}(H) \oplus L^2_{\text{cont}}(H).$$

It happens that, in the case $H = \text{SL}_2(\mathbb{Q}) \backslash \text{SL}_2(\mathbb{A})$, the Eisenstein series

$$\Sigma(g; s, f) = \sum_{\gamma \in \mathbb{B}(\mathbb{Q}) \backslash \mathbb{G}(\mathbb{Q})} \text{Im}(\gamma g, i)^{s-1/2} f(\gamma g)$$

parameterize the "continuous spectrum" of $L^2(\text{SL}_2(\mathbb{Q}) \backslash \text{SL}_2(\mathbb{A}))$, much as real numbers $\gamma \in \mathbb{R}$ do for $L^2(\mathbb{R})$.

Cusp forms complete the picture by describing $L^2_{disc}(SL_2(\mathbb{Q}) \backslash SL_2(\mathbb{A}))$; we sketch the reasons why.

First of all: let $\mathcal{L}_0^{(SL_2(\mathbb{Q}) \backslash SL_2(\mathbb{A}))}$ denote the ^{Hilbert} space of all functions $f: SL_2(\mathbb{A}) \rightarrow \mathbb{C}$ satisfying:

- 1. $f(\gamma g) = f(g)$ for $\gamma \in SL_2(\mathbb{Q})$;
- 2. $\int_{G(\mathbb{Q}) \backslash G(\mathbb{A})} |f(g)|^2 dg < \infty$; and

3. For all unipotent radicals U of all proper parabolic subgroups of $SL_2(\mathbb{A})$, one has

$$\int_{U(\mathbb{Q}) \backslash U(\mathbb{A})} f(ug) du = 0 \quad \text{or} \quad \int_{\mathbb{Q} \backslash \mathbb{A}} f\left(\begin{pmatrix} 1 & x \\ & 1 \end{pmatrix} g\right) dx = 0$$

We call this the space of (automorphic) cuspidal forms of $\mathbb{Q} \backslash SL_2(\mathbb{A})$. One can show that, if $f \in S_{2k}$, then $\phi_f \in \mathcal{L}_0(SL_2(\mathbb{Q}) \backslash SL_2(\mathbb{A}))$.

Ex

Given $\phi \in \mathcal{L}_0$, let V_ϕ be the Hilbert space generated by all right-translates $\{x \mapsto \phi(xg) : g \in G(\mathbb{A})\}$.

An action of $G(\mathbb{A})$ on V_ϕ is also given

$$(g \cdot f)(x) = f(xg),$$

making V_ϕ a unitary representation of $G(\mathbb{A})$, called a cuspidal representation. The collection

$L^2_{cusp}(SL_2(\mathbb{Q}) \backslash SL_2(\mathbb{A}))$ of all cuspidal representations sits inside $L^2_{disc}(SL_2(\mathbb{Q}) \backslash SL_2(\mathbb{A}))$.

