

Fractals

(1)

↳ complex geometric shapes with fine structure at arbitrarily small scales [some exhibit self-similarity]

Aim: to understand / describe strange attractors, e.g. fractal dimension.

Countable and uncountable sets

- Sets X and Y are said to have the same cardinality (number of elements)

if there is an invertible mapping that pairs each element of X with precisely one element of Y → one-to-one correspondence

- If a set X can be put into a one-to-one correspondence with the natural numbers $\mathbb{N} = \{1, 2, 3, \dots\}$ then X is said to be countable.

Otherwise, X is uncountable.

↳ in other words, can we list the elements of X as $X = \{x_1, x_2, x_3, \dots\}$.

Example: The set of even numbers $E = \{2, 4, 6, \dots\}$ is countable as we can define an invertible mapping for each $n \in \mathbb{N}$ with $2n \in E$.

Example: The integers are countable as we can express them as the list $\{0, 1, -1, 2, -2, 3, -3, \dots\}$.

Example: To show that the positive rationals are countable, we make a table such that the (p,q) entry is p/q

i.e.	$\frac{1}{1}$	$\rightarrow \frac{1}{2}$	$\rightarrow \frac{1}{3}$	$\rightarrow \frac{1}{4}$...
	$\frac{2}{1}$	$\frac{2}{2}$	$\frac{2}{3}$	$\frac{2}{4}$...
	$\frac{3}{1}$	$\frac{3}{2}$	$\frac{3}{3}$	$\frac{3}{4}$...
	
	

We can weave through all the elements of the table, meaning that each positive rational appears after a finite number of steps, i.e. we have a list.

Example: Show that the set of real numbers between 0 and 1 is uncountable, (2)

→ Proof by contradiction: Suppose the set were countable with elements $\{x_1, x_2, \dots\}$.

In decimal form: $x_1 = 0.x_{11}x_{12}x_{13} \dots$

$x_2 = 0.x_{21}x_{22}x_{23} \dots$

$x_3 = 0.x_{31}x_{32}x_{33} \dots$

Aim: construct $r \in [0, 1]$ that is not on the list. ∵ the list is incomplete!

{ Define $r = 0.\bar{x}_{11}\bar{x}_{22}\bar{x}_{33} \dots$

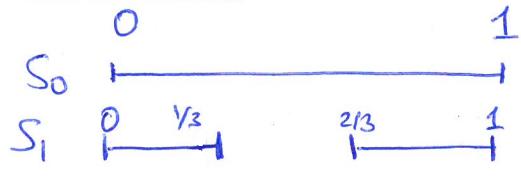
where \bar{x}_{nn} is any digit other than x_{nn} .

- Hence, $r \neq x_1$ as first digit differs

... $r \neq x_n$ as n^{th} digit differs!

i.e. r not on the list \Rightarrow uncountable set.

Cantor set - a fractal!



← delete the middle third from each interval.

$S_\infty = C$. ← Cantor set:

← an infinite number of infinitesimal pieces separated by gaps of various sizes.

Fractal properties of C.

① structure at arbitrarily small scales - as we zoom in on C, we see more and more structure.

② self-similarity - it contains smaller copies of itself at all scales
i.e. the left portion of S_{n+1} (or right portion) looks like S_n ,
but scaled by a factor of 3.

③ non-integer dimension! (see later)

Non-fractal properties of C.

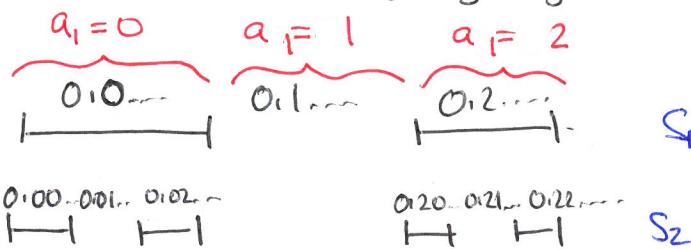
① C has measure zero, i.e. it can be covered by intervals whose total length is arbitrarily small.

↳ Why? C can be covered by each of the sets S_n , where S_n has length $L_n = \left(\frac{2}{3}\right)^n$.
 But $L_n \rightarrow 0$ as $n \rightarrow \infty$, ∴ C has total length zero.

② C is uncountable.

First note that C contains all points $x \in [0,1]$ that have no ones in their base 3 expansion.

Note: $x \in [0,1]$ in base 3 is $x = \frac{a_1}{3} + \frac{a_2}{3^2} + \frac{a_3}{3^3} + \dots$ where $a_j \in \{0,1,2\}$.



Note: the $a_n=1$ entries are deleted at each stage

Note: what about endpoints like $0.\overline{1} = \frac{1}{3}$?

Write $0.\overline{1} = 0.0222\dots$ ← no ones appear in base 3 expansion

↳ Application: apply the diagonal argument to a list $\{c_1, c_2, \dots\}$ of all points in C → contradiction. Write c_i with digits c_{ij} .

Define $\bar{c} = 0.\bar{c}_{11} \bar{c}_{22} \bar{c}_{33} \dots$

↳ where $\bar{c}_{nn} = \begin{cases} 0 & \text{if } c_{nn} = 2 \\ 2 & \text{if } c_{nn} = 0 \end{cases}$

Hence, \bar{c} is in C as it is written in terms of zeros and twos, but \bar{c} is not on the list $\cancel{\cancel{\cancel{\quad}}}$

Hence, C is uncountable \square

Dimension of self-similar fractals.

(see later for non-self-similar case)

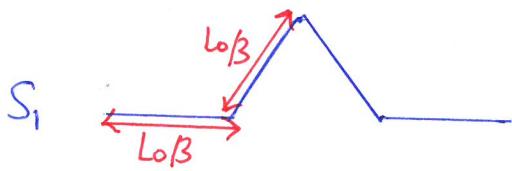
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↳ the definition for familiar geometric objects breaks down

- we cannot easily define a fractal dimension as the minimum number of coordinates needed to describe every point in the set

Example: von Koch curve.

S_0 —————— length L_0



each edge segment has length $L_0/3$

S_2 —————— ↓

each length is $L_0/3^2$.

To get S_n , replace the middle third of each line segment of S_{n-1} by an equilateral triangle

$S_0 = K$: von Koch curve.

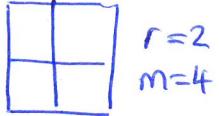
↳ Arclength: $l_1 = \frac{4}{3}l_0$, etc $\Rightarrow l_n = \left(\frac{4}{3}\right)^n l_0 \rightarrow \infty$ as $n \rightarrow \infty$. \Rightarrow Infinite Arclength
 \Rightarrow length between any two points is infinite! Hence, K appears more than one-dimensional

Similarity dimension

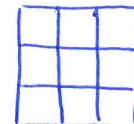
• Consider a square of length l

↳ to reconstruct the square out of m squares each of length l/r (where r is an integer), we need $m=r^2$

e.g.



$$r=2 \\ m=4$$



$$r=3 \\ m=9$$

Note: $m=r^d$ where $d=2$

• For a cube, we need $m=r^3$ cubes of length l/r to reconstruct the original cube ($d=3$)

• To generalize this definition to non-integers, we consider self-similar set that is composed of m copies of itself scaled down by a factor of r

Then the similarity dimension d satisfies $m=r^d$, i.e. $d = \frac{\log m}{\log r}$

Example: Cantor Set.

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C is composed of two copies of itself, each scaled down by a factor of 3

$$\therefore m=2, r=3 \Rightarrow d = \frac{\log 2}{\log 3} = 0.63\dots$$

Example: von Koch curve.

K is composed of four pieces,^{of itself} each scaled down by a factor of 3

$$\therefore m=4, r=3 \Rightarrow d = \frac{\log 4}{\log 3} = 1.26\dots \quad [\text{Note: } d>1, \text{ as anticipated}]$$

Example: even-fifths Cantor set.

So

$$S_1: 0 \quad \frac{1}{5} \quad \frac{2}{5} \quad \frac{3}{5} \quad \frac{4}{5} \quad 1$$

$$S_2: \dots \quad \dots \quad \dots \quad \dots \quad \dots$$

S_2 is ^{m=3} copies of S_1 , each scaled by $r=5$ $\Rightarrow d = \frac{\log 3}{\log 5} = 0.68\dots$

Cantor set - abstraction!

A closed set S is called a topological Cantor set if it satisfies

- ① S is "totally-disconnected", i.e. all points in S are separated from each other.
(S contains no intervals)
- ② S has no "isolated points": for any $p \in S$ and $\forall \varepsilon > 0$, $\exists q \in S$ ($q \neq p$) such that q lies within a distance ε of p.

↳ Paradox: points are spread apart but also packed together!

↳ Note: these properties are topological rather than geometric, so no mention of self-similarity.

→ Topological Cantor Sets often appear as strange attractors!

Box dimension → for non-self-similar fractals.

↳ measure the set in a way that ignores irregularities of size less than $\varepsilon > 0$, and then consider the dependence as $\varepsilon \rightarrow 0$.

Definition of Box dimension.

Let S be a subset of a D -dimensional Euclidean space, and let $N(\varepsilon)$ be the minimum number of D -dimensional cubes of size ε needed to cover S .

Define the box dimension d so that $N(\varepsilon) \propto \varepsilon^{-d}$ as $\varepsilon \rightarrow 0$, (power law)

$$\text{i.e. } d = \lim_{\varepsilon \rightarrow 0} \frac{\log N(\varepsilon)}{\log(1/\varepsilon)} \quad [\text{if the limit exists}]$$

e.g. for a smooth curve of length L ; $N(\varepsilon) \propto L/\varepsilon$ ($d=1$)

e.g. for a planar region of area A bounded by a smooth curve; $N(\varepsilon) \propto A/\varepsilon^2$ ($d=2$)

Example: Cantor set.

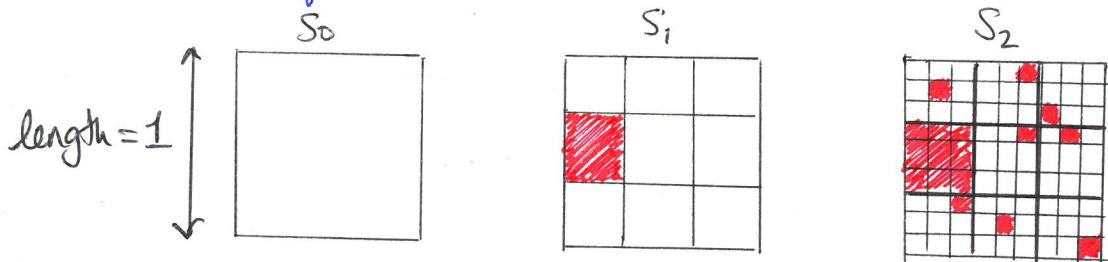
Note: C is covered by each set S_n , where S_n has 2^n intervals each of length 3^{-n}

∴ if $\varepsilon = 3^{-n}$, we need $N = 2^n$ intervals to cover the Cantor set

$$\Rightarrow d = \lim_{\varepsilon \rightarrow 0} \frac{\log N}{\log(1/\varepsilon)} = \lim_{n \rightarrow \infty} \frac{\log(2^n)}{\log(3^{-n})} = \frac{n \log 2}{-n \log 3} = \frac{\log 2}{\log 3} \quad [\text{agrees with similarity dimension}]$$

Example: a non-self-similar fractal.

Construction: Divide a square region into 9 equal squares and discard one of the squares at random. Then repeat the process on each of the remaining squares.



Note: S_1 covered by $N=8$ squares of side $\varepsilon = 1/3$ } no wastage!
 S_2 $N=8^2$ $\varepsilon = 1/3^2$
 S_n $N=8^n$ $\varepsilon = 3^{-n}$

$$\Rightarrow d = \lim_{\varepsilon \rightarrow 0} \frac{\log N(\varepsilon)}{\log(1/\varepsilon)} = \lim_{n \rightarrow \infty} \frac{\log(8^n)}{\log(3^{-n})} = \frac{\log 8}{\log 3}$$

Downsides of Box Dimension.

- difficult to find a minimal cover
- computation requires too much storage space
- some mathematical irregularities, e.g. set of rational numbers between 0 and 1 has a box dimension of 1, yet only has countably many points.

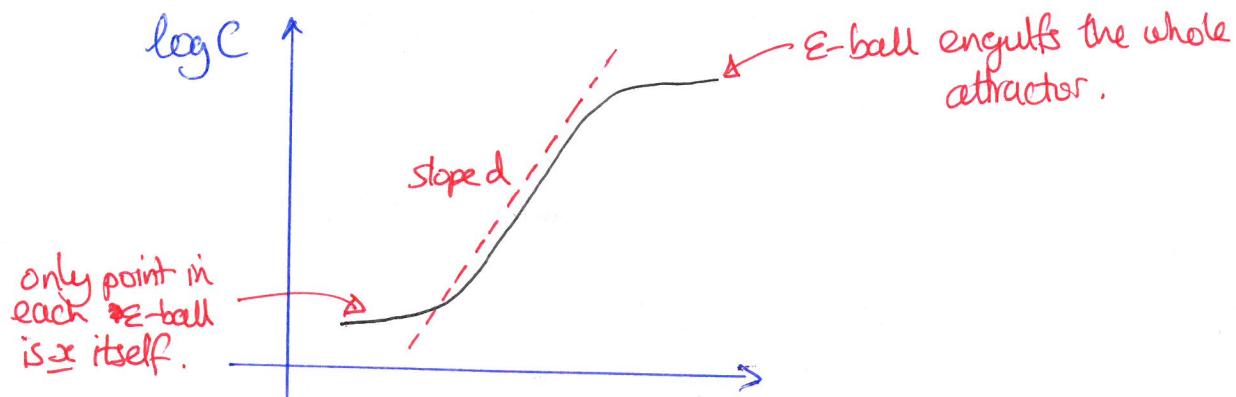
Alternative: Hausdorff dimension - coverings of small sets with varying size, not equal size.
 ↳ nicer mathematically, but even harder to compute numerically.

Pointwise and Correlation Dimensions → how to characterize strange attractors!

Consider a chaotic system that settles down to a strange attractor in phase space.
 ↳ if the strange attractor exhibits fractal structure, what is the fractal dimension?

Method: Fix a point \underline{x} on the attractor A

- Let $N_{\underline{x}}(\varepsilon)$ denote the number of points on A inside a ball of radius ε about \underline{x} .
 the points \underline{x}_i are generated from a simulation, discarding initial transients
- As most of the points in the ball come from earlier/later portions of the trajectory, $N_{\underline{x}}(\varepsilon)$ measures how frequently a typical trajectory visits an ε neighbourhood of \underline{x} .
- As ε increases, $N_{\underline{x}}(\varepsilon)$ grows as a power law $N_{\underline{x}}(\varepsilon) \propto \varepsilon^d$, where d is the pointwise dimension at \underline{x} ,
- To account for variations in \underline{x} , we average $N_{\underline{x}}(\varepsilon)$ over many \underline{x} , where the resulting quantity $C(\varepsilon) \propto \varepsilon^d$ and d is the correlation dimension.



∴ we only expect the power law to hold in the scaling region
 [minimum separation of points in A] $\ll \varepsilon \ll$ [diameter of A]

Correlation vs. Box dimension.

- Correlation dimension takes account of the density of points of the attractor
- Box dimension weighs all boxes equally, no matter how many points they contain.
- In general, $d_{\text{correlation}} \leq d_{\text{box}}$, but they are usually close.

Example (Lorenz attractor - standard parameter values)

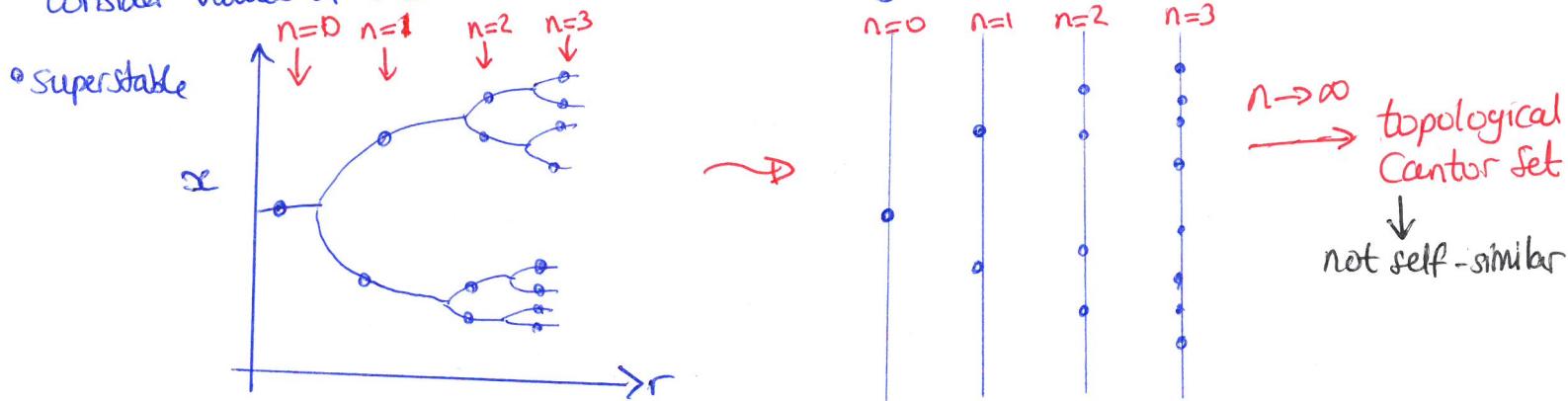
$$d_{\text{corr}} = 2.05 \pm 0.01 \quad \text{non-integer!}$$

↳ estimated to within 5% using only a few thousand points. ∴ rapid convergence!

Example (logistic map).

$$x_{n+1} = r x_n (1 - x_n), \quad r = r_\infty \text{ (onset of chaos)}$$

Consider values of r and x at which each 2^n -cycle is superstable



Find that $d_{\text{corr}} = 0.500 \pm 0.005$ at $r = r_\infty$.

Note: $d_{\text{box}} \approx 0.538 \geq d_{\text{corr}}$, as expected.

Note: At the onset of chaos, one might obtain residual correlations along a single trajectory

Multifractals.

↳ when the dimension varies across the attractor.

- For a multifractal A , let S_α be the subset of A consisting of all points with pointwise dimension α . ↳ S_α is itself a fractal with fractal dimension $f(\alpha)$.

- an interwoven set of fractals on different dimensions α , where $f(\alpha)$ determines their weight.
- useful for characterizing systems at the onset of chaos.

