It is interesting to compare this result with corresponding 2D solution

$$u_i(\boldsymbol{x}) = J_{ij}(\boldsymbol{x})F_j \ , \qquad p = \frac{F_j x_j}{2\pi|\boldsymbol{x}|^2} + \partial_\infty \ , \qquad \boldsymbol{x} = (x, y) \tag{5.51a}$$

where

$$J_{ij}(\boldsymbol{x}) = \frac{1}{4\pi\mu}\left[-\delta_{ij}\ln\left(\frac{|\boldsymbol{x}|}{a}\right) + \frac{x_i x_j}{|\boldsymbol{x}|^2}\right] \tag{5.51b}$$

with a being an arbitrary constant fixed by some intermediate flow normalization condition. Note that (5.51) decays much more slowly than (5.50), implying that hydrodynamic interactions in 2D freestanding films are much stronger than in 3D bulk solutions.