

18.155 LECTURE 3: 12 SEPTEMBER, 2013

- We showed $\mathcal{F} : \mathcal{S}(\mathbb{R}^n) \rightarrow \mathcal{S}(\mathbb{R}^n)$ is continuous, in fact

$$\|\hat{\phi}\|_N \leq C\|\phi\|_{N+n+1}.$$

- To prove it is an isomorphism we start with two Lemmas – the first is very standard

Lemma 1. *There exists $\psi \in \mathcal{S}(\mathbb{R}^n)$, $\psi(x) \geq 0$, $\psi(x) = 1$ in $|x| < \frac{1}{2}$ $\psi(x) = 0$ in $|x| > 1$.*

Proof. In one variable consider the function

$$\mu(x) = \begin{cases} \exp(-1/x) & x > 0 \\ 0 & x \leq 0. \end{cases}$$

This is an example of a non-analytic but infinitely differentiable function. It is certainly smooth in $x > 0$ and the derivatives are of the form

$$(1) \quad \frac{d^k}{dx^k} \mu(x) = \frac{p_k(x)}{x^{2k}} \exp(-1/x),$$

with p_k a polynomial. The convergence of the power series for e^s , all terms in which are positive for $s > 0$ shows that for each N , $e^s \geq s^N/n!$ – and hence

$$\frac{d^k}{dx^k} \mu(x) \leq C_{k,N} x^N \text{ in } 0 < x \leq 1.$$

Thus all the derivatives, defined in $x > 0$, extend by continuity down to 0 where they vanish to infinite order. This includes μ itself which is therefore given by the integral from -1 of its own derivative extended to be 0 in $x < 0$. This argument iterates to show – or you could do it directly anyway – that μ is infinitely differentiable across 0.

From this we can construct

$$(2) \quad \psi'(x) = \mu(1 - |x|^2) \in \mathcal{S}(\mathbb{R}^n), \quad \eta(x) = \mu(|x|^2 - \frac{1}{2}) \in \mathcal{C}^\infty(\mathbb{R}^n)$$

which are respectively positive in $|x| < 1$ but zero in $|x| \geq 1$ and zero near $x = 0$ and positive in $|x| > 3/4$. Thus $\psi + \eta$ is strictly positive everywhere and then

$$(3) \quad 0 \leq \psi(x) = \frac{\psi'(x)}{\psi'(x) + \eta(x)} \in \mathcal{S}(\mathbb{R}^n)$$

satisfies the conditions we want, since it is equal to 1 where $\eta(x) = 0$ and 0 where $\psi'(x) = 0$. \square

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Lemma 2. *Any $\phi \in \mathcal{S}(\mathbb{R}^n)$ can be written in the form*

$$(4) \quad \phi(x) = \phi(0) \exp(-\frac{|x|^2}{2}) + \sum_{j=1}^n x_j \psi_j(x), \quad \psi_j \in \mathcal{S}(\mathbb{R}^n).$$

Proof. Taylor's formula and then use of the cutoff from the previous lemma. \square

- Now, to back to the Forier transform. We consider the composite and then restrict to 0 and claim that

$$(5) \quad \mathcal{G}(\hat{\phi})(0) = c\phi(0) \quad \forall \phi \in \mathcal{S}(\mathbb{R}^n)$$

where c is a fixed constant independent of ϕ . Indeed to see this insert (4) into the left side to get

$$(6) \quad \mathcal{G}(\hat{\phi})(0) = c\phi(0) + \sum_j \mathcal{G}(\widehat{x_j \psi_j})(0), \quad c = \mathcal{G}(\hat{\gamma})(0), \quad \gamma = \exp(-\frac{|x|^2}{2}).$$

However, $\widehat{x_j \psi_j}(\xi) = i\partial_{\xi_j} \hat{\psi}_j(\xi)$ as we showed last time and since

$$(7) \quad \mathcal{G}(f)(0) = (2\pi)^{-n} \int f,$$

so each of the terms in the sum vanishes. Thus we arrive at (5) where we even no the constant in terms of the Gaussian.

- Now, we can also work out formulæ for the Fourier transform of translates and multiples by exponentials:-

$$(8) \quad \mathcal{F}(\phi(\bullet + y))(\xi) = e^{iy \cdot \xi} \hat{\psi}, \quad \mathcal{G}(e^{-iy \cdot \bullet} f)(x) = \mathcal{G}(f)(x + y).$$

Combining these two and (5) shows that

$$(9) \quad \mathcal{G}\mathcal{F} = c \text{Id}.$$

- So we are reduced to working out the constant. This amounts to working out the Fourier transform of the Gaussian and this is pretty standard. First we only need do the 1-D case since

$$(10) \quad \mathcal{F}(\exp(-\frac{|x|^2}{2})) = \prod_j \mathcal{F}(\exp(-x_j^2/2))$$

and then we can check that

$$(11) \quad \left(\frac{d}{d\xi} + \xi\right)\mathcal{F}(\exp(-x^2/2))(\xi) = 0 \implies \mathcal{F}(\exp(-x^2/2))(\xi) = c' \exp(-\xi^2/2)$$

where the constant is the value of the Fourier transform at 0, i.e.

$$(12) \quad c' \int_{\mathbb{R}} \exp(-x^2/2) = \sqrt{2\pi}.$$

Going back to (5) it follows that $c = 1$ and we have the Fourier inversion formula

$$(13) \quad \mathcal{G} \circ \mathcal{F} = \text{Id} = \mathcal{F} \circ \mathcal{G} \text{ on } \mathcal{S}(\mathbb{R}^n)$$

where the second follows from the first by changing signs.

- So the Fourier transform is an isomorphism on $\mathcal{S}(\mathbb{R}^n)$, a continuous linear bijection with a continuous inverse.
- Computing with absolutely convergent Lebesgue integrals it also follows directly that for any $\phi, \psi \in \mathcal{S}(\mathbb{R}^n)$,

$$(14) \quad \int_{\mathbb{R}^n} \hat{\phi} \psi = \int_{\mathbb{R}^{2n}} e^{-ix \cdot \xi} \phi(x) \psi(\xi) = \int_{\mathbb{R}^n} \phi \hat{\psi}.$$

This gives a weak formulation of the Fourier transform which we can write

$$(15) \quad I(\hat{\phi})(\psi) = I(\phi)(\hat{\psi})$$

and so for general $u \in \mathcal{S}'(\mathbb{R}^n)$ it is consistent to *define*

$$(16) \quad \hat{u}(\psi) = u(\hat{\psi}) \quad \forall u \in \mathcal{S}'(\mathbb{R}^n), \psi \in \mathcal{S}(\mathbb{R}^n).$$

This gives a linear bijection (so far we don't have a topology to measure continuity)

$$(17) \quad \mathcal{F} : \mathcal{S}'(\mathbb{R}^n) \longrightarrow \mathcal{S}'(\mathbb{R}^n).$$

- If we apply the identity (14) to $\psi = \bar{\eta}$ for some $\eta \in \mathcal{S}(\mathbb{R}^n)$ then $\hat{\eta} = \bar{\psi}$ so by the inversion formula,

$$\eta(\xi) = \mathcal{G}(\bar{\psi}) = (2\pi)^{-n} \mathcal{F}(\bar{\psi})(-\xi) \implies \bar{\eta}(\xi) = (2\pi)^{-n} \hat{\psi}(\xi).$$

The result is Parseval's formula

$$(18) \quad \int_{\mathbb{R}^n} \hat{\phi} \bar{\eta} = (2\pi)^{-n} \int_{\mathbb{R}^n} \phi \bar{\eta}$$

the essential unitarity of the Fourier transform, i.e. $(2\pi)^{-n/2} \mathcal{F}$ preserves the L^2 inner product on $\mathcal{S}(\mathbb{R}^n)$ and so, if we accept that $\mathcal{S}(\mathbb{R}^n)$ is dense in $L^2(\mathbb{R}^n)$,

Proposition 1. *The operators $(2\pi)^{-n/2} \mathcal{F}$ extends by continuity to be unitary on $L^2(\mathbb{R}^n)$.*

- So we have seen that the Fourier transform is an isomorphism on $\mathcal{S}(\mathbb{R}^n)$, and on $\mathcal{S}'(\mathbb{R}^n)$ which restricts to an isomorphism on the subspace $L^2(\mathbb{R}^n)$.
- The L^2 -based Sobolev spaces are then defined for each *real* number $s \in \mathbb{R}$:

$$(19) \quad H^s(\mathbb{R}^n) = \{u \in \mathcal{S}'(\mathbb{R}^n); (1 + |\xi|^2)^{s/2} \hat{u} \in L^2(\mathbb{R}^n)\}.$$

Note that in the homework this week you showed that $(1 + |x|^2)^{s/2}$ is a 'multiplier' on $\mathcal{S}(\mathbb{R}^n)$, and hence on $\mathcal{S}'(\mathbb{R}^n)$ so the definition makes sense, $L^2(\mathbb{R}^n)$ being a well-defined subspace of $\mathcal{S}'(\mathbb{R}^n)$.

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Lemma 3. *For $s = k \in \mathbb{N}$ a positive integer,*

$$(20) \quad H^k(\mathbb{R}^n) = \{u \in L^2(\mathbb{R}^n); D^\alpha u \in L^2(\mathbb{R}^n) \subset \mathcal{S}'(\mathbb{R}^n), |\alpha| \leq k\}$$

is also the space of L^2 functions with strong derivatives up to order k in L^2 where successive strong derivatives are defined by convergence of the difference quotient in L^2 :

$$(21) \quad \partial_j u = \lim_{s \rightarrow 0} \frac{u(x + se_j) - u(x)}{s} \text{ in } L^2(\mathbb{R}^n).$$