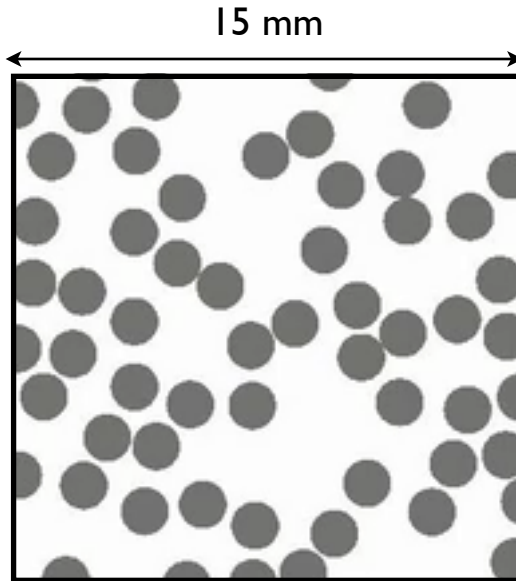


L04:
Continuum theory
&
diffusion

18.354

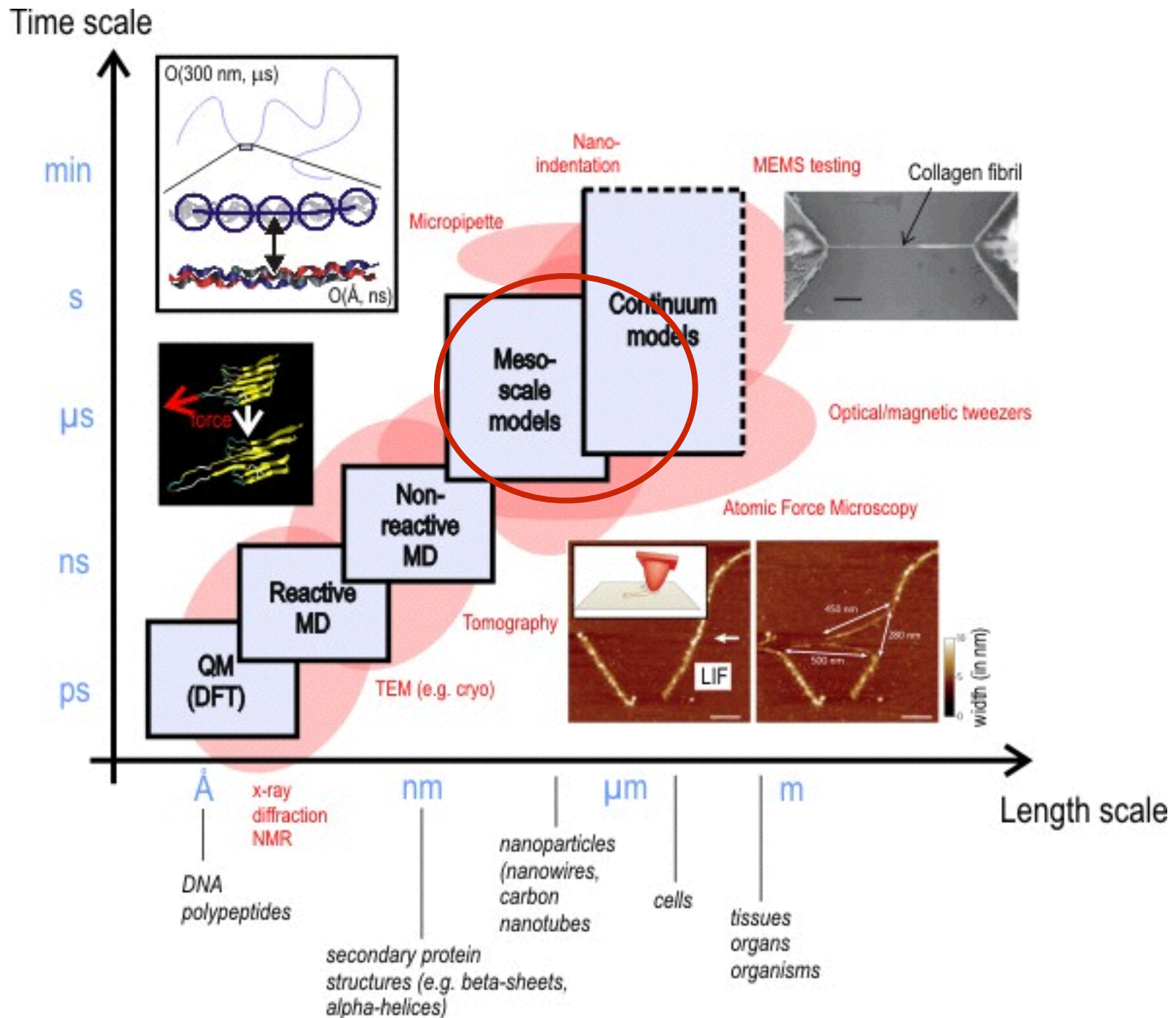
Last class: Hamiltonian dynamics



Dilute Fluid

$$\phi = 0.34$$

Molecular Dynamics Simulations



How do we describe systems with large N of particles?

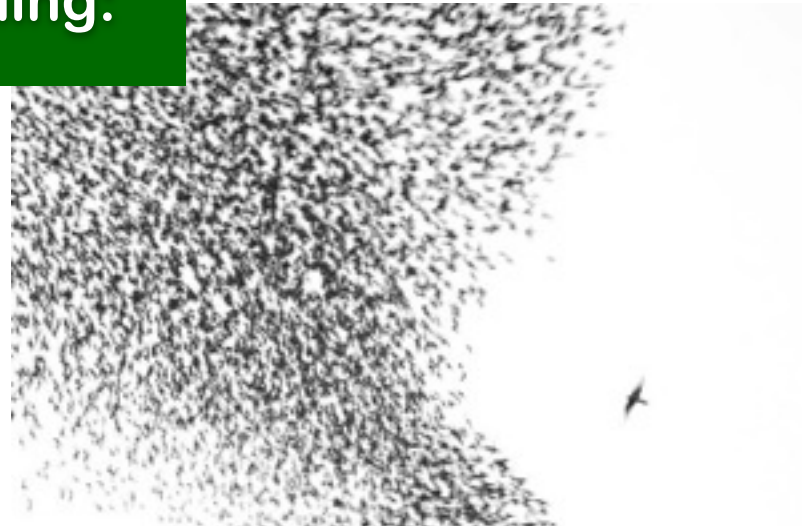


Falcon attacks a flock of starlings

How do we describe systems with large N of particles?



Swarming:



How do we describe systems with large N of particles?



Traffic flow:

How do we describe systems with large N of particles?



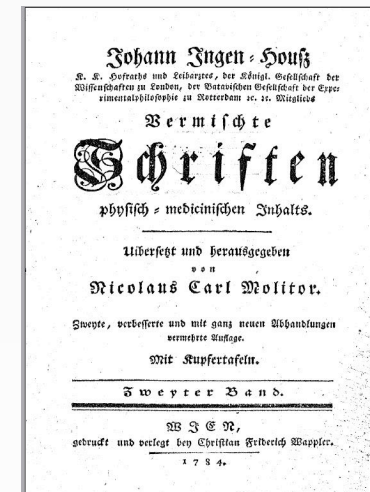
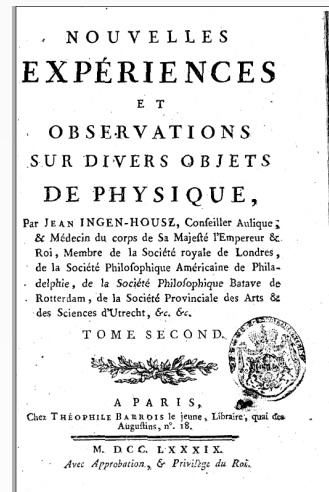
1st example:
Brownian motion

Brownian motion



“Brownian” motion

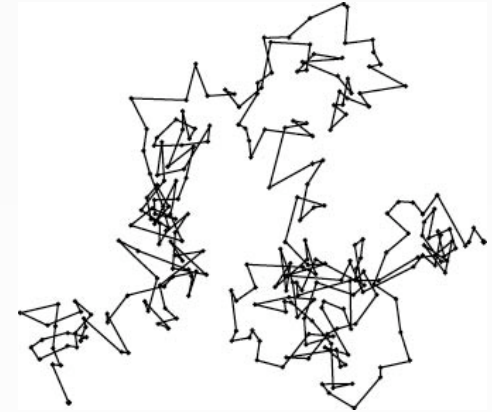
Jan Ingen-Housz (1730-1799)



1784/1785:

über betrügen könnte, darf man nur in den Brennpunct eines Mikroskops einen Tropfen Weingelst sammt etwas gestoßener Kohle setzen; man wird diese Körperchen in einer verwirrten beständigen und heftigen Bewegung erblicken, als wenn es Thierchen wären, die sich reißend unter einander fortbewegen.

Robert Brown (1773-1858)

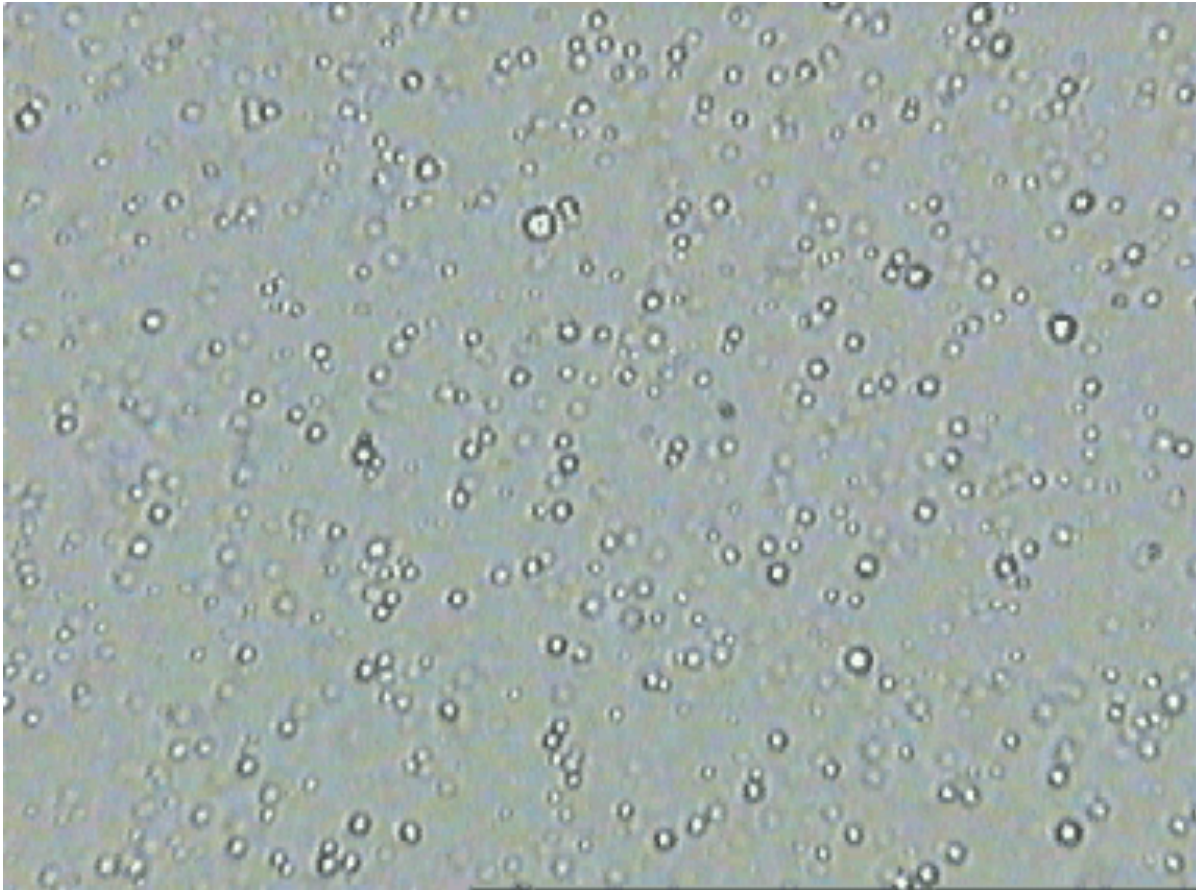


Linnean Society (London)

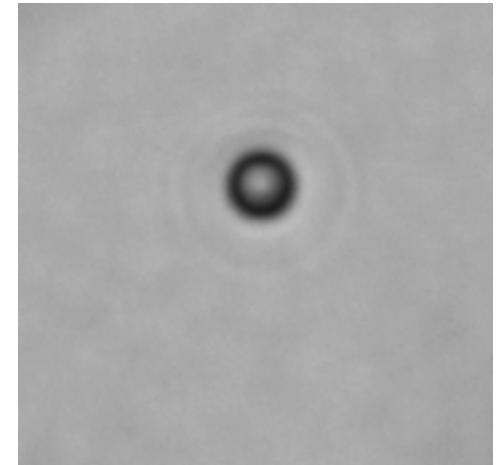
1827: irregular motion of pollen in fluid

<http://www.brianjford.com/wbbrownc.htm>

Brownian motion



David Walker

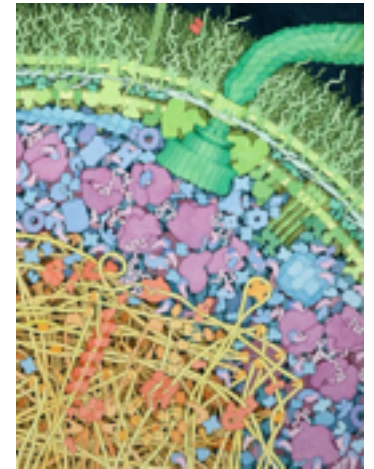


Mark Haw

Polymer in a fluid

Dogic lab
(Brandeis)

18.S995

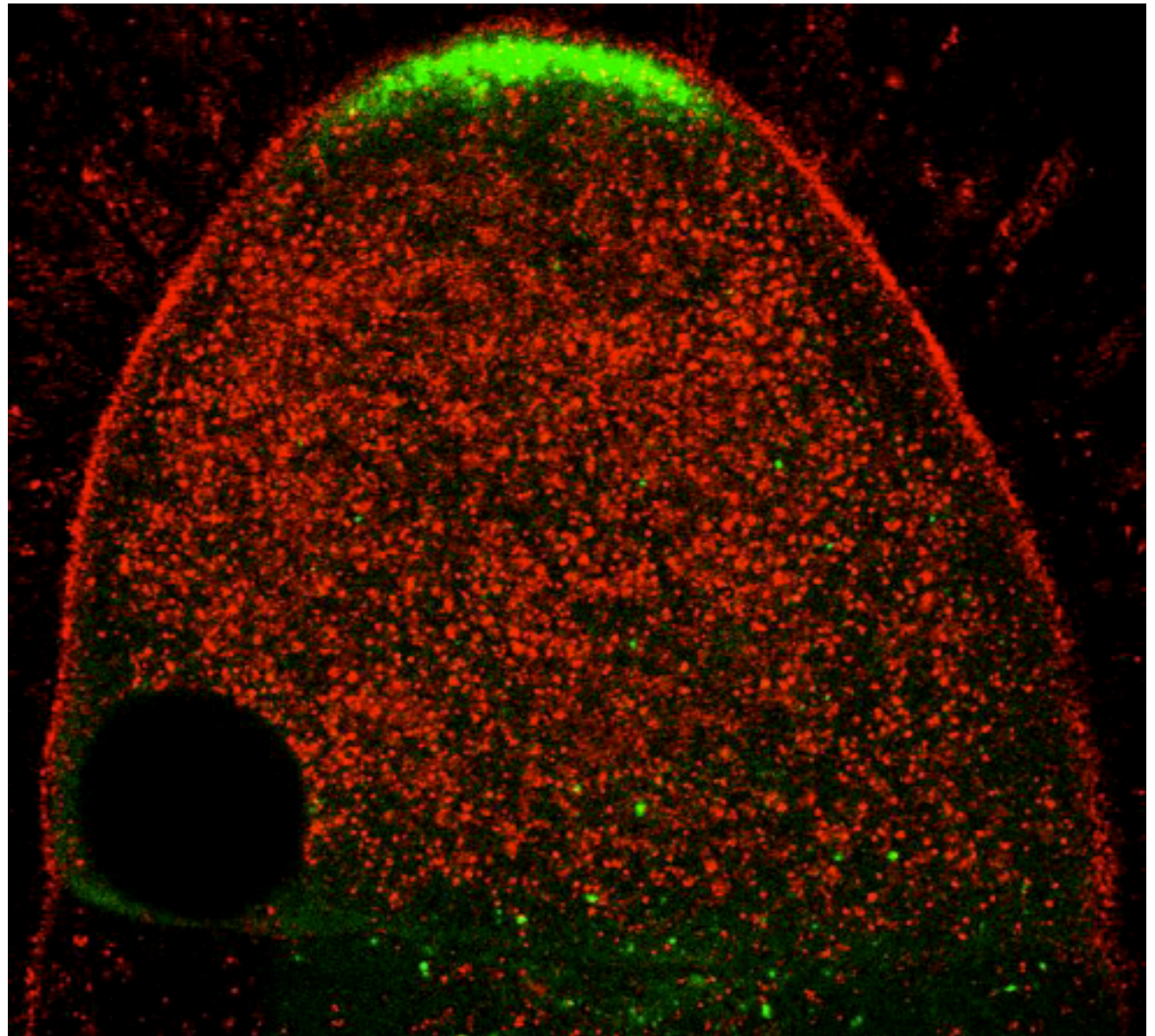


$< 1\mu m$

Flow & transport in cells



Drosophila
embryo



Goldstein lab (Cambridge)

Mathematical description of standard BM

Adolf Eugen Fick (1829-1901)

first law

$$\frac{\partial}{\partial t} \varrho(t, x) = -\nabla J(t, x)$$

second law

$$J(t, x) = -\mathcal{D} \nabla \varrho(t, x)$$

$$\Rightarrow \frac{\partial}{\partial t} \varrho = \mathcal{D} \nabla^2 \varrho$$



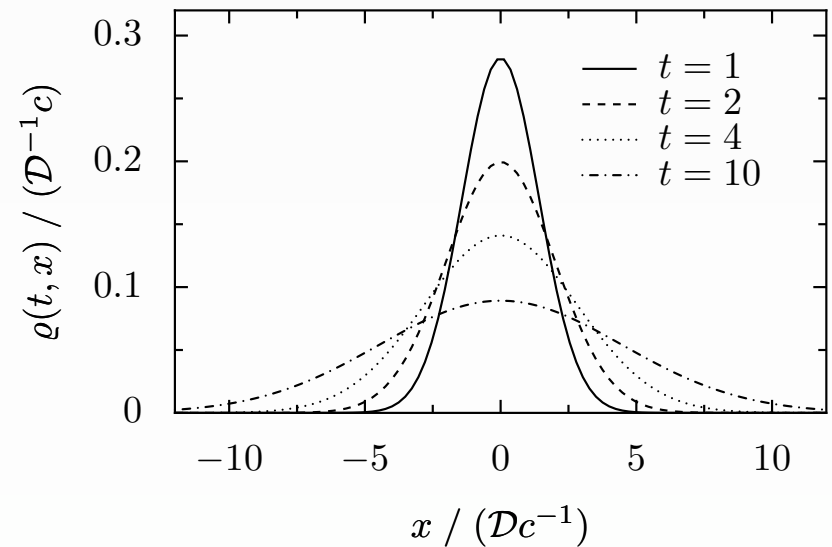
Phil. Mag. 10: 30 (1855)



Louis Bachelier, *Theorie de la speculation*,
Ann. Sci. l'École Norm. Sup. 3 (17): 21–86 (1900)

$$\langle x^2(t) \rangle = 2Dt$$

Goal: $D=f(\text{microscopic parameters})$



1870 – 1946



Louis Bachelier, aged 15

Annales scientifiques de l'É.N.S. 3^e série, tome 17 (1900), p. 21-86.

THÉORIE
DE
LA SPÉCULATION,

PAR M. L. BACHELIER.

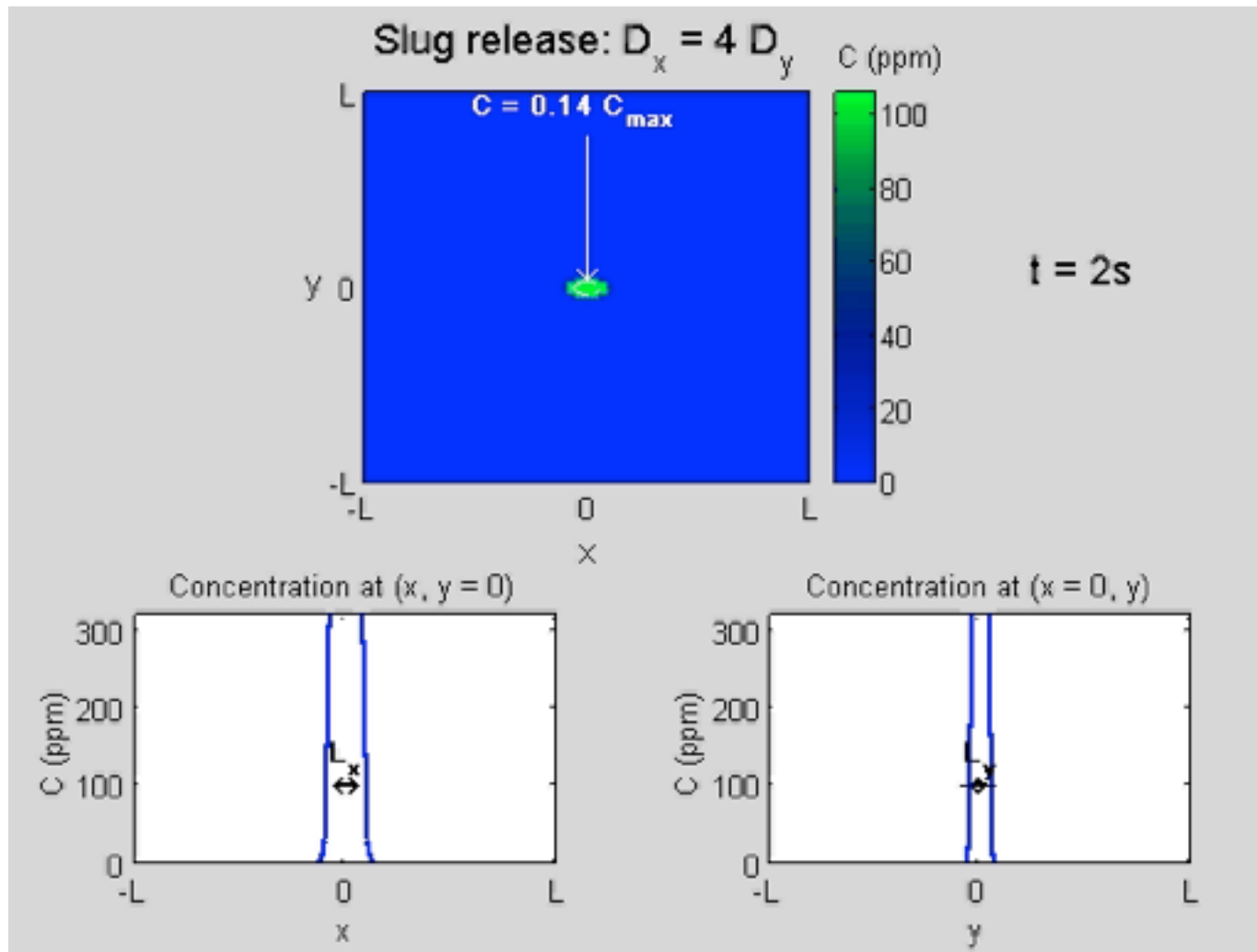
38

L. BACHELIER.

Nous avons donc pour expression de la probabilité

$$p = \frac{H}{\sqrt{t}} e^{-\frac{\pi H^2 x^2}{t}}.$$

Diffusion of an Instantaneous Point Source



This animation depicts the diffusion of a discrete mass released at $(x = 0, y = 0, t = 0)$. The diffusion is anisotropic, $D_x = 4 D_y$. The length scales grow in proportion to the square root of the diffusion, such that the dimensions of the cloud are anisotropic, with $L_x = 2 L_y$. Note that the profiles of concentration along the x - and y -axes are Gaussian in shape.

W. Sutherland (1858-1911)



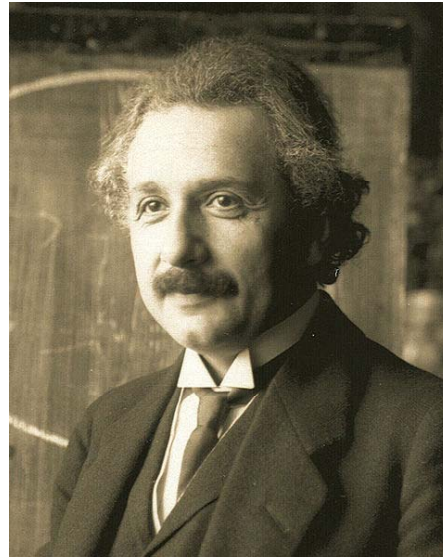
Source: www.theage.com.au

$$D = \frac{RT}{6\pi\eta aC}$$

Phil. Mag. **9**, 781 (1905)

R : molar gas constant
 k : viscosity

A. Einstein (1879-1955)



Source: wikipedia.org

$$\langle x^2(t) \rangle = 2Dt$$
$$D = \frac{RT}{N} \frac{1}{6\pi kP}$$

Ann. Phys. **17**, 549 (1905)

T : Temperature
 P : ParticleRadius

M. Smoluchowski
(1872-1917)



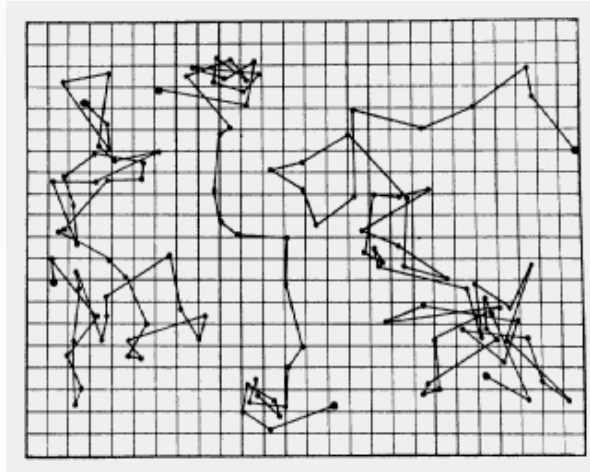
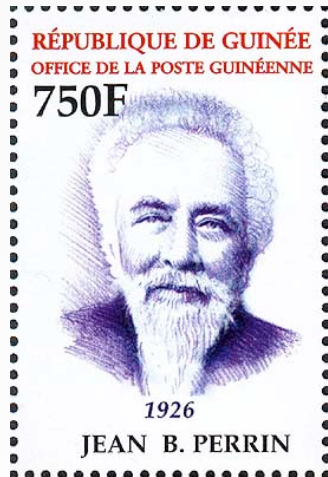
Source: wikipedia.org

$$D = \frac{32}{243} \frac{mc^2}{\pi\mu R}$$

Ann. Phys. **21**, 756 (1906)

N : Avogadro's number

Jean Baptiste Perrin (1870-1942, Nobel prize 1926)



- ▶ colloidal particles of radius $0.53\mu\text{m}$
- ▶ successive positions every 30 seconds joined by straight line segments
- ▶ mesh size is $3.2\mu\text{m}$

Mouvement brownien et réalité moléculaire, Annales de chimie et de physique VIII 18, 5-114 (1909)

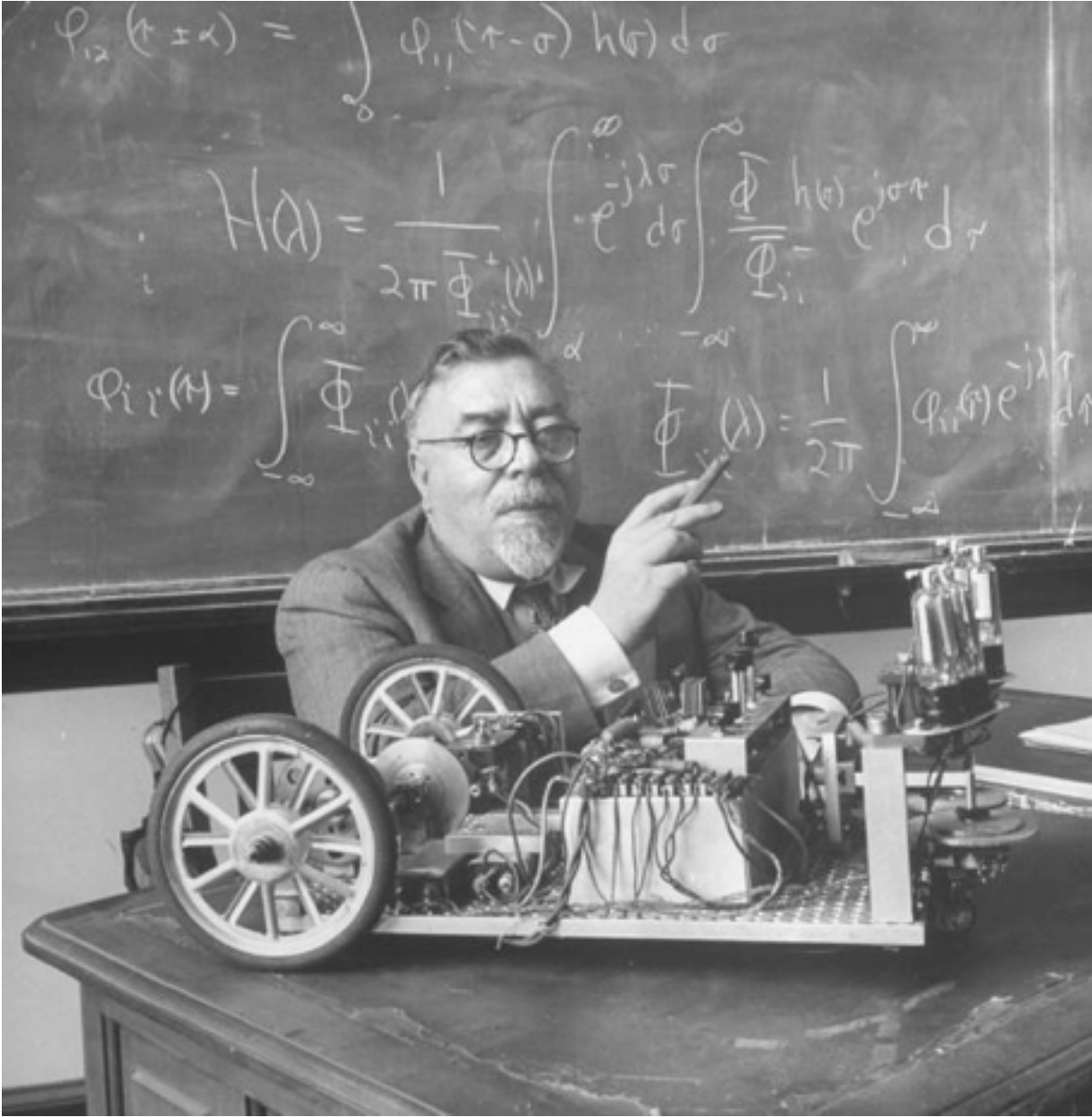
Les Atomes, Paris, Alcan (1913)

experimental evidence for
atomistic structure of matter

Norbert Wiener

(1894-1964)

MIT



How fast must a cell swim to beat Brownian motion?

$$\langle x^2 \rangle = 2Dt$$

$$D = \frac{kT}{6\pi\eta_0 a}$$

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Hence, we find for the diffusion constant

$$D \sim 0.2 \mu\text{m}^2/\text{s}$$

Assuming a run length ~ 1 s, Brownian motion would move a micron-sized bacterium by approximately $0.5 \mu\text{m}$ per second. Thus a bacterium should swim at last $5\text{-}10 \mu\text{m/s}$, which is close to typical bacterial swimming speeds.

Sedimentation

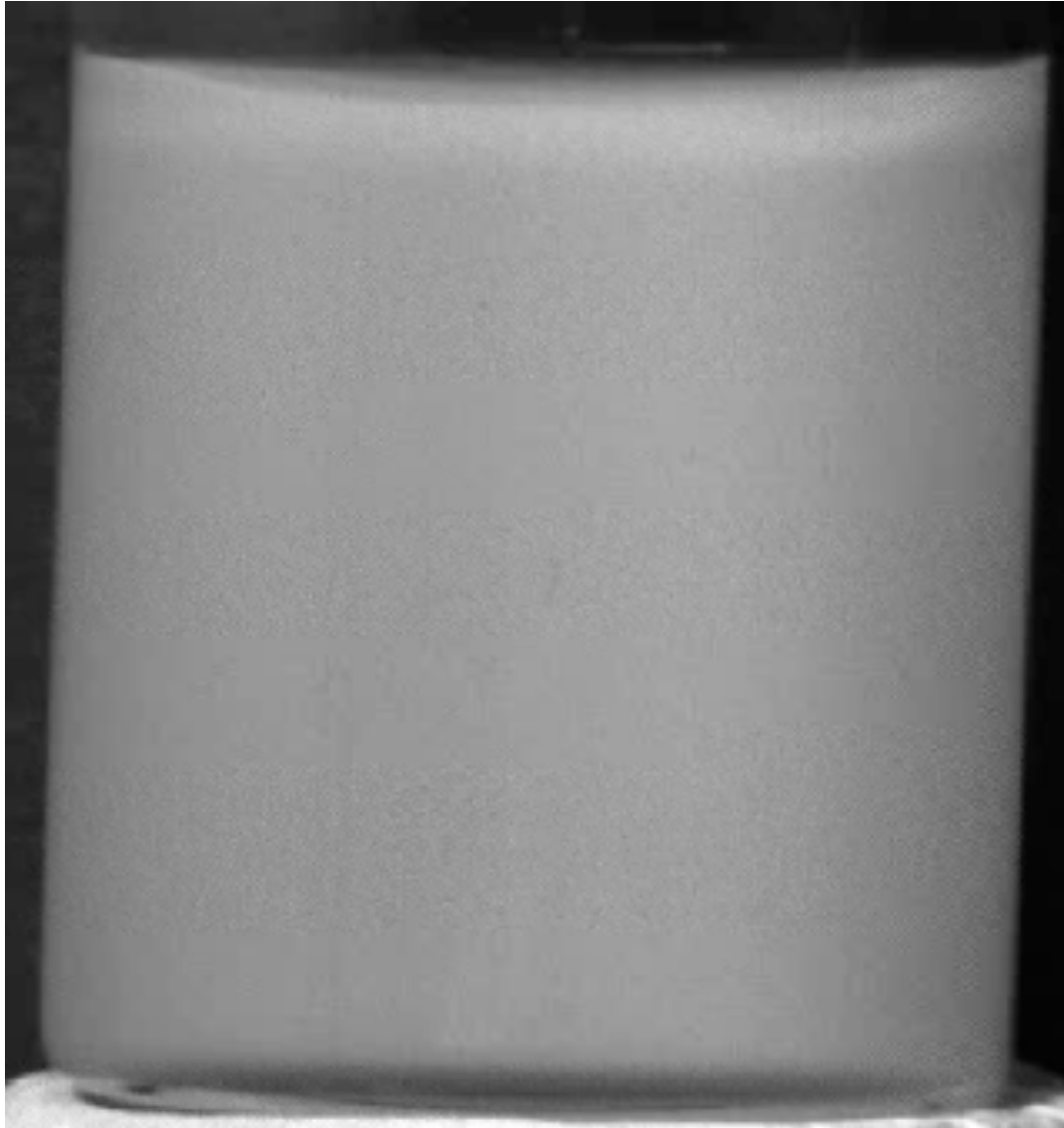


Mississippi

NASA Earth Observatory



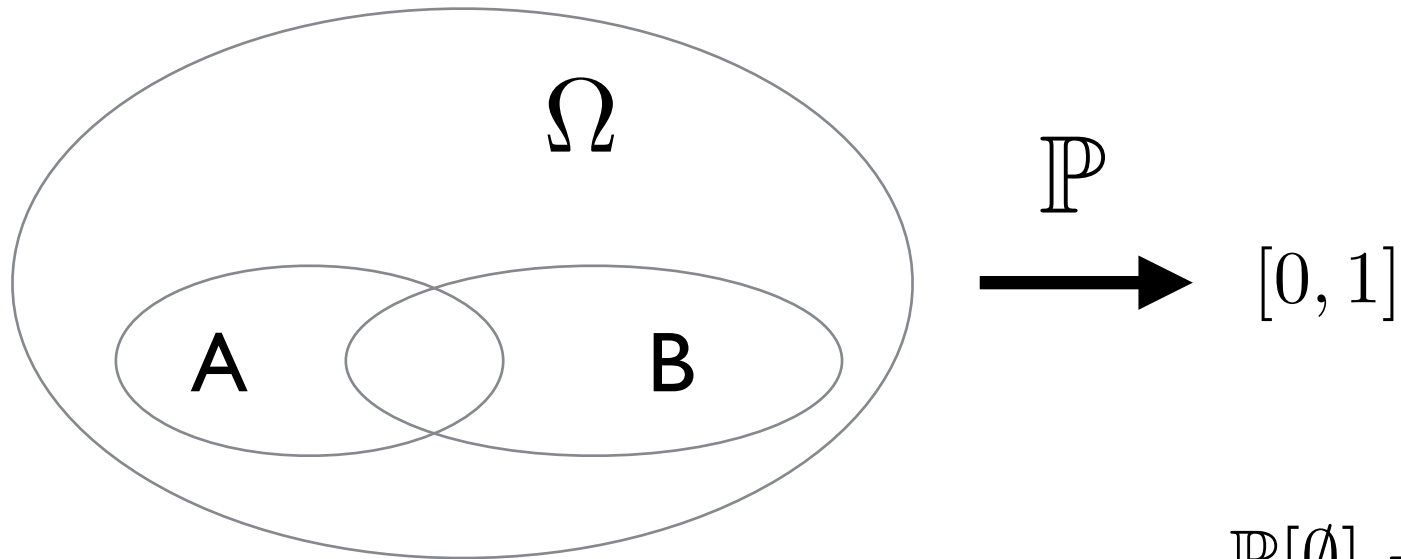
Sedimentation



Particle
separation

Probability space $(\Omega, \mathcal{F}, \mathbb{P})$

$$\mathcal{F} = \{\emptyset, A, B, A \cap B, A \cup B, \dots, \Omega\}$$



$$\mathbb{P}[\emptyset] = 0$$

$$\mathbb{P}[\Omega] = 1$$

$$\mathbb{P}[A \cup B] = \mathbb{P}[A] + \mathbb{P}[B] - \mathbb{P}[A \cap B]$$

Expectation values of random variables

$$X : \Omega \rightarrow \mathbb{R}^n$$

$$\mathbb{E}[f(X)] = \int d\mathbb{P} f(x) = \int dx p(x) f(x)$$

$$p(x) \geq 0, \quad \int dx p(x) = 1$$

$$\mathbb{E}[\alpha f(X) + \beta g(X)] = \alpha \mathbb{E}[f(X)] + \beta \mathbb{E}[g(X)]$$

This course simply:

$$p(x) \geq 0, \quad \int dx p(x) = 1$$

$$\mathbb{E}[f(X)] = \int d\mathbb{P} f(x) = \int dx p(x) f(x)$$

$$\mathbb{E}[\alpha f(X) + \beta g(X)] = \alpha \mathbb{E}[f(X)] + \beta \mathbb{E}[g(X)]$$