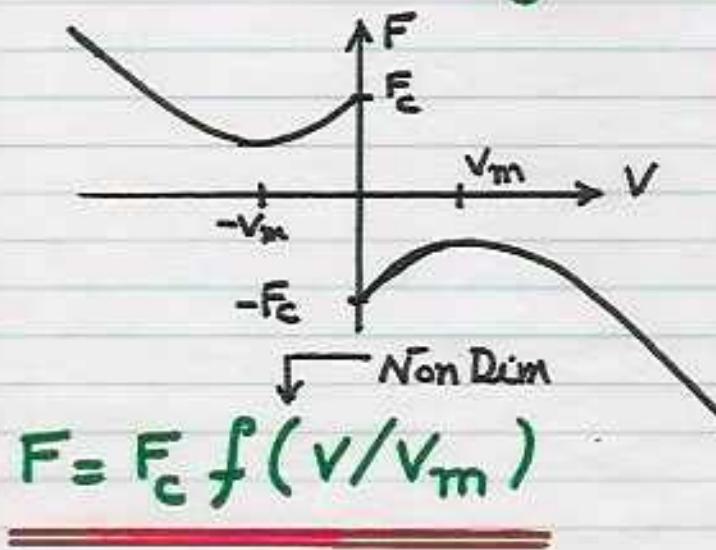


Relaxation oscillations due to friction



Friction function as function of velocity. The existence of a static critical load F_c and a critical velocity v_m are important. $x=0$: spring relaxed

Equation $M \ddot{x} = -kx + F(\dot{x} - U)$

Equilibrium $x = x^*$, where $kx^* = F(-U)$

so, write $x = x^* + X$

Then $M \ddot{X} = -kX + F(\dot{X} - U) - kx^*$

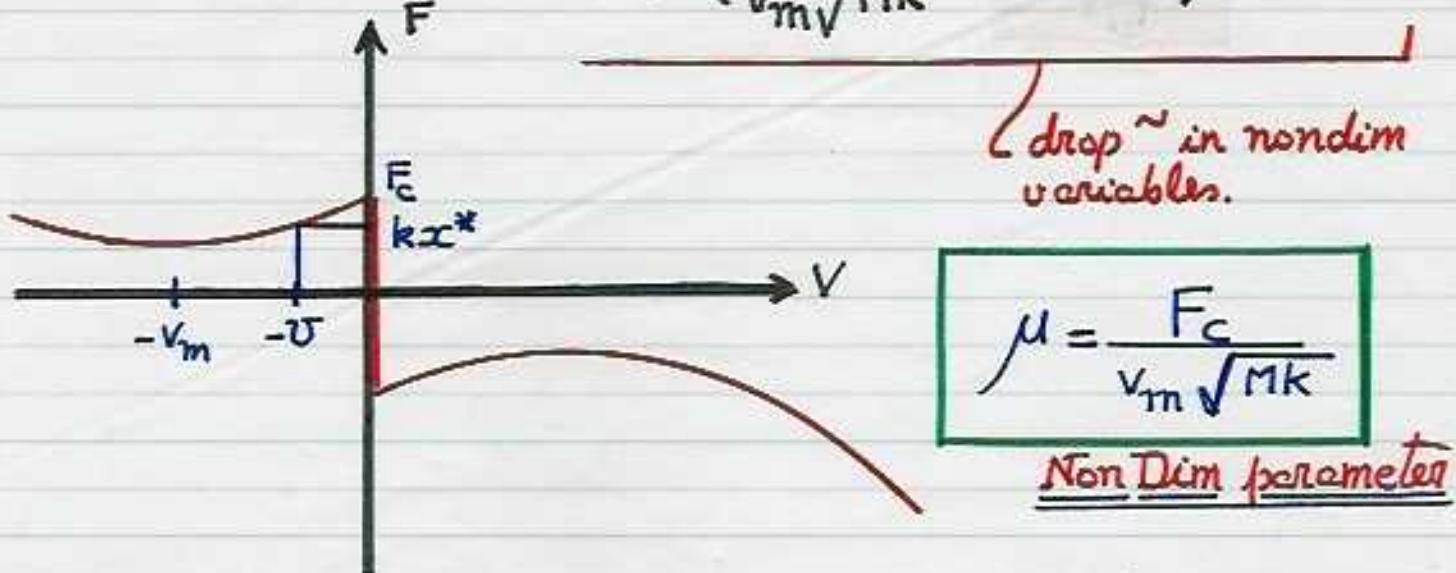
Assume now: $0 < U < v_m$, $U = O(v_m)$

and nondim. $x = (F_c/k) \tilde{x}$; $x^* = \frac{F_c}{k} \tilde{x}^*$,

$t = \sqrt{\frac{M}{k}} \tilde{t}$; $U = v_m \tilde{U}$

mass-spring response ^{time.} ✓

Then: $\ddot{x} = -x + f\left(\frac{F_c}{v_m \sqrt{Mk}} \dot{x} - U\right) - x^*$

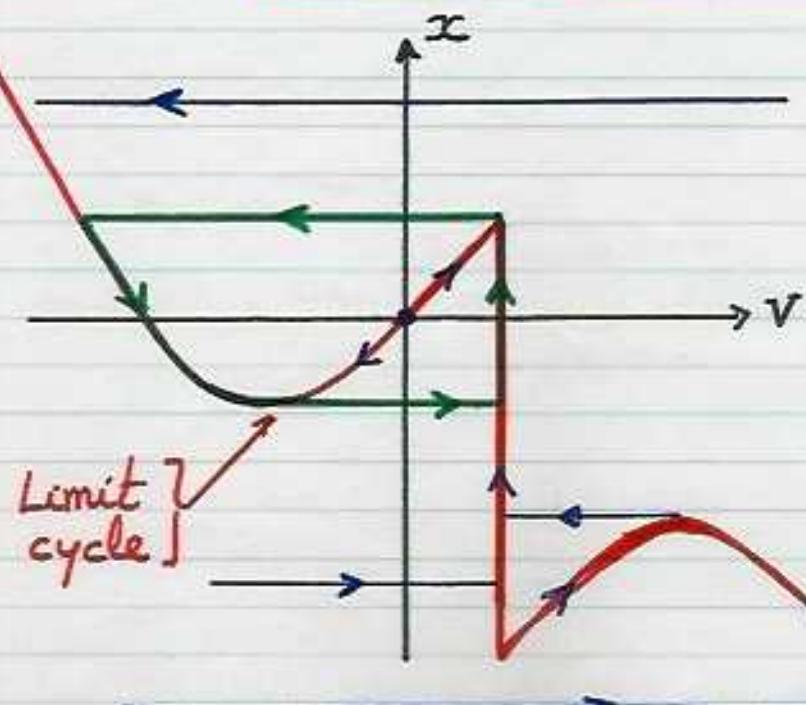


$$\mu = \frac{F_c}{v_m \sqrt{Mk}}$$

Non Dim parameters

In non-dim version: $F_c = 1, v_m = 1, F = f$
 $k = 1, 0 < U < 1.$

Let now $g(v) = f(v-U) - x^*$



$$\ddot{x} + x = g(\mu \dot{x})$$

$$\begin{cases} \dot{x} = \frac{1}{\mu} V \\ \dot{V} = \mu(g(V) - x) \end{cases}$$

For $\mu \gg 1$

$$x = g(v)$$

Thus, the conditions under which a relaxation limit cycle occurs are:

$$0 < \bar{v} < v_m \quad \left\{ \begin{array}{l} \text{Mean pushing} \\ \text{velocity in} \\ \text{critical range} \end{array} \right.$$

and

$$F_c \gg v_m \sqrt{Mk}$$

This seems to indicate that the mass cannot be very large. However $F_c \approx M a$

$\left. \begin{array}{l} \text{a typical} \\ \text{acceleration} \end{array} \right\} \uparrow$

So get $\underline{\underline{\sqrt{M} \gg v_m \sqrt{k/a}}}$

F_c depends on weight!