

• Radiation field (far from antenna near the origin)

$$u_{inc}(x, \omega) \approx \int \frac{e^{i\omega \|x\|}}{4\pi \|x\|} e^{-i\omega \hat{x} \cdot y} j(y, \omega) dy$$

$$= \frac{e^{i\omega \|x\|}}{4\pi \|x\|} \hat{j}(\omega \hat{x}, \omega)$$

current density (vector) $J(y) f(\omega)$
 \times pulse f.
 called radiation vector, $= J(\hat{x}, \omega)$
 or its scalar analogue.

when antenna at x_0 :

$$u_{inc}(x, \omega) = \frac{e^{i\omega \|x-x_0\|}}{4\pi \|x-x_0\|} J(\hat{x}-\hat{x}_0, \omega)$$

radiation beam pattern & transmitted waveform

• Reception: $S_{rec}(\omega) = \int_{\text{antenna near } x_0} u_{RE}(y, \omega) \hat{w}(y, \omega) dy$
 (supp V far from antenna x_0)
 \hat{w} window (rec \rightarrow real)

$$= \int \int \frac{e^{i\omega \|y-y_0\|}}{4\pi \|y-y_0\|} (-\omega^2) V(y) u_{inc}(y, \omega) \hat{w}(y, \omega) dy dz$$

$$\approx \int \frac{e^{i\omega \|x_0-y\|}}{4\pi \|x_0-y\|} (-\omega^2) V(y) u_{inc}(y, \omega) \hat{w}(\omega(y-\hat{x}_0), \omega) dy$$


$$V(\hat{x}-\hat{x}_0, \omega)$$

reception beam pattern

(usually, same as radiation beam pattern, disregarding pulse ($f(\omega)=1$))

$$\approx \int e^{2i\omega \|x_0-y\|} A(\omega, x_0, y) V(y) dy$$

where $A(\omega, x_0, y) = -\omega^2 \frac{J(\omega, \hat{y}-\hat{x}_0) V(\omega, \hat{y}-\hat{x}_0)}{(4\pi \|x_0-y\|)^2}$


 speed targets < aircraft < speed of light
 (Vigdet) (x. rad = x. sec.)

Continuum model for slow time, $x = x(s)$

$$d(s, \omega) = \int_{\text{rec B}} (x(s), \omega)$$

$$d(s, \omega) = \int e^{2i\omega \|x(s) - y\|} A(\omega, s, y) V(y) dy = FV$$

$$\Rightarrow d(t, s) = \int e^{-i\omega [t - 2\|x(s) - y\|/c]} A(\omega, s, y) V(y) dy d\omega$$

$$= \int \hat{A}(t - 2\frac{\|x(s) - y\|}{c}, s, y) V(y) dy$$

Formal adjoint:

$$F^* d = \int e^{-2i\omega \|x(s) - x\|} \overline{A(\omega, s, x)} d(s, \omega) ds$$

\rightarrow Backprojection.

10/20

Compl.: $F: X \rightarrow \mathbb{R}$ functional.

$f \rightarrow F(f)$

Begin appendix

F is Fréchet-differentiable at $f_0 \in X$ if

$$\lim_{h \rightarrow 0} \frac{\|F(f_0 + h) - F(f_0) - A_{f_0}(h)\|}{\|h\|} = 0$$

for some bounded lin. op. A_{f_0} :
(functional)

Call $A_{f_0} = \frac{\delta F}{\delta f}[f_0]$ the derivative of F at f_0 .

also called 1st variation of F at f_0 .

e.g. if $F: \mathbb{R}^m \rightarrow \mathbb{R}$,

$$\frac{\delta F}{\delta f}[f_0](h) = (\nabla F)(f_0) \cdot h$$

\hookrightarrow inner product

Formula for $\frac{\delta F}{\delta f}$. Gateaux derivative

$$\frac{\delta F}{\delta f}[f_0](h) = \lim_{t \rightarrow 0} \frac{F(f_0 + th) - F(f_0)}{t}$$

(directional derivative in the direction h).

Ex. $F(f) = \langle f, g \rangle$

Then $\frac{\delta F}{\delta f}[f_0] = g$ independent of f_0

$$\frac{\delta F}{\delta f}[f_0](h) = \langle h, g \rangle$$

Log. mult. $\left(\min_x f(x) \text{ s.t. } g(x) = c \right)$

$\Rightarrow d(x, \lambda) = f(x) - \lambda \cdot (g(x) - c)$

then $\nabla_x d = 0$
 $\nabla_\lambda d = 0$

Ex. of functionals and their derivatives

$$F(f) = \langle f, g \rangle \Rightarrow \frac{\delta F}{\delta f} = g(x)$$

$$F(f) = f(0) \Rightarrow \frac{\delta F}{\delta f} = \delta(x)$$

$$F(f) = f'(0) \Rightarrow \frac{\delta F}{\delta f} = -\delta'(x)$$

Nonlinear functional: $F(f) = \langle f^2, g \rangle$
 $\Rightarrow \frac{\delta F}{\delta f} = 2fg$

Remark: any bounded linear functional can be represented as the inner product with a distribution (Riesz representⁿ thm)

end
appendix

19/22/09

Intro.
(end of Day 3)

Time formul.

Linearized field $u_{x,B}(\alpha, t)$

Compare to (processed) data

$$d(\alpha, t) = rec(\alpha, t) - u_{inc}(\alpha, t)$$

↳ actual recorded

Linearized inverse problem:

$$J[u_{x,B}, V] = \frac{1}{2} \|u_{x,B} - d\|_2^2$$

$$s.t. \quad u_{x,B} = FV.$$

↳ linearized forward map

a) least-squares:

$$\Rightarrow -\frac{\delta J}{\delta V} = F^*(d - u_{x,B})$$

b) Variations of $J = 0$:

$$\left(\frac{\delta J}{\delta V} = 0\right) \rightarrow -\frac{\delta J}{\delta V} = \int_0^T q(\alpha, t) \frac{\partial^2 u_{inc}(\alpha, t)}{\partial t^2} dt \quad (RTM)$$

$$\left(\frac{\delta J}{\delta u_{x,B}(\alpha, t)} = 0\right) \text{ with } \left(\frac{1}{c_0^2} \frac{\partial^2}{\partial t^2} - \Delta\right) q = d_{ext}(\alpha, t) - u_{x,B,ext}(\alpha, t)$$

$$\left(\frac{\delta J}{\delta u_{x,B}(\alpha, T)} = 0\right) \rightarrow \frac{\partial q}{\partial t}(\alpha, T) = 0$$

$$\left(\frac{\delta J}{\delta \frac{\partial u_{x,B}(\alpha, T)}{\partial t}} = 0\right) \rightarrow q(\alpha, T) = 0$$

$q =$ Lagrange multiplier
 $=$ adjoint state.

Comp. (a) & (b):

$$F^*(d - u_{x,B}) = \int_0^T q(\alpha, t) \frac{\partial^2 u_{inc}(\alpha, t)}{\partial t^2} dt \quad (RTM)$$

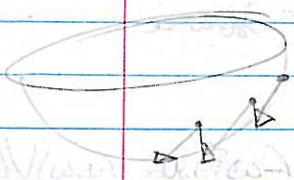
$$\text{Imaging } F = \tilde{V}(x) \text{ approx. to } V(x)$$

Go further? $d = rec - u_{inc}$

a) Fix c_0 , obtain the min of J :

$$F^{-1} = (F^* F)^{-1} F^*$$

Could be obtained by iteration:



① $d \rightarrow V_0 = \alpha_0 F^* d$
↳ small number

② $V_0 \rightarrow u_{n,B,0} = FV_0$

$$\rightarrow V_1 = V_0 + \alpha_1 F^*(d - u_{n,B,0})$$

③ $V_1 \rightarrow u_{n,B,1} = FV_1$

$$\rightarrow V_2 = V_1 + \alpha_2 F^*(d - u_{n,B,1})$$

or linear solver (GMRES) or...

b) Update $c(x)$ at each step: $d_0 = rec - u_{inc,0}$

① $d_0 \rightarrow V_0 = \alpha_0 F^* d_0$

② $V_0 \rightarrow \frac{1}{c_1^2} = \frac{1}{c_0^2} - V_0$

$\rightarrow u_{inc,1}$ simulated in c_1^2

$$\rightarrow d_1 = rec - u_{inc,1}$$

$$\rightarrow V_1 = \alpha_1 F^* d_1$$

③ $V_1 \rightarrow \frac{1}{c_2^2} = \frac{1}{c_1^2} - V_1$ etc.

Called full waveform inversion

* Here: talked about Tikhonov regularization.

(Chap 3/4) 2) Freq. formulation.

10/27/09 Kinematics of F, F^* in radar imaging:
Backprojection (Stripmap SAR)

• Radiation field

$$u_{inc}(x, \omega) = \int \frac{e^{i\omega \|x-y\|}}{4\pi \|x-y\|} j(y, \omega) dy$$

source

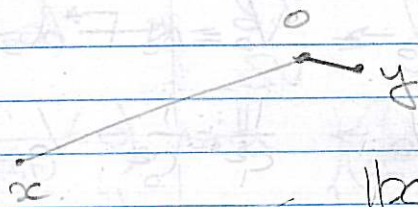
• Scattered field

$$u_{sc,B}(x, \omega) = \int \frac{e^{i\omega \|x-y\|}}{4\pi \|x-y\|} (-\omega^2) u_{inc}(y, \omega) V(y) dy$$

• Far field expansion:

Rad field: antenna at $y=0$.

$$\Rightarrow \|x\| \gg \|y\|$$



$$\|x-y\| = \dots = \|x\| - \hat{x} \cdot y + O\left(\frac{\|y\|^2}{\|x\|}\right)$$

geometry:
spherical
wave fronts
look planar
from a distance

$$e^{i\omega \|x-y\|} = e^{i\omega \|x\|} e^{-i\omega \hat{x} \cdot y}$$

x remainder

$$\text{with } |\text{remainder}| \leq \left| e^{i\omega O\left(\frac{\|y\|^2}{\|x\|}\right)} \right|$$

$$\leq C \frac{\omega \|y\|^2}{\|x\|}$$

$$= O\left(\frac{\omega \|y\|^2}{\|x\|}\right)$$

$$\frac{1}{\|x-y\|} = \frac{1}{\|x\|} \left(\frac{1}{1 - \frac{\hat{x} \cdot y}{\|x\|} + O\left(\frac{\|y\|^2}{\|x\|^2}\right)} \right)$$

$$= \frac{1}{\|x\|} \left(1 + O\left(\frac{\|y\|}{\|x\|}\right) \right)$$

$$\Rightarrow G(r, y, \omega) \approx \frac{e^{i\omega \|x\|}}{4\pi \|x\|} e^{-i\omega \hat{x} \cdot y}$$

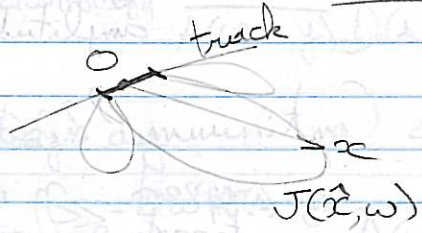
when $\left. \begin{array}{l} \|y\| \ll \|x\| \\ \omega \|y\|^2 \ll \|x\| \end{array} \right\}$

• Radiation field far from antenna:

$$\begin{aligned} u_{inc}(r, \omega) &\approx \int_{\text{antenna}} \frac{e^{i\omega \|x\|}}{4\pi \|x\|} e^{-i\omega \hat{x} \cdot y} j(y, \omega) dy \\ &= \frac{e^{i\omega \|x\|}}{4\pi \|x\|} \hat{j}'(\omega \hat{x}, \omega) \end{aligned}$$

$j(y, \omega) = c d n s(y) P(\omega)$
 ↳ "current density vector" (antenna shape) (scalar analogue) → pulse (signal sent) (in freq.)

$\hat{j}'(\omega \hat{x}, \omega) = J(\hat{x}, \omega)$
 ↳ "radiation vector" (scalar analogue) gives the radiation beam pattern.



Change coord: antenna at x_0

$$\Rightarrow u_{inc}(r, \omega) \approx \frac{e^{i\omega \|x-x_0\|}}{4\pi \|x-x_0\|} J(\hat{x}-\hat{x}_0, \omega)$$

Scattered field / reception

Stripmap SAR:
antenna fixed
along track.

$$S_{rec, B}(\omega) = \int_{\text{antenna near } x_0} u_{rec, B}(\vec{y}, \omega) w(\vec{y}, \omega) dy$$

window
(real)



far field

$$= \int_{\text{supp } V} \frac{e^{i\omega \|y - y_0\|}}{4\pi \|y - y_0\|} (-\omega^2) u_{inc}(y, \omega) V(y) w(\vec{y}, \omega) dy$$

$$\approx \int_{\text{supp } V} \frac{e^{i\omega \|y - x_0\|}}{4\pi \|y - x_0\|} (-\omega^2) u_{inc}(y, \omega) V(y) \hat{w}'(\omega(\hat{y} - \hat{x}_0), \omega) dy$$

Call $\hat{w}'(\omega(\hat{y} - \hat{x}_0), \omega) = W(\hat{y} - \hat{x}_0, \omega)$
the reception beam pattern

Plug

u_{inc}
(same x_0
for rad/rec)

Same as J, usually, up to phase =

$$J(\hat{y} - \hat{x}_0, \omega) = W(\hat{y} - \hat{x}_0, \omega) P(\omega)$$

$$S_{rec, B}(\omega) \approx \int_{\text{supp } V} e^{2i\omega \|y - x_0\|} A(y, x_0, \omega) V(y) dy$$

with $A(y, x_0, \omega) = -\omega^2 \frac{J(\hat{y} - \hat{x}_0, \omega) W(\hat{y} - \hat{x}_0, \omega)}{4\pi \|x_0 - y\|^2}$

↓
amplitude

△ Continuum model for slow time: $x_0 = x(t)$

speed targets \ll speed aircraft \ll speed light
($V \pm t$) ($x_0, rec = x_0, rad$)

$$D(s, \omega) = S_{rec, B}(x(s), \omega)$$

in this approx,

light
rotational
approx.

$$D(s, \omega) = \int e^{2i\omega \|x(s) - y\|} A(\omega, s, y) V(y) dy$$

$$= (FV)(s, \omega)$$

Back to t:

$$D(s, \omega) = d(t, s) = \iint e^{-i\omega [t - 2\|x(s) - y\|]} A(\omega, s, y) V(y) dy d\omega$$

$$= \int \hat{A}\left(t - \frac{2\|x(s) - y\|}{c}, s, y\right) V(y) dy$$

$c \rightarrow \omega c = 1$

$$\approx \int \delta\left(t - \frac{2\|x(s) - y\|}{c}\right) \cdot \frac{V(y)}{4\pi \|x(s) - y\|^2} dy$$

△ $\int_{\omega} \frac{1}{\omega} d\omega$
 $\int_{\omega} \frac{1}{\omega} d\omega$

located
near $t=0$

Geometry: time for signal to come back:
 $t = \frac{2\|x(s) - y\|}{c}$

Adjoint of F: F^* (imaging of $\omega \rightarrow 0$)

$$\langle D, FV \rangle_{(s, \omega)} = \langle F^*D, V \rangle_y$$

$$\iint D(s, \omega) \overline{\int e^{2i\omega \|x(s) - y\|} A(\omega, s, y) V(y) dy} ds d\omega$$

$$\int V(y) \iint e^{-2i\omega \|x(s) - y\|} \overline{A(\omega, s, y)} D(s, \omega) ds d\omega$$

$$\int V(y) F^*D(y) dy$$

$$\Rightarrow F^*D(y) = \int e^{-2i\omega \|x(s) - y\|} \overline{A(\omega, s, y)} D(s, \omega) ds d\omega$$

$$= \int e^{-2i\omega \|x(s) - y\|} \overline{A(\omega, s, y)} \int e^{+i\omega t} d(s, t) dt ds d\omega$$

$$= \int \hat{A}\left(\frac{2\|x(s) - y\|}{c} - t, s, y\right) d(s, t) ds dt$$

high
low
circles

Because

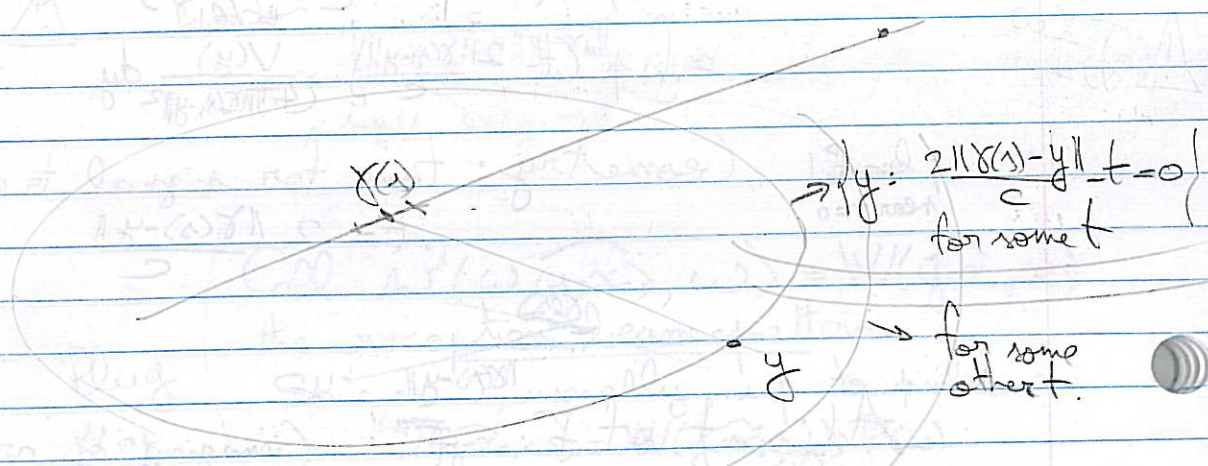
$$f(\omega) \rightarrow \hat{f}(t) = \int e^{-i\omega t} f(\omega) d\omega$$

$$\bar{f}(\omega) \rightarrow g(t) = \int e^{-i\omega t} \bar{f}(\omega) d\omega$$

$$= \int e^{i\omega t} f(\omega) d\omega$$

$$= \hat{f}(-t)$$

$$\approx \iint f''\left(\frac{2\|x(s)-y\|}{c} - t\right) \frac{d(s,t)}{(\|x(s)-y\|^2)} ds dt$$



Consider circles in y weighted by data $d(s,t)$,
 corresp. to geom. laws $\frac{2\|x(s)-y\|}{c} = t$,
 then add up.

\Rightarrow backprojection

Similar to inverting a Radon trf.

Rmk: possible project.

$y_T = \begin{pmatrix} y \\ s \end{pmatrix}$ | Recall $D(s, \omega) = FV(s, \omega) = \int e^{2i\omega \|x(s) - y_T\|} A(\omega, s, y_T) V(y) dy$ ① $y \in \mathbb{R}^2$

10/29/09

Imaging operator w. I.O.:

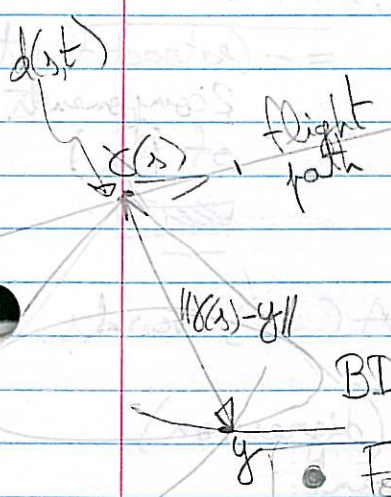
Find better amplitude $Q(\omega, s, y_T)$ in

$$F^*D \rightarrow BD(y) = \int e^{-2i\omega \|x(s) - y_T\|} Q(\omega, s, y_T) D(s, \omega) ds d\omega$$

and again, $y_T = \begin{pmatrix} y \\ 0 \end{pmatrix}$

$y = \frac{1}{2} y_T$ (extract first 2 compon.)

- If $Q=1$, then pure backprojection. because (idealize w. \mathbb{R}).



$$\begin{aligned} & \int e^{-2i\omega \|x(s) - y_T\|} D(s, \omega) ds d\omega \\ &= \iint e^{-2i\omega \|x(s) - y_T\|} \int e^{i\omega t} d(s, t) dt ds d\omega \\ &= \iint \delta\left(t - \frac{2 \|x(s) - y_T\|}{c}\right) d(s, t) ds dt \end{aligned}$$

$$BD(y) = \int d\left(s, \frac{2 \|x(s) - y_T\|}{c}\right) ds$$

- For inversion,

$$BD(y) = \iint_{s, \omega} e^{-2i\omega \|x(s) - y_T\| + 2i\omega \|x(s) - x\|} Q(\omega, s, y_T) A(\omega, s, x) ds d\omega \times V(x) dx$$

would like $\approx V(y)$

\rightarrow choose Q s.t. $K(y, x) \approx \delta(y - x)$ $x, y \in \mathbb{R}^2$

by reducing $K(y, x)$ to $\int_{\mathbb{R}^2} e^{i(y-x) \cdot \xi} d\xi$ for some ξ .

Here $\iint ds d\omega$ over a data collection manifold M
 \rightarrow idealize over \mathbb{R}^2 . \triangle

Call $R_{y,s} = \delta(s) - y_T$

For ref, $\frac{\partial R_{y,s}}{\partial s} = \delta'(s)$ $\begin{matrix} 1 \\ 3 \end{matrix}$ $\frac{\partial R_{y,s}}{\partial y} = - \begin{pmatrix} 1 & 0 \\ 0 & 1 \\ 0 & 0 \end{pmatrix}$ $\begin{matrix} 1 \\ 3 \end{matrix}$

$= -P_2$

$\frac{\partial \|R_{y,s}\|}{\partial s} = \frac{x}{\|R_{y,s}\|}$

$\square \frac{\partial \|R_{y,s}\|}{\partial s} = \frac{R_{y,s}^T}{\|R_{y,s}\|} \frac{\partial R_{y,s}}{\partial s} = \hat{R}_{y,s} \cdot \delta'(s)$

\downarrow
unit vector

$\frac{\partial \|R_{y,s}\|}{\partial y} = \frac{R_{y,s}}{\|R_{y,s}\|} \frac{\partial R_{y,s}}{\partial y} = - \hat{R}_{y,s} \cdot P_2$

$= -$ (extract first 2 components of $\hat{R}_{y,s}$)

$K(y,x) = \int e^{2i\omega(\|R_{y,s}\| + \|R_{x,s}\|)} Q A(\dots) ds ds$

Stationary phase: (digression)

a) $\phi(x)$ has no crit. point:
 $\nabla \phi(x^*) \neq 0$

$a(x)$ smooth, compactly supported,

$\lim_{\beta \rightarrow \infty} \int_{\mathbb{R}^n} e^{i\beta\phi(x)} a(x) dx = I(\beta)$, $\beta \rightarrow \infty$

then $|I(\beta)| \leq C \cdot \beta^{-N}$ as $\beta \rightarrow \infty$
 $\forall N \geq 0$.

Pf. $\alpha = \frac{1}{i\beta\phi'(x)} \frac{d}{dx}$

$\alpha \int e^{i\beta\phi(x)} = e^{i\beta\phi(x)}$

$|\int_{\mathbb{R}^n} \alpha a(x)| \leq \beta^{-N}$ \square

b) $\phi(x)$ has a crit. point: x^*

$$\nabla \phi(x^*) = 0$$

$$D^2 \phi(x^*) \neq 0.$$

$a(x)$ smooth, comp. supp.

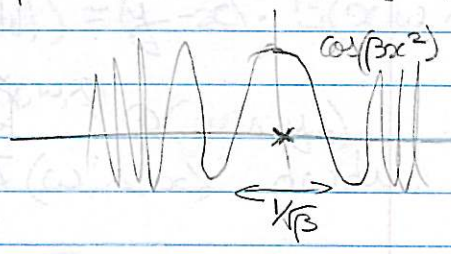
Thm. $\int_{\mathbb{R}^n} e^{i\beta \phi(x)} a(x) dx$

$$= \left(\frac{2\pi}{\beta}\right)^{n/2} a(x^*) e^{i\beta \phi(x^*)} \frac{e^{i\frac{\pi}{4} \text{sgn } D^2 \phi(x^*)}}{\sqrt{|\det D^2 \phi(x^*)|}}$$

$$+ O(\beta^{-n/2-1}) \text{ as } \beta \rightarrow \infty.$$

sgn $M = \# \text{ pos eig } \nu(M) - \# \text{ neg eig } \nu(M)$

Ex. $\int_{\mathbb{R}} e^{i\beta x^2} dx$



$$\sim \frac{1}{\sqrt{\beta}}$$

Extension: $\left| \frac{\partial^k \phi}{\partial x^k}(x^*) \right| = 0 \quad k=1, \dots, l$
 $\frac{\partial^l \phi}{\partial x^l}(x^*) \neq 0.$
 $\rightarrow O(\beta^{-l/2})$

Ex. $\iint e^{i(y-x) \cdot \xi} d\xi f(x) dx \stackrel{?}{=} f(y)$ X
 $\iint e^{i(y-x) \cdot \xi} d\xi \left(\int e^{ix \cdot \eta} F(\eta) d\eta \right) dx$

\rightarrow stat. phase in pair (x, ξ) of variables!

$$\phi(x, \xi) = (y-x) \cdot \xi + x \cdot \eta$$

y, η fixed param

$$\frac{\partial \phi}{\partial x} = \eta - \xi = 0 \Rightarrow \xi = \eta$$

$$\frac{\partial \phi}{\partial \xi} = y - x = 0 \Rightarrow x = y$$

canonical relation.

$$\frac{\partial \phi}{\partial \xi} \frac{\partial \xi}{\partial x} = -Id \text{ non-sing}$$

$$\phi(y, \eta) = (y-y) \cdot \eta + y \cdot \eta = y \cdot \eta$$

$$\rightarrow \text{expect } \sim \int e^{iy \cdot \eta} F(\eta) d\eta = f(y)$$

Back to $K(y, x)$: $e^{2i\omega(|R_{ys}| - |R_{xs}|)}$

Large parameter: $\omega = \beta\omega'$
 $d\omega = \beta d\omega'$

Phase: $\Phi(\omega', s) = 2\omega'(|R_{ys}| - |R_{xs}|)$ x, y fixed.

$$\frac{\partial \Phi}{\partial \omega'} = 2(|R_{ys}| - |R_{xs}|)$$

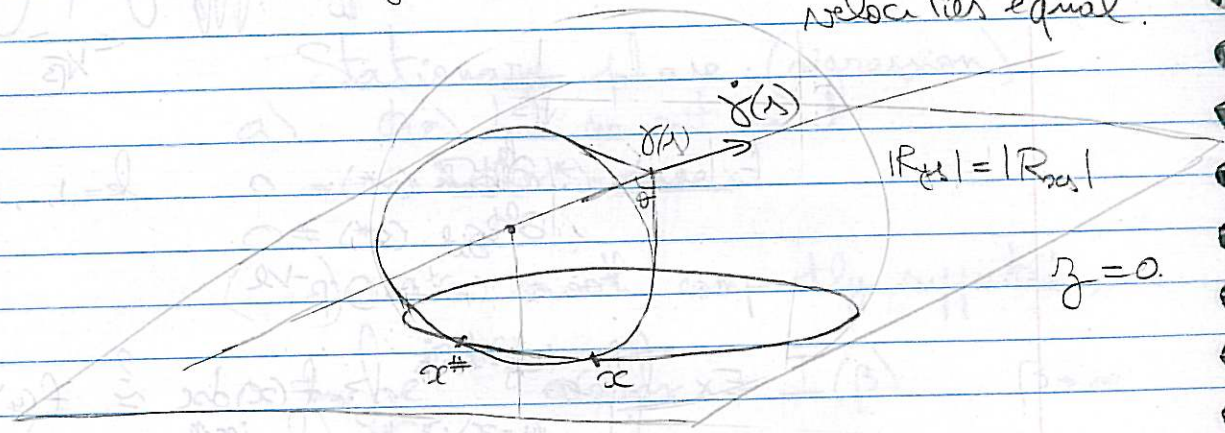
$$\frac{\partial \Phi}{\partial s} = 2\omega'(\hat{R}_{ys} \cdot \dot{\gamma}(s) - \hat{R}_{xs} \cdot \dot{\gamma}(s))$$

$\frac{\partial^2 \Phi}{\partial \omega' \partial s}$ singular but OK.

Stat. phase points:

$|R_{ys}| = |R_{xs}|$ ranges equal.

$\hat{R}_{ys} \cdot \dot{\gamma}(s) = \hat{R}_{xs} \cdot \dot{\gamma}(s)$ down-range velocities equal.



\Rightarrow 2 solutions: x and its mirror onto flight plane

Side(ways) looking antenna:
 $y = x$ only stat phase point.

Transform $2w |R_{yy}| - |R_{xx}| = (y-x) \cdot \dots$

something

Trick: $f(y) - f(x) = \int_0^1 \frac{d}{d\mu} f(x + \mu(y-x)) d\mu = (y-x) \cdot \int_0^1 (\nabla f)(x + \mu(y-x)) d\mu$

Here: $f(x) = 2w \|R_{xx}\|$
 $\nabla f(x) = -2w R_{xx} \cdot P_2$

$\int_0^1 (\nabla f)(x + \mu(y-x)) d\mu \approx -2w R_{xx} \cdot P_2$ when $y \approx x$.
 \Rightarrow call this $E(x, w, \lambda)$. (capital xi)
 $\Rightarrow 2w (|R_{yy}| - |R_{xx}|) \approx (y-x) \cdot E(x, w, \lambda)$

$K(y, x) = \int e^{i(y-x) \cdot E(x, w, \lambda)} Q(\omega, \lambda, y_T) \times \bar{A}(\omega, \lambda, x_T) d\omega d\lambda$

Change var $(\omega, \lambda) \rightarrow \xi = E(x, \omega, \lambda) \in \mathbb{R}^2$.
 $Q(\omega, \lambda, y_T) \rightarrow Q(\xi, y_T)$ (slight abuse)
 $\bar{A}(\omega, \lambda, x_T) \rightarrow \bar{A}(\xi, x_T)$ (abuse)
 $d\omega d\lambda = \left| \frac{\partial(\omega, \lambda)}{\partial \xi} \right| d\xi$

\hookrightarrow Beyerlein determinant.

$K(y, x) = \int e^{i(y-x) \cdot \xi} Q(\xi, y_T) \bar{A}(\xi, x_T) \left| \frac{\partial(\omega, \lambda)}{\partial \xi} \right| d\xi$

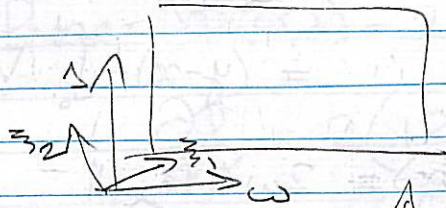
\triangle intro χ .

Want $Q(\xi, x_T) = \frac{1}{A(\xi, x_T) \left| \frac{\partial(\omega, \lambda)}{\partial \xi} \right|} \chi(\xi, x)$ prevent division by zero.

$K(y, x) \approx \int e^{i(y-x) \cdot \xi} d\xi \approx \delta(y-x)$
Rmk about matched filtering: see next page.

quintons

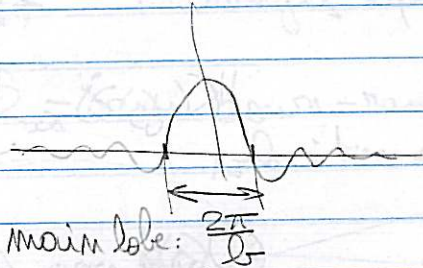
$$\text{Resolution} = K(y, x) = \int_{\text{data manifold } M} e^{i(y-x) \cdot \vec{z}} d\vec{z}$$



Still $M \approx M_1 \times M_2$

$$\int_M e^{i(y-x) \cdot \vec{z}} d\vec{z} = \int_{M_1} e^{i(y_1-x_1) \cdot \vec{z}_1} d\vec{z}_1 \times \int_{M_2} e^{i(y_2-x_2) \cdot \vec{z}_2} d\vec{z}_2$$

$$\text{and } \int_{-b}^b e^{iak} dk = \frac{2 \sin bx}{x} = 2b \text{sinc}(bx)$$



$[-b, b]$ in k domain \Rightarrow resolution $\frac{2\pi}{b}$ in x domain
(Shannon sampling theorem)

Imp. matched filtering

\bar{A} includes $\bar{P}(\omega)$

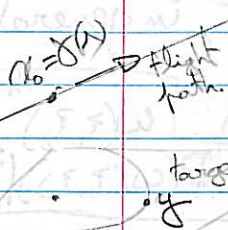
$$\begin{aligned} &\rightarrow \int p(t-t) d(t') dt \\ &= p(t) \star d(t) \end{aligned}$$

Today:
X-ray tomography

(1)

$$\frac{e^{i\omega \|x-y\|}}{\|x-y\|} \approx \frac{e^{i\omega \|x\|}}{\|x\|} e^{-i\omega \hat{x} \cdot y} \text{ when } \|x\| \gg \|y\|$$

11/03/01 Far field approx $\Rightarrow y_{\mp} = \begin{pmatrix} y \\ 0 \end{pmatrix}$



$$D(s, \omega) = \int_{\mathbb{R}^2} e^{2i\omega \|x_0 - y_{\mp}\|} A(y, s, \omega) V(y) dy$$

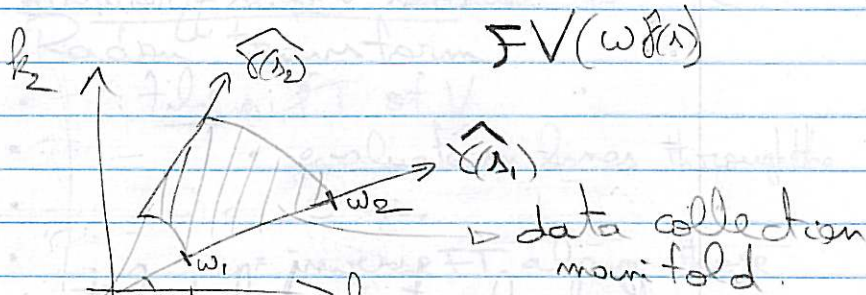
have knocked off two integrals over the antenna \rightarrow antenna beam patterns in A . and $x_0 = \delta(s)$



(1) Go further: $e^{2i\omega \|x_0 - y_{\mp}\|} \approx e^{2i\omega \|x_0\|} e^{-2i\omega \hat{x}_0 \cdot y_{\mp}}$
 also $A(y, s, \omega) \perp y$, say $A = P(\omega)$
 (localized target, nondirectional beam)

$$D(s, \omega) \sim e^{2i\omega \|x_0\|} P(\omega) \int_{\mathbb{R}^2} e^{-i\omega \hat{x}_0 \cdot y} V(y) dy = FV$$

$\underbrace{\hspace{10em}}_{F_{\text{approx}}}$



Adjoint $(F^*D)(y) = \int_{\mathcal{M}} e^{-2i\omega \|x_0\|} P(\omega) e^{i\omega \hat{x}_0 \cdot y} D(s, \omega) ds d\omega$

Composition: $F^*FV =$

$$\int ds d\omega |P(\omega)|^2 \int e^{i\omega \hat{x}_0 \cdot (y-z)} V(z) dz \stackrel{?}{=} V(y)$$

let $\omega \hat{x}_0 = \begin{pmatrix} \xi_1 \\ \xi_2 \end{pmatrix} = \Xi(\omega, s)$

$$d\xi_1 d\xi_2 = \left| \frac{\partial(\xi_1, \xi_2)}{\partial(s, \omega)} \right| ds d\omega$$



ex. $s = \theta$, and $\hat{x}_0 = \begin{pmatrix} \cos \theta \\ \sin \theta \end{pmatrix}$. $\theta \in [0, 2\pi)$

① $y \in \mathbb{R}^m$, here $m=2$.

then $\left| \frac{\partial(x, \omega)}{\partial(\xi_1, \xi_2)} \right| = \omega^{m-1}$ (polar to cartesian)

⚠ Data manifold $d = \mathbb{R}^2$. \hookrightarrow Bessel's determinant in general

then $\int d\xi_1, d\xi_2 e^{i\xi \cdot (y - \xi)} V(\xi) \frac{\omega(\xi_1, \xi_2)^{m-1}}{|P(\omega(\xi_1, \xi_2))|^2}$
not too far from $V(y)$

\rightarrow instead of F^* , use

$$(BD(x, \omega))(y) = \int e^{-2i\omega \|\hat{\gamma}(x)\|} \bar{P}(\omega) \cdot \left(\frac{\omega^{1-m}}{|P(\omega)|^2} \right)$$

then $\hat{\gamma}(x) = \begin{pmatrix} \cos \theta \\ \sin \theta \end{pmatrix}$, get $\times e^{i\omega \hat{\gamma}(x) \cdot y} \int D(x, \omega) dx d\omega$

$$BFV(y) = \int d\xi_1, d\xi_2 e^{i\xi \cdot (y - \xi)} V(\xi) = V(y)$$

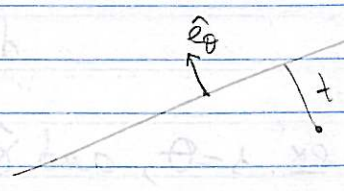
- Simplify:
- full data manifold
 - for field, twice
 - ok to divide by $|P(\omega)|^2$
 - $\hat{\gamma}(x) = \begin{pmatrix} \cos \theta \\ \sin \theta \end{pmatrix}$

Important special case $\hat{\gamma}(x) = \begin{pmatrix} \cos \theta \\ \sin \theta \end{pmatrix} = \hat{e}_\theta$, $e^{2i\omega \|\hat{\gamma}(x)\|} P(\omega) = 1$

$$D(\theta, \omega) = (FV)(x, \omega) = \left[e^{2i\omega \|\hat{\gamma}(x)\|} P(\omega) \right] (FV)\left(\omega \begin{pmatrix} \cos \theta \\ \sin \theta \end{pmatrix}\right) = 1$$

$$d(\theta, t) = \int e^{+i\omega t} \left(\int e^{-i\omega \hat{e}_\theta \cdot y} V(y) dy \right) d\omega$$

$$= \int \delta(t - \hat{e}_\theta \cdot y) V(y) dy$$



S = n meters $\omega = \omega$

Integration along a line with

- angle θ
- offset t

→ Radon transform ($n=2$)
 (or X-ray transform ($n=3$))

$$d(\theta, t) = (R V(x))(\theta, t)$$

we have $(F_t d(\theta, t))^{(\theta, \omega)} = D(\theta, \omega) = (F_y V)(\omega(\frac{\cos \theta}{\sin \theta}))$

$$\Rightarrow (F_t R V(\theta, t))(\theta, \omega) = (F_y V)(\omega(\frac{\cos \theta}{\sin \theta}))$$

$$R V(\theta, t) = F_t^* \left(F_y V \left(\omega \left(\frac{\cos \theta}{\sin \theta} \right) \right) \right) (\theta, t)$$

→ Fourier-slice theorem for the Radon transform

- FT of V
- evaluate on lines through the origin
- inverse FT along those lines

$$F^* D(y) = \int \int e^{-2i\omega \|\hat{x}(s)\|} \overline{P(\omega)} e^{i\omega \hat{s}(s) \cdot y} D(\theta, \omega) d\theta d\omega$$

$$= \int \int e^{i\omega \hat{e}_\theta \cdot y} \left(\int e^{-i\omega t} d(\theta, t) dt \right) d\theta d\omega$$

$$= \int \int \delta(t - \hat{e}_\theta \cdot y) d(\theta, t) d\theta dt$$

! full data

→ Backprojection

$$\omega^{1-m} = \omega^{-1} \text{ when } m=2 \quad (4)$$

$$\begin{aligned} \text{(BD)}(y) &= \int e^{i\omega \hat{e}_\theta \cdot y} \cdot \omega^{-1} D(\theta, \omega) d\theta d\omega \\ &= \int \frac{e^{i\omega \hat{e}_\theta \cdot y}}{\omega} \int e^{-i\omega t} d(\theta, t) d\theta dt \\ &= \int F\left[\frac{1}{\omega}\right](t - \hat{e}_\theta \cdot y) d(\theta, t) d\theta dt \\ &= V(y) \quad \text{with } F\left[\frac{1}{\omega}\right](t) = 2i \operatorname{sgn}(t) \end{aligned}$$

ω^{-1} is a filter \rightarrow Filtered backprojection.

Since $d(\theta, t) = (R V(y))(\theta, t)$,
filtered backprojection is R^{-1} .

$$(R^{-1} d(\theta, t))(y) = \int e^{i\omega \hat{e}_\theta \cdot y} \frac{1}{\omega} \int e^{-i\omega t} d(\theta, t) d\theta dt$$

Rank: R, R^{-1} basis for X-ray tomography
(3D: $\omega^{1-m} = \omega^{-2}$)

\rightarrow CT scans.

• $FV(\omega \hat{e}_\theta)$ like in crystallography
also quantum scattering

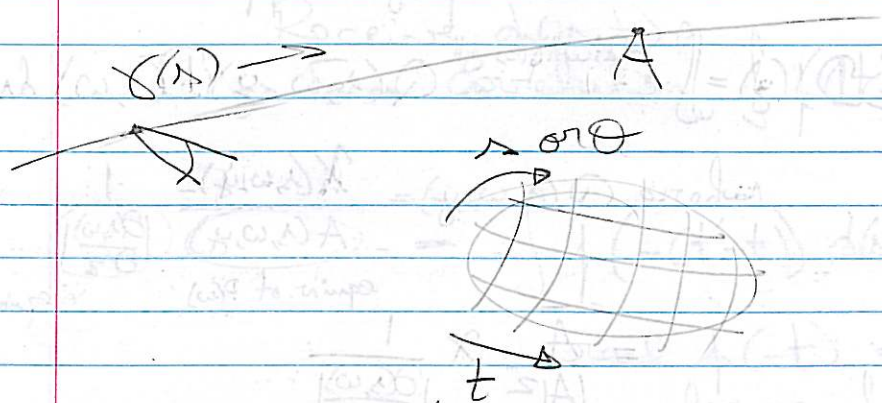
• Coincidentally, similar equations in MRI although not same physics
(absorption of radio waves)

• Radar:

$$e^{2i\omega \|\delta(s) - y\|} \approx e^{2i\omega \|\delta(s)\|} e^{-2i\omega \delta(s) \cdot y}$$

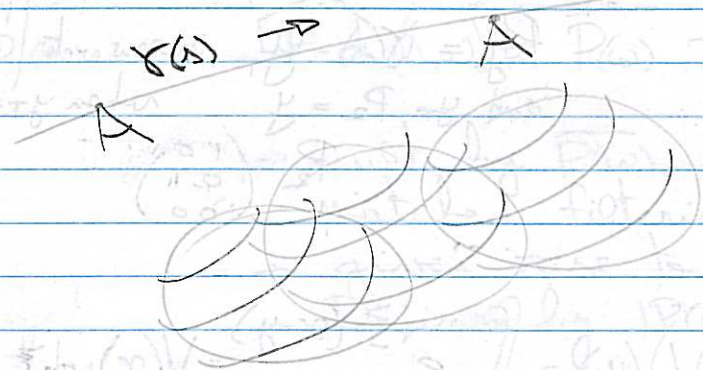
(also: ultrasound and sonar)

OK for spotlight SAR. (stores)



circles approx by lines

also OK for moving-targets SAR. (air traffic control, etc.)
not OK for stripmap SAR.



→ back to formula
$$D(s, \omega) = \int_{\mathbb{R}^2} e^{2i\omega \|x(s) - y\|} A(y, s, \omega) V(y) dy$$

then
$$d(s, t) = \int f''(t - 2\|x(s) - y\|) V(y) dy$$

$\Delta A = P(\omega)$
 $f(A) = F^{-1} P(\omega) \Delta$

integration along circles

→ generalized Radon transform
if $f'' = \delta$

still, F^* is backprojection
 B is filtered backprojection
 (good approx. to inverse)

$$(BD)(y) = \int e^{-2i\omega \|x-y\|} Q(s, \omega, y) D(s, \omega) ds d\omega$$

$$\text{where } Q(s, \omega, y) = \frac{X(s, \omega, y)}{A(s, \omega, y)} \frac{1}{\left| \frac{\partial(s, \omega)}{\partial \xi} \right|}$$

equiv. of $P(\omega)$ equiv. of ω^{n-1}

$$= \frac{\bar{A}}{|A|^2} X \frac{1}{\left| \frac{\partial(s, \omega)}{\partial \xi} \right|}$$

and here $\xi = \Xi(s, \omega, y) = -2\omega \hat{R}_{ys} \cdot P_2$
 (right approx.)

$R_{ys} = x(s) - y_T$ in row vector (extract 1st two components)
 and $y_T P_2 = y$ when $y_T = (y, 0)$

$$P_2 = \begin{pmatrix} 1 & 0 \\ 0 & 1 \\ & & 0 & 0 \end{pmatrix}$$

then

$$(BFV)(y) = \int_{\text{data manifold } M} e^{i\xi \cdot (x-y)} V(x) d\xi dx$$

(restriction created by $X(s, \omega)$)

$$= \int K(y, x) V(x) dx$$

where $K(y, x) = \int_M e^{i\xi \cdot (x-y)} d\xi$

- ⚠
- OK to divide by $A(s, \omega, y)$, i.e. $P(\omega) \neq 0$, nondirectional beams
 - (far-field)
 - (Ξ approx.)
 - sampling of M .

Rank about matched filtering, (somewhere earlier)

Send $p(t)$

Receive $d(s, t)$

→ Form correlation

$$\int p(t-t') d(s, t') dt' = (MFd)(s, t)$$

$$= \int p(-(t'-t)) d(s, t) dt' = (p(t) \star d)(s, t)$$

real $p(-t) \leftrightarrow \overline{P(\omega)}$

because $\int e^{-i\omega t} p(-t) dt$

$$= \int e^{i\omega t} p(t) dt$$

$$= P(-\omega) \text{ or } \overline{P(\omega)}$$

Convolution theorem \Rightarrow

$$\widehat{MFd}(s, \omega) = \overline{P(\omega)} D(s, \omega)$$

→ mult by $\overline{P(\omega)}$ is the effect of matched filtering, present in F^*

→ gives rise to $|P(\omega)|^2$

Division by $|P(\omega)|^2$ present in B

11/05/09

Error in the last few pages!

Radon transform:

$$(FV)(\theta, \omega) = \int_{\mathbb{R}^2} e^{-i\omega \hat{e}_\theta \cdot y} V(y) dy$$

$$(F^*D)(y) = \int_{\mathbb{R}_+ \times [0, 2\pi]} e^{i\omega \hat{e}_\theta \cdot y} D(\theta, \omega) d\theta d\omega$$

$$(F^*FV)(y) = \int e^{i\omega \hat{e}_\theta \cdot (y-z)} V(z) d\theta d\omega dz$$

$$= \int \frac{1}{\omega(\xi)} e^{i\xi \cdot (y-z)} V(z) d\xi dz$$

$$\left(\left| \frac{\partial \xi}{\partial \omega} \right| = \omega \text{ or } \left| \frac{\partial \omega}{\partial \xi} \right| = \frac{1}{\omega} \right) \text{ with } \omega(\xi) = \sqrt{\xi_1^2 + \xi_2^2}$$

ω and not $1/\omega$

Def. $(BD)(y) = \int_{\mathbb{R}_+ \times [0, 2\pi]} \omega e^{i\omega \hat{e}_\theta \cdot y} D(\theta, \omega) d\theta d\omega$

then $(BFV)(y) = \int \frac{\omega}{\omega} e^{i\xi \cdot (y-z)} V(z) d\xi dz$

$$= V(y)$$

$\Rightarrow B = F^{-1}$

$$RV(\theta, t) = \left[\int_{\omega \rightarrow t} (FV) \right](\theta, t) \quad R = \int_{\omega \rightarrow t} F$$

$$R^{-1} = F^{-1} \int_{\omega \rightarrow t}^{-1} = \left[B \int_{\omega \rightarrow t}^* \right]$$

filtered backprojection
 $\omega = \text{filter}$

Alternative $R^{-1} = (R^*R)^{-1} R^*$
 with $R^* = F^* \int_{\omega \rightarrow t}^*$

$$R^*R = F^* \int_{\omega \rightarrow t}^* \int_{\omega \rightarrow t} F = F^*F$$

with $F^*F = \int_{\xi \rightarrow y}^* \left(\frac{1}{\sqrt{\xi_1^2 + \xi_2^2}} \right) \int_{y \rightarrow \xi}$

$$= (-\Delta)^{-1/2} \quad \text{because } (-\Delta) \leftrightarrow \xi_1^2 + \xi_2^2$$

also, $F^*Fu(y) = \int \frac{c}{\|x-y\|} u(y) dy$

$$F^*F = (-\Delta)^{1/2}, \quad R^{-1} = \left[(-\Delta)^{1/2} F^* \int_{\omega \rightarrow t}^* \right]$$

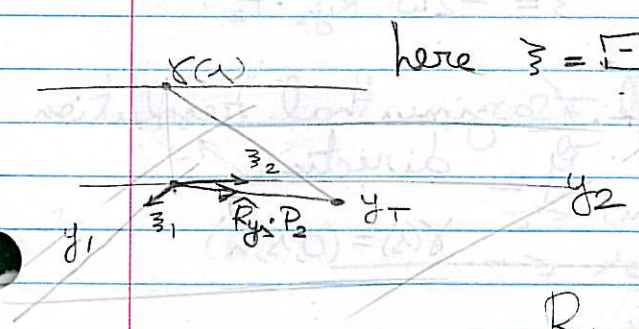
filter backpr.

Resolution: Back to radon

$$(BFV)(y) = \int_{\substack{\text{data} \\ \mathbb{R}^2 \\ \text{manifold } M}} e^{i\vec{\omega} \cdot (y-x)} V(x) d\vec{\omega} dx$$

$$= \int_{\mathbb{R}^2} K(y,x) V(x) dx$$

$$\text{kernel } K(y,x) = \int_M e^{i\vec{\omega} \cdot (y-x)} d\vec{\omega}$$

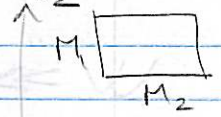


here $\vec{\omega} = \Gamma(1, \omega, y) = -2\omega \hat{R}_{y_2} \cdot P_2 = \begin{pmatrix} \omega_1 \\ \omega_2 \end{pmatrix}$

$-2\omega \times$ (first 2 components of \hat{R}_{y_2})

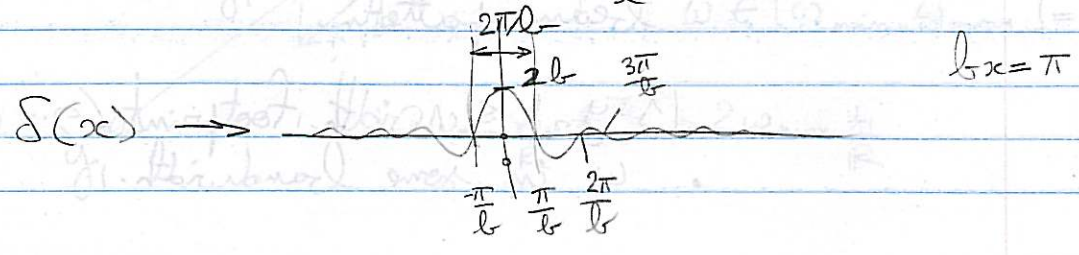
$$\hat{R}_{y_2} = \frac{R_{y_2}}{\|R_{y_2}\|} = \frac{y_2 - y_1}{R} = \begin{pmatrix} y_2 - y_1 \\ 0 \end{pmatrix}, R = \text{range}$$

Assume $M = M_1 \times M_2$



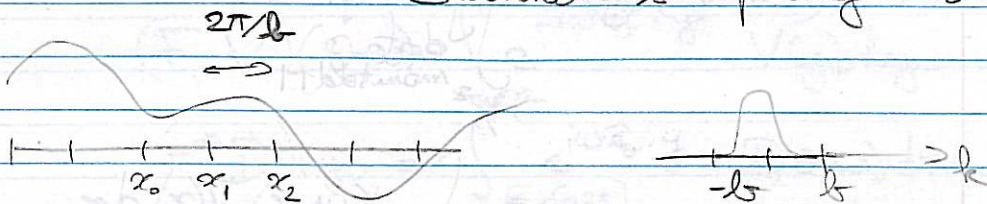
$$\int_M e^{i\vec{\omega} \cdot (y-x)} d\vec{\omega} = \int_{M_1} e^{i\vec{\omega}_1 \cdot (y_1-x_1)} d\vec{\omega}_1 \times \int_{M_2} e^{i\vec{\omega}_2 \cdot (y_2-x_2)} d\vec{\omega}_2$$

with $\int_{-b}^b e^{ixk} dk = \int_{-b}^b \cos(xk) dk = \frac{1}{x} \sin xk \Big|_{-b}^b = \frac{2 \sin xb}{x} = 2b \text{sinc}(bx)$



→ resolution of $\frac{2\pi}{\delta}$ when \int_{-b}^b

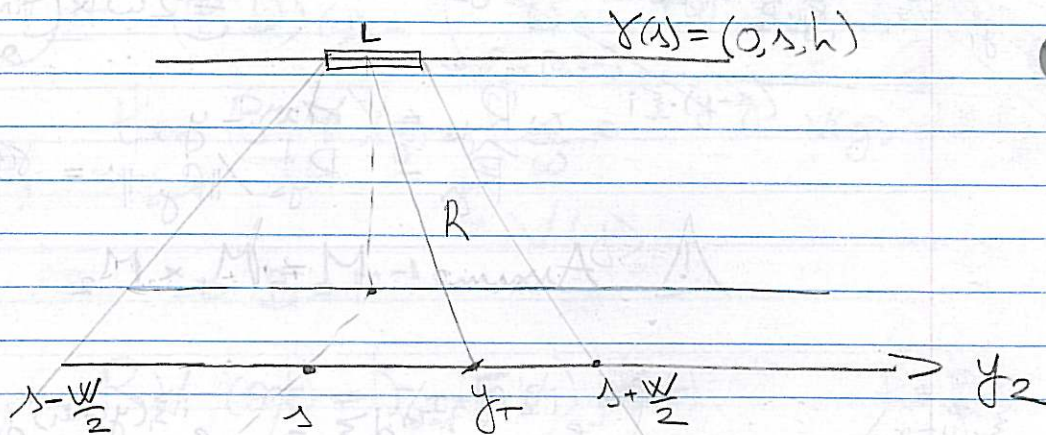
(Nyquist rate -
Shannon sampling theorem)



more generally, resolution of $\frac{4\pi}{\|M_1\|} \times \frac{4\pi}{\|M_2\|}$

Back to M: $\xi = -2\omega \hat{R}_{y_2} - P_2$

① Along-track / azimuthal resolution
→ in the y_2 direction.



$$\xi_2 = -2\omega \frac{(\delta(s) - y_T)_2}{\|\delta(s) - y_T\|} = -2\omega \frac{(s - y_T)}{R}$$

orange

Data manifold in ξ_2 :

- s and y_2 both in the antenna beam pattern.

$$|s - y_2| \leq \text{width footprint} / 2$$

- ω in some bandwidth.

Width of an antenna footprint:

- wavelength $\lambda = 2\pi/k$ $f_k = \omega$
($c=1$)
- length L
- range R

$$W = R \left(\frac{2\lambda}{L} \right)$$

$$\Rightarrow |\tilde{m}_2| = 2\omega \frac{|1 - y_2|}{R} / 2$$

$$b = \max_{\tilde{m}_2} |\tilde{m}_2| = \omega \frac{W}{R} = \frac{\omega}{R} R \left(\frac{2\lambda}{L} \right)$$

$$= \frac{2\omega}{L} \frac{2\pi}{\omega} = \frac{4\pi}{L}$$

$$\Rightarrow \frac{2\pi}{b} = \boxed{\frac{L}{2}} \text{ azimuthal resolution.}$$

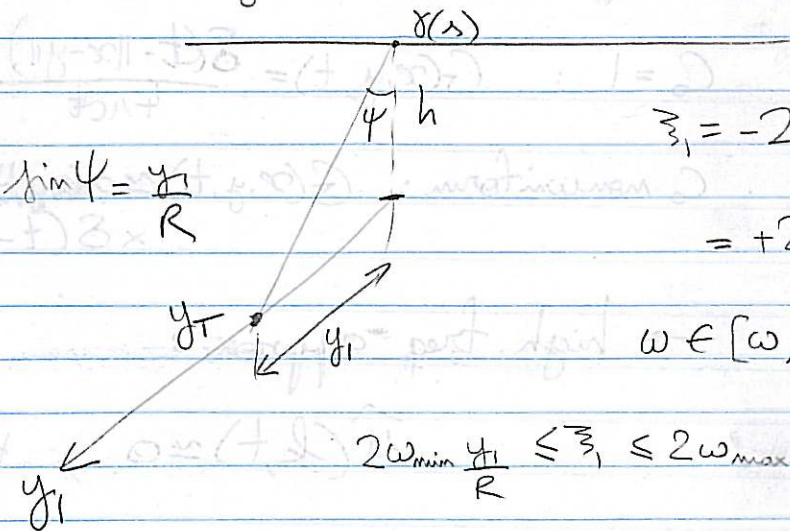
→ indep ω, R

→ small antenna is better so that the beam pattern is wide very counter-intuitive

$$\text{Aperture} = W = R \left(\frac{2\lambda}{L} \right)$$

$$\text{Synthetic aperture: } L/2 \ll W$$

② Range resolution



$$\sin \phi = \frac{y_1}{R}$$

$$\tilde{z}_1 = -2\omega \frac{(x(s) - y_T)}{R}$$

$$= +2\omega \frac{y_1}{R}$$

$$\omega \in [\omega_{\min}, \omega_{\max}] = \text{supp } P(\omega)$$

$$2\omega_{\min} \frac{y_1}{R} \leq \tilde{z}_1 \leq 2\omega_{\max} \frac{y_1}{R}$$

$$2\omega_{\min} \sin \psi \leq \omega_1 \leq 2\omega_{\max} \sin \psi$$

$$2\Delta\omega = \Delta\omega_1 = 2\Delta\omega \sin \psi$$

$$\frac{2\pi}{\Delta\omega} = \frac{2\pi}{(\Delta\omega) \sin \psi}$$

→ better when wider freq. band

→ better when far from antenna

$$(\psi \approx \frac{\pi}{2})$$

(proj. of the range vector
is longer)

(amplitude decreases though)

Chap. 5, Geom. optics & GRT.

Back to

$$u_{\text{inc}}(x, t) = \iint G(x, y, t-s) f(y, s) dy ds$$

$$u_{\text{ref}}(x, t) = \iint G(x, y, t-s) \frac{\partial^2 u_{\text{inc}}(y, s)}{\partial t^2} V(y) dy ds$$

$$c_0 = 1 : G(x, y, t) = \frac{\delta(t - \|x - y\|)}{4\pi \|x - y\|^2}$$

$$c_0 \text{ nonuniform: } G(x, y, t) \approx a(x, y, t) \times \delta(t - \tau(x, y))$$

→ high freq. approx.:

$$\hat{u}(k, t) \approx 0, \quad |k| \leq N$$

18.325. 11/10/09.

$$\left(\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \Delta\right) u = f$$

$$\Rightarrow u_{\text{inc}}(x, t) = \int_0^t \int_{\mathbb{R}^3} G(x, y, t-s) f(y, s) dy ds$$

$$u_{\text{scat}}(x, t) = \int_0^t \int_{\mathbb{R}^3} G(x, y, t-s) \left(\frac{\partial^2}{\partial t^2} V\right)(y, s) dy ds$$

Structure of G ? with $\left(\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \Delta\right) G(x, y, t) = \delta(x-y) \delta(t)$

$$c=1 \quad G(x, y, t) = \frac{\delta(ct - |x-y|)}{4\pi c^2 t} \quad (x \in \mathbb{R}^3, ct = |x-y|)$$

$$= \frac{1}{2\pi} \frac{1}{\sqrt{c^2 t^2 - |x-y|^2}} \quad (x \in \mathbb{R}^2)$$

Progressing wave expansion / geometrical optics (Goursat & Hilbert)

$\tau(x, y) = \frac{|x-y|}{c}$ is the traveltime between x and y

$$\text{Now } G(x, y, t) = a(x, y) S(t - \tau(x, y)) + R(x, y, t)$$

R smoother than G , not small

a, τ smooth in an as-yet unspecified region.

$$\left(\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \Delta\right) G = \delta(x-y) \delta(t) \quad (\text{def.})$$

$$= 0 \quad (t \neq 0)$$

$$\frac{\partial^2}{\partial t^2} (a S(t-\tau)) = a S''(t-\tau)$$

$$\Delta = \nabla \cdot \nabla \quad \nabla(a S(t-\tau)) = (\nabla a) S - a(\nabla \tau) S'$$

$$\nabla \cdot (\nabla(a S(t-\tau))) = \Delta a S - (\nabla a) \cdot (\nabla \tau) S' - (\nabla a) \cdot (\nabla \tau) S' - a(\Delta \tau) S' + a|\nabla \tau|^2 S''$$

$$\left(\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \Delta\right) G = a \left(\frac{1}{c^2} - |\nabla \tau|^2\right) S'' + (2\nabla \tau \cdot \nabla a + a \Delta \tau) S' + \Delta a S + \left(\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \Delta\right) R = 0.$$

In decreasing order of singularity,

$$(2) \quad \frac{1}{c^2} - |\nabla \tau|^2 = 0 \quad \rightarrow \text{det. } \tau$$

$$(1) \quad 2\nabla a \cdot \nabla \tau + a \Delta \tau = 0 \quad \rightarrow \text{det. } a$$

(0) choose R such that $\Delta a S + \left(\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \Delta\right) R = 0$
 f. (4) here Remark: Geom. optics & Helmholtz equation. (*)

Eikonal equation of geometrical optics.

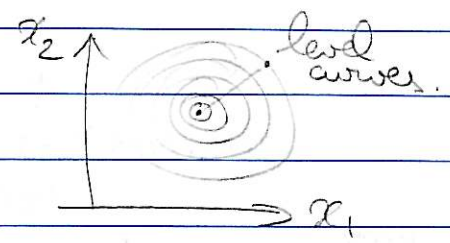
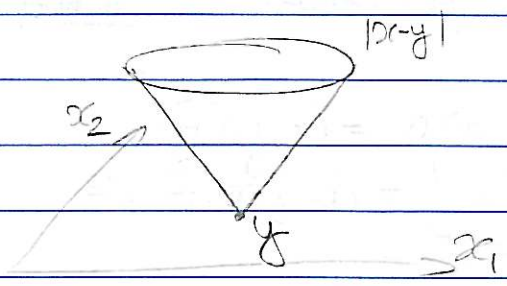
Fix y, PDE in x:

$$|\nabla_x \tau(x, y)| = \frac{1}{c(x)}, \quad \tau(y, y) = 0$$

$$\text{ex } |\nabla_x \tau| = 1 \quad \tau(y, y) = 0$$

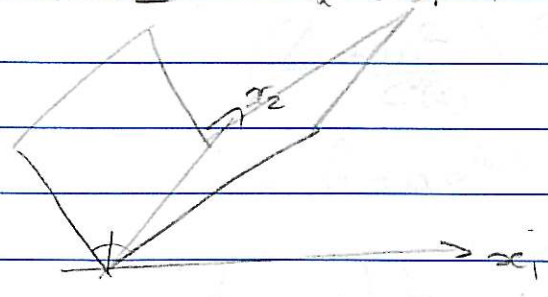
$$\Rightarrow \tau(x, y) = |x - y| \quad \text{distance function}$$

$$\nabla_x \tau = \frac{x - y}{|x - y|} \quad \text{obeys } |\nabla_x \tau| = 1$$



ex $|\nabla_x \tau(x)| = 1$

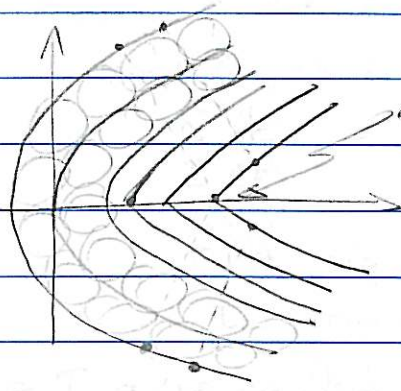
$\tau(x) = 0$ when $x_1 = 0$



ex. $|\nabla_x \tau(x)| = 1$

$\tau(x) = 0$ along $x_1 = x_2^2$

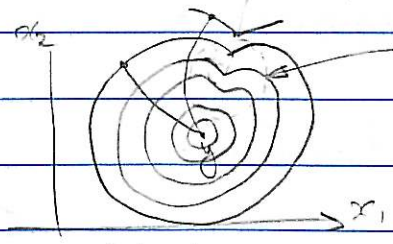
Huygen's principle.



- discontinuity of $\nabla \tau$.
- jump in the point on the parabola as $\epsilon \rightarrow 0$ to which the distance is measured.
- when normal lines cross. (rays).

ex. $|\nabla_x \tau(x)| = \frac{1}{c(x)}$

locally looks like $\frac{|x-y|}{c(y)}$



ex (1D) $\left| \frac{d\tau}{dx} \right| = \frac{1}{c(x)} \Rightarrow \tau(x,y) = \int_x^y \frac{1}{c(z)} dz$ ✓
 $\tau(y,y) = 0$

Rank: Geom. optics

$$G(x,y,t) = a(x,y) \delta(t - \tau(x,y))$$

$$\Rightarrow \hat{G}(x,y,\omega) = \int e^{i\omega t} a(x,y) \delta(t - \tau(x,y)) dt$$

$$= a(x,y) e^{i\omega \tau(x,y)}$$

(generalization of $\frac{e^{i\omega|x-y|}}{4\pi|x-y|}$)

Green's function for the Helmholtz equation:

$$\left(\Delta + \frac{\omega^2}{c^2(x)}\right) \hat{G}(x,y,\omega) = \delta(x-y)$$

Try $\hat{G} = ae^{i\omega \tau}$ and equate like-powers of ω
 \Rightarrow get (2), (1).

Asymptotic matching as $\omega \rightarrow \infty$.

High-frequency approximation to the actual \hat{G}
(more accurate as $\omega \rightarrow \infty$)

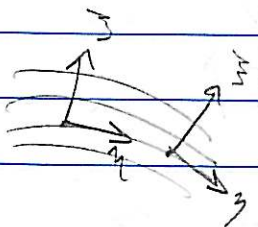
Smoothness in t	\Leftrightarrow	Asymptotics as $\omega \rightarrow \infty$
$\delta''(t)$		ω^2
$\delta'(t)$		ω
$\delta(t)$		1
$H(t)$ (step)		$1/\omega$
t_+ (ramp)		$1/\omega^2$

$\Rightarrow \tau$ also called a phase

Rank: Characteristic surfaces:

- wave eq. $\left(\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \Delta\right) u = 0$, $-\frac{1}{c^2} \left(\frac{\partial \phi}{\partial t}\right)^2 + |\nabla \phi|^2 = 0$ (Ham-Jac.)

- Helmholtz, $\left(-\frac{\omega^2}{c^2} - \Delta\right) u = 0$, $-\frac{\omega^2}{c^2} + |\nabla \phi|^2 = 0$
then $\phi = \omega \tau$



Ans: Show that under (eikonal)

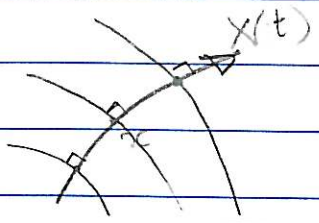
$(x,y) \rightarrow (\phi, \eta)$, such that $\phi = \omega \tau$, there exist (approx.) solutions of HE that oscillate in \bar{z} and are const. in η ,

locally: $\left[\frac{\partial^2 u}{\partial \bar{z}^2} + \omega^2 u = 0\right]$ and $\left[\frac{\partial^2 u}{\partial \eta^2} = 0\right]$ (approx.)

⊗ Back to 1.2

Method of characteristics

Ray method for determining τ :
curves along which the traveltime is calculated.



$$|\nabla\tau| = \frac{1}{c(x)}$$

consider $X(t)$: ray
↳ parameter

$$\frac{d}{dt} \tau(X(t)) = \nabla\tau(X(t)) \cdot \dot{X}(t)$$

$\dot{X}(t) \parallel \nabla\tau$
 $\Rightarrow X(t)$ perpendicular
to level lines of τ
locally polar
coordinates

Choose $\dot{X}(t) = c(X(t)) \cdot \frac{\nabla\tau}{|\nabla\tau|}$

$$\dot{X}(t) = c^2(X(t)) \cdot \nabla\tau(X(t)) \quad \text{ODE}$$

then $\nabla\tau \cdot \dot{X} = c^2 |\nabla\tau|^2 = 1$.

$$\Rightarrow \tau(X(t)) - \tau(X(t_0)) = t - t_0$$

and τ is identified with t

$X(t)$ still depends on τ though.

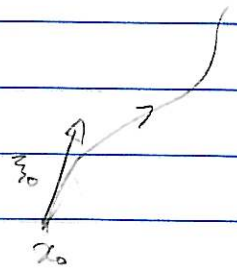
Record direction:

Put $\xi(t) = \nabla\tau(X(t))$ (other possible choice: $\frac{\nabla\tau}{|\nabla\tau|}$)

Then $\dot{\xi}(t) = \nabla\nabla\tau(X(t)) \cdot \dot{X}(t)$

$$\begin{aligned} &= \nabla\nabla\tau \cdot \nabla\tau \cdot c^2 \quad (\text{at } X(t)) \\ &= \frac{1}{2} c^2 \nabla \cdot |\nabla\tau|^2 \\ &= \frac{1}{2} c^2 \nabla(c^{-2}) \\ &= -\frac{1}{2} c^{-2} \nabla(c^2) \\ &= -\frac{1}{2} |\nabla\tau|^2 \nabla(c^2) \\ &= -\frac{1}{2} |\xi(t)|^2 \nabla(c^2(X(t))) \end{aligned}$$

index of τ



$$\Rightarrow \begin{cases} \dot{X}(t) = c^2(X(t)) \xi(t) & X(0) = x_0 \\ \dot{\xi}(t) = -\frac{\nabla c^2(X(t))}{2} |\xi(t)|^2 & \xi(0) = \xi_0 \end{cases}$$

\Rightarrow Solve for the rays about y
 $\tau = \text{arclength along the rays.}$

Remark: $\tau =$ action for the Fermat principle

Remark: PDE in x : Eulerian viewpoint
2ODEs in t : Lagrangian viewpoint.

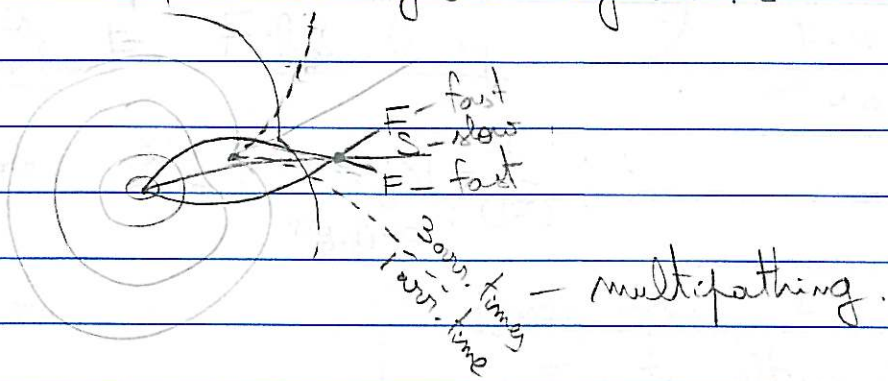
Remark: Rays are charact. curves for the eikonal equation. They are bi-variant curves for HEDWE.

Remark: $H(x, \xi) = \frac{1}{2} c^2(x) |\xi|^2$ (optics acoustics)

then $\dot{x}(t) = \nabla_{\xi} H(x(t), \xi(t))$
 $\dot{\xi}(t) = -\nabla_x H(x(t), \xi(t))$

Hamiltonian system

Remark: Validity of τ ? $\Omega(y)$
 $= \{x: \text{only one ray links } x \text{ to } y\}$



In general, need τ_j as many functions as there are arrival times
 \rightarrow multiple-valued solution to $|\nabla\tau| = 1/c$ (locally)

Otherwise, τ records the 1st arrival. (in which sense? viscosity solution \rightarrow PDE)

Remark: What about a ? Have $2\nabla a \cdot \nabla\tau - a \Delta\tau = 0$
Called a transport equation: along rays

$$\frac{d}{dt} a(x(t)) = \dot{x}(t) \cdot \nabla a(x(t)) = c^2(x(t)) \nabla\tau \cdot \nabla a(x(t)) = c^2(x(t)) a(x(t)) \Delta\tau(x(t)) \neq$$

Next time: a , caustics, slowly, remainder R , then GRT migration. Talk about wavefronts

on ray path

ex: express in terms of $x, \xi, \frac{\partial x}{\partial t}, \frac{\partial \xi}{\partial t}$

\Rightarrow ODE along rays. (do properly next)

11/12/09 Geometrical optics: amplitude

eq. Wave equation Helmholtz eq.

$$\left(\frac{1}{c^2(x)} \frac{\partial^2}{\partial t^2} - \Delta \right) u = 0$$

$$\left(\Delta + \frac{\omega^2}{c^2(x)} \right) u = 0$$

$G(x,y,t) = a(x,y) \times \delta(t - \tau(x,y))$
 Prop. wave expansion
 $t - \tau(x,y)$ is a wave front

Eikonal eq

$\phi(x,y,\omega) \approx a(x,y) e^{i\omega \tau(x,y)}$
 geom. optics
 τ is a phase

Char. eq. Ham-Jacobi here Osher

$$|\nabla_x \tau(x,y)| = \frac{1}{c(x)}$$

$\tau(x,y) = 0$

connection

$\tau(x(t)) - \tau(x(0)) = t$ $x(0) = y$
 τ is traveltime along rays
 $\nabla \tau = \vec{\beta}$ called slowness vector
 τ is the action for the Fermat principle.

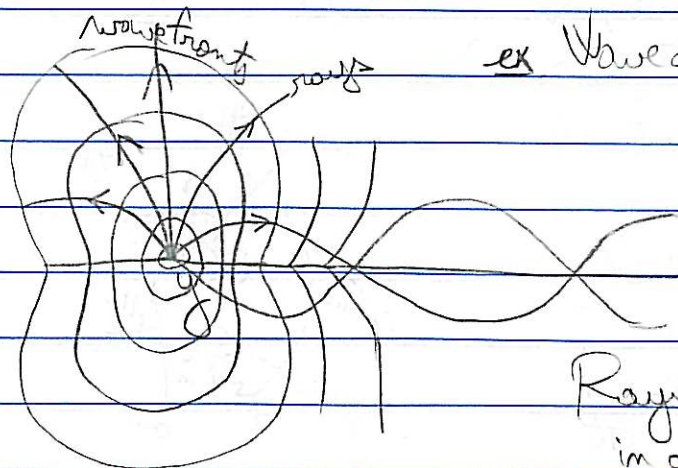
Bichar. eq.

Ray equations:

$$H(x, \vec{\beta}) = \frac{1}{2} c^2(x) |\vec{\beta}|^2$$

$$c^2(x(t)) \vec{\beta}(t) = \begin{cases} \dot{x}(t) = \nabla_{\vec{\beta}} H(x(t), \vec{\beta}(t)) \\ -\frac{1}{2} \nabla_x c^2(x(t)) |\vec{\beta}(t)|^2 = \dot{\vec{\beta}}(t) = -\nabla_x H(x(t), \vec{\beta}(t)) \end{cases}$$

$x(0) = x_0$
 $\vec{\beta}(0) = \vec{\beta}_0$



ex Waveguide:

Rays \perp Wave fronts: only in an isotropic medium

simple geom. optics

Validity: no multiple arrivals. / τ records only 1st arrival otherwise branches of a multivalued function.
 $\nabla \tau$ discontinuous. $\xi = \nabla \tau$ not defined uniquely

Viscosity solution: $|\nabla \tau_\epsilon(x)|^2 + \epsilon \Delta \tau_\epsilon(x) = 0$
 then let $\epsilon \rightarrow 0$.

Eikonal eq. valid for $|y-x|$ not too large.
 Ray equations are valid for all times.

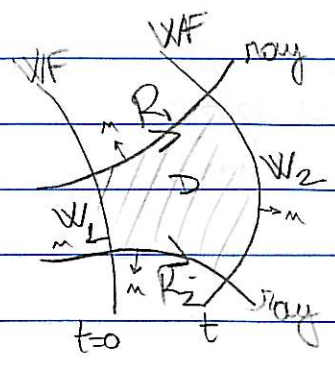
⚠ when using τ , simple geom. optics.

Amplitude: $2 \nabla a \cdot \nabla \tau + a \Delta \tau = 0$. (1)
 dir. der of a in the dir. of $\nabla \tau = \xi$

$$\begin{aligned} \frac{d}{dt} a(x(t), y) &= \dot{x}(t) \cdot \nabla a && \text{(Fix } y) \\ &= c^2 \xi \cdot \nabla a \\ &= c^2 \nabla \tau \cdot \nabla a \\ &= -\frac{c^2}{2} a \Delta \tau && \text{by (1)} \end{aligned}$$

\rightarrow transport equation along rays

Or (1) $\Leftrightarrow \nabla \cdot (a^2 \nabla \tau) = 0$ div. form
 ($2 a \nabla a \cdot \nabla \tau + a^2 \Delta \tau = 0$)



Use the div. thm in \square (ray tube)
 $\iint_D \nabla \cdot (a^2 \nabla \tau) dA$ (surface)
 \parallel
 $\int_{\partial D} a^2 \nabla \tau \cdot n ds$ (line-circulation)

$\partial D = W_1, U, W_2, U, R_1, U, R_2$

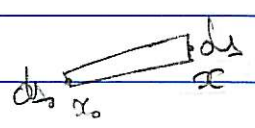
$n \cdot \nabla \tau = 0$ on R_1, R_2

$n \cdot \nabla \tau = -|\nabla \tau|$ on W_1
 $+|\nabla \tau|$ on W_2

$\Rightarrow \int_{W_1} a^2 |\nabla \tau| ds = \int_{W_2} a^2 |\nabla \tau| ds$

$\int_{W_1} \frac{a^2}{c} ds = \int_{W_2} \frac{a^2}{c} ds$

limit $|W_1|, |W_2| \rightarrow 0$, $\frac{d^2(x)}{a^2(x)} \frac{c(x_0)}{c(x)} = \frac{ds}{ds_0}$



$a(x) = a(x_0) \cdot \sqrt{\frac{c(x_0)}{c(x)} \frac{ds}{ds_0}}$

→ conservation of energy (flow down tube)

ex. Plane waves: $\frac{ds}{ds_0} = 1$, $a = \text{const.}$

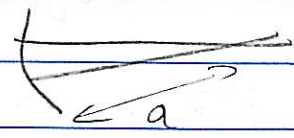
Circ. waves / cylindrical: $\frac{ds}{ds_0} = r$, $a \sim \frac{1}{\sqrt{r}}$

Spherical waves: $\frac{dA}{dA_0} = r^2$, $a \sim \frac{1}{r}$

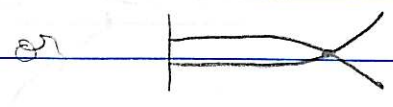
(compare to $\frac{e^{i\omega|x-y|}}{4\pi|x-y|}$)



Focus / Caustic: $a \sim \frac{1}{\sqrt{r-a}}$



→ blow up of the amplitude.



→ brighter spots

$$G(x, y, t) = a(x, y) S(t - \tau(x, y)) + R(x, y, t) \quad (4)$$

Source. (ω^2, S'') $|\nabla \tau| = \frac{1}{c(x)} \Rightarrow$ eq. for τ

(ω', S') : $2 \nabla a \cdot \nabla \tau + a \Delta \tau = 0 \Rightarrow$ eq. for a

(ω^0, S) : $\Delta a S + \left(\frac{1}{c^2} \frac{\partial^2}{\partial t^2} - \Delta \right) R = 0 \Rightarrow$ eq. for R

- Ex.
- R is no worse than discontinuous if $S = \delta$
 - $\int dx \int dt |R(x, y, t)|^2 \leq C \|G\|^4$

Hint: energy estimate, $\begin{cases} \frac{\partial w}{\partial t} - Lw = f \\ \Rightarrow \frac{dE}{dt} = 2 \langle f, w \rangle \\ E(t) = \langle w, w \rangle \end{cases}$

Concl: R is smoother than S

Forward / Imaging operators

$$u_{inc}(x, t) = \iint G(x, y, t-s) f(y, s) dy ds$$

! $f(y, s) = \delta(y - x_0) \delta(s - t)$

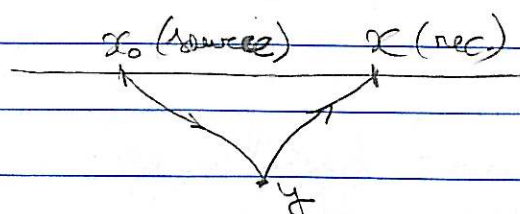
$$= G(x, x_0, t)$$

$$u_{refl}(x, t) = \iint G(x, y, t-s) V(y) \frac{\partial^2 u_{inc}(y, s)}{\partial t^2} dy ds$$

with $V(y) = \frac{2\epsilon c}{c^3} - n$

$$= \int dy V(y) \int ds G(x, y, t-s) \frac{\partial^2 G}{\partial t^2}(y, x_0, s) ds$$

! $G(x, y, t) \approx a(x, y) \delta(t - \tau(x, y))$



! simple geom. optics

$$u_{A,B}(x,t) = \int dy V(y) K(x,y,t) = FV(x,t)$$

$$\text{formal } K(x,y,t) = \int G(x,y,t-s) \frac{\partial^2 G}{\partial t^2}(y,x_0,s) ds$$

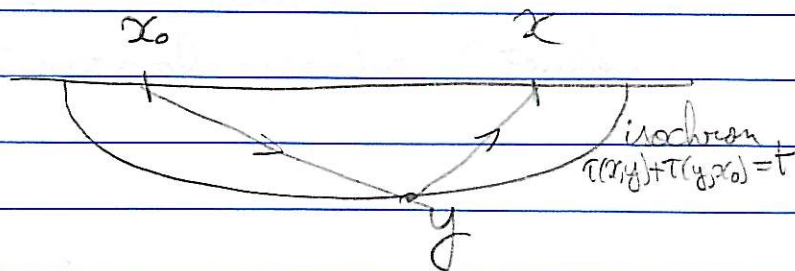
$$\approx \int a(x,y) a(y,x_0) \delta(t-s-\tau(x,y)) \times \delta''(s-\tau(y,x_0)) ds$$

$$= \boxed{a(x,y) a(y,x_0) \delta''(t-\tau(x,y)-\tau(y,x_0))}$$

unclear (Possible when $s = t - \tau(x,y)$ and $s = \tau(y,x_0)$ do not intersect transversally $\Rightarrow \nabla_x \tau(x,y) \neq -\nabla_x \tau(y,x_0)$ no forward scattering \Leftarrow simple geom. optics)

$$\Rightarrow u_{A,B}(x,t) = \underbrace{\frac{\partial^2}{\partial t^2}}_{\text{filtering}} \int \underbrace{a(x,y) a(y,x_0)}_{\text{weighting}} \underbrace{\delta(t-\tau(x,y)-\tau(y,x_0))}_{\text{integral over isochron}} V(y) dy$$

$\{ y = t - \tau(x,y) - \tau(y,x_0) \}$



ellipse/ellipsoid when $c_0 = \text{const.}$
 $\tau(x,y) = \frac{|x-y|}{c_0}$

\rightarrow Generalized Radon trf.,
Kirchhoff modeling

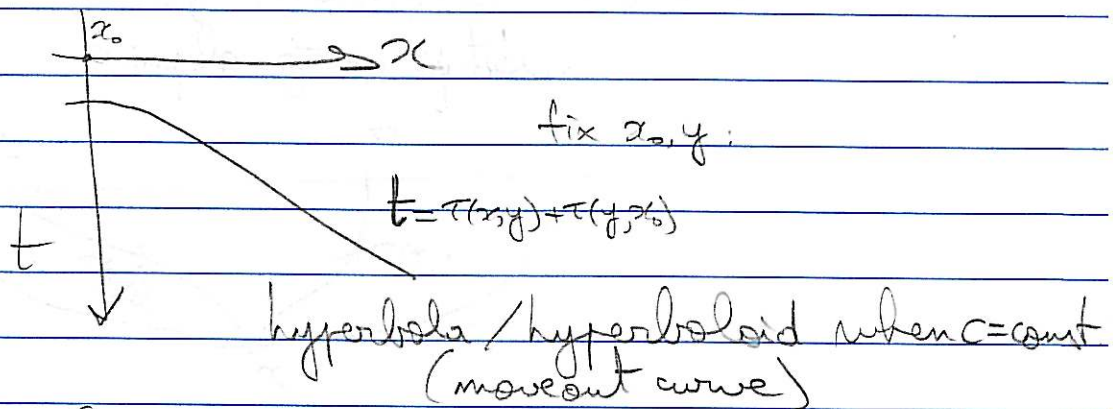
Adjoint: $\langle d, FV \rangle_{(x,t)} = \langle F^*d, V \rangle_{(y)}$

$$\int d(x,t) \int V(y) K(x,y,t) dy dx dt$$

$$\int V(y) \int d(x,t) K(x,y,t) dx dt dy$$

$$\Rightarrow F^*d(y) = \int K(x,y,t) d(x,t) dx dt$$

$$= \int \int a(x,y) a(y,z_0) \delta''(t - \tau(x,y) - \tau(y,z_0)) d(x,t) dx dt$$



→ GRT as well
Kirchhoff migration

Next: multi-offset/sources

GRT normal operator (Beulkin)

11/17/09

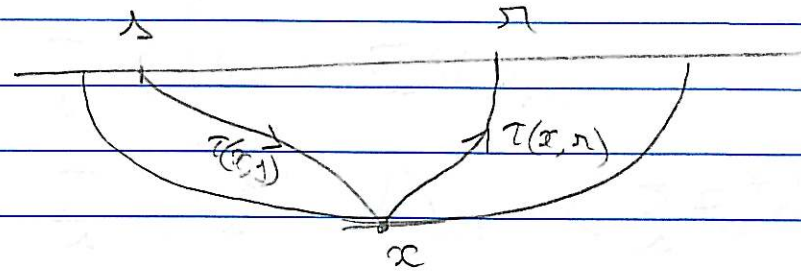
r = receiver s = source, fixed
 t, t' = time x = subsurface position

$$u_{R,B}(r, t) = \int K(r, x; t) V(x) dx$$

with $K(r, x; t) = \int dt' G(r, x, t-t')$
 $\frac{\partial^2 G}{\partial t'^2}(x, s, t')$

⚠ $f(x, t) = \delta(x-s) \delta(t)$

⚠ $G(x, y, t) = a(x, y) \delta(t - \tau(x, y))$
 Simple geom optics



$$\Rightarrow K(r, x, t) = a(x, s) a(x, r) \delta''(t - \tau(x, r) - \tau(x, s))$$

and $\int \delta'' f = \int \delta f'' \quad \forall f$

$$u_{R,B}(r, t) = \frac{\partial^2}{\partial t^2} \int \underbrace{a(x, s) a(x, r)}_{\text{weight}} \underbrace{\delta(t - \tau(x, r) - \tau(x, s))}_{\text{integral over } \{x = t = \tau(x, r) + \tau(x, s)\}} V(x) dx$$

isochron (one for each r, t)

\Rightarrow generalized Radon transform along ellipse / ellipsoid when $c_0 = \text{const.}$

\Rightarrow Kirchhoff modeling. $\tau(x, y) = \frac{|x-y|}{c_0}$

Adjoint: see p. 6 prev. lecture
Kirchhoff migration.

Multiple sources: $u_{x,B}(r,s,t) = F_s(V(x))(r,t)$

Modeling: $FV = \{ F_s V ; s \text{ sources} \}$

$$\langle d, FV \rangle_{(r,t;s)} = \langle F^*d, V \rangle_{(x)}$$

$$\sum_{r,s,t} d(r,s,t) (FV)(r,t)$$

$$\sum_s \langle d(\cdot,s,\cdot), F_s V \rangle_{(r,t)}$$

$$\sum_s \langle F_s^* d, V \rangle_{(x)} = \langle \sum_s F_s^* d, V \rangle_{(x)}$$

$$\Rightarrow F^* = \sum_s F_s^*$$

Imaging operator: sum over ^{redundant coord.} sources:
= stack

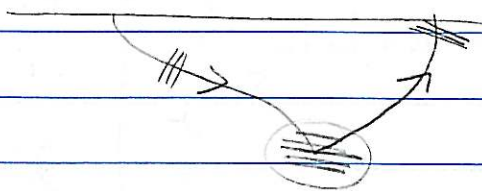
Other ways to cut up data space
gather eg common source (r,s)
Common midpoint $(\frac{r-s}{2}, \frac{r+s}{2})$
offset

$\{ F_s^* d \}$ prestack migration
 $\sum_s F_s^* d$ poststack migration

Algo for K mig:
loop over s
loop over x
loop over r
at $r \pm \dots$

Kinematics of migration (instead of studying the normal of like before)

$$u_{\alpha, \beta}(\sigma, t) = \int K(\sigma, x, t) V(x) dx$$



Goal: understand relationship between singular/oscillatory behaviours of $u_{\alpha, \beta}$ and V (microlocal)

(eg $\delta(x) = \int |e^{i\tilde{x}} d\tilde{x}$)
 \Rightarrow very rich high freq. content)

a) Singular supports: $V(x) = \delta(x - x_0)$

$$\Rightarrow u_{\alpha, \beta}(\sigma, t) = \text{amp} \times \delta(t - \tau(x_0, \sigma) - \tau(x_0, \lambda))$$

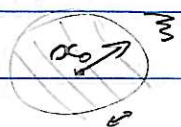
map $x_0 \rightarrow \{(\sigma, t) : t = \tau(x_0, \sigma) + \tau(x_0, \lambda)\} = \Sigma_{x_0}$

Then $\text{sing supp}(u_{\alpha, \beta}) \subseteq \bigcup_{x_0 \in \text{sing supp } V} \Sigma_{x_0}$

But = find a way to index $x_0 \in \text{sing supp } V \rightarrow e^{\tilde{x} \cdot x}$ no confusion with Hom. flow

b) (local) Oscillations = $V(x) = \chi(x - x_0) e^{i\tilde{x} \cdot x}$

$\lambda = \frac{2\pi}{|\tilde{x}|}$ tiny
 diam(x) small



(complex OK)

\rightarrow study oscillations in $u_{\alpha, \beta}$
 \rightarrow compare them to $\psi(\sigma, t) = e^{i(\omega t + \tilde{x} \cdot \sigma)}$ (detector)

on $\vec{\xi}, \vec{\xi}_n, \omega$

Find conditions under which

$\langle \bar{\Psi}, FV \rangle_{(x,t)}$ is large
 "

$\int_{x,t} e^{-i(\omega t + \vec{\xi}_n \cdot \vec{r})} \int_{\alpha} \text{amp}(-) \chi_1 \chi_2 e^{i\vec{\xi} \cdot \vec{r}} \delta''(t - \tau(\alpha, x) - \tau(\alpha, \Lambda)) dx d\alpha dt.$

or in other words find $\vec{\xi}, \vec{\xi}_n, \omega$ s.t. $K(\vec{\xi}, \vec{\xi}_n, \omega)$ is large (F.T. in all arguments)

int by parts in t: $\frac{\partial^2}{\partial t^2} (e^{-i\omega t} \chi_2(t-t_0, r-r_0))$
 $\approx -\omega^2 e^{-i\omega t} \chi_2(t-t_0, r-r_0)$

→ replace δ'' by δ

$\int_{\alpha, x} e^{i\vec{\xi} \cdot \vec{r}} e^{-i\vec{\xi}_n \cdot \vec{r}} e^{-i\omega(\tau(\alpha, x) + \tau(\alpha, \Lambda))} [\text{amp } \chi_1 \chi_2] dx d\alpha$

Stationary phase: (check nondegen Hessian: see homework!)

• in α : $\sum -i\omega \nabla_x \tau(\alpha, x) - i\omega \nabla_x \tau(\alpha, \Lambda) = 0$

$\vec{\xi} = i\omega (\nabla_x \tau(\alpha, x) + \nabla_x \tau(\alpha, \Lambda))$

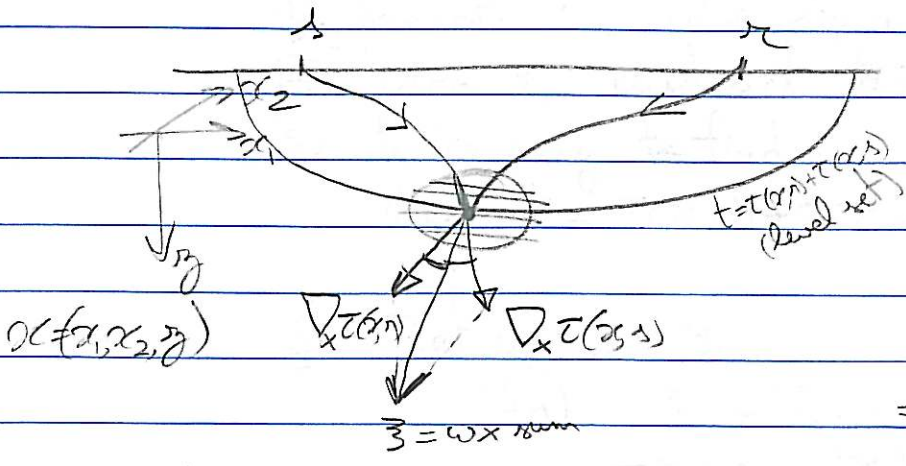
Δ chosen s.t. $\nabla_x \tau(\alpha, x) + \nabla_x \tau(\alpha, \Lambda) \neq 0$

• in \vec{r} : $\vec{\xi}_n + \omega \nabla_r \tau(\alpha, r) = 0$

$\eta \cdot (\quad) = 0$

$$\begin{pmatrix} \vec{\xi}_n \\ \omega \end{pmatrix} \cdot \begin{pmatrix} \eta \\ \eta \cdot \nabla_r \tau(\alpha, r) \end{pmatrix} = 0$$

• in x : (model space)



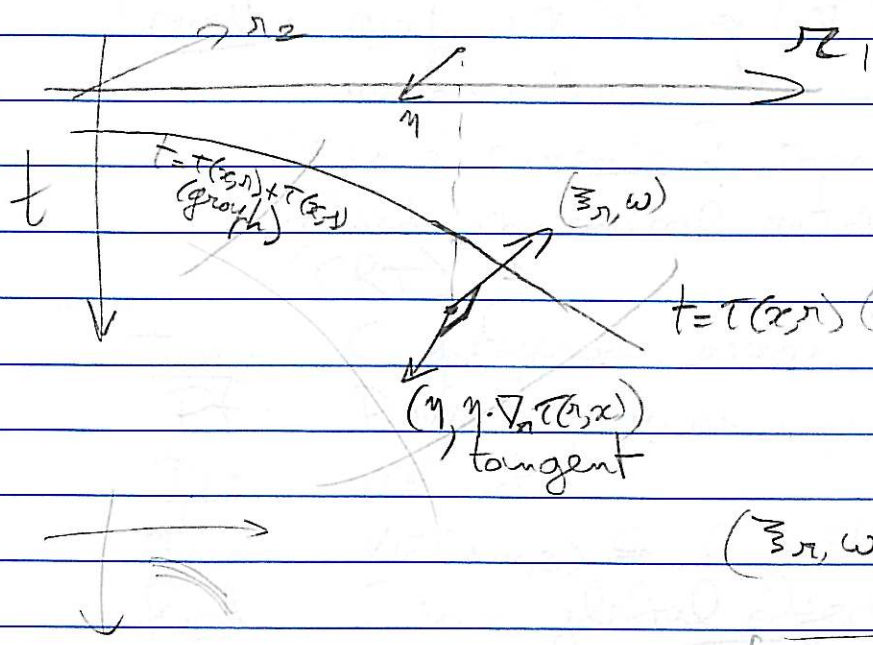
$\vec{z} = (\omega)$ normal to reflector

$$|\nabla_x T(x_1)| = |\nabla_x T(x_2)| = 1/c(x)$$

$\Rightarrow \vec{z}$ bisector of $\nabla_x T(x_1)$ and $\nabla_x T(x_2)$

Snell's law of reflection
Also, \vec{z} normal to isochron!
(again, $\nabla_x T + \nabla_x T \neq 0$)

• in η : (data space)



$(\vec{\zeta}_\eta, \omega) = (\omega)$ normal to oscillations in data

$t = T(x_1) + T(x_2)$ movement curve/surface

$(\vec{\zeta}_\eta, \omega) \perp$ to all tang. vectors.

$\Rightarrow (\vec{\zeta}_\eta, \omega) = (\omega)$ -normal to movement curve/surface

How to det. $(\vec{\zeta}_\eta, \omega)$?
Fix η . Consider $(x, \vec{\zeta})$.

- ① connect x_1 and x by a ray, find $\nabla_x \tau(x, x_1)$
- ② from ξ , construct take-off dir. $-\nabla_x \tau(x, x_1)$
- ③ trace ray to find r ($y=0$ crossing),
then $t = \tau(x, r) + \tau(x, x_1)$

④ Det ω and ξ_r :

④a) $\xi = \omega (\nabla_x \tau(x, r) + \nabla_x \tau(x, x_1))$

$\Rightarrow \omega = |\xi| / |\nabla_x \tau + \nabla_x \tau|$ (denom. $\neq 0$)

④b) $\xi_r + \omega \nabla_r \tau(x, r) = 0$

$\Rightarrow \xi_r = -\omega \nabla_r \tau(x, r)$

Concl = map $(x, \xi) \rightarrow (x, t; \xi_r, \omega) = C(x, \xi)$
removes the multi-valuedness of $x \rightarrow (x, t)$

It is a canonical transformation,
called canonical reflection transformation
(CRT)

Thm: C preserves areas (symplectomorphism)
(PF: Pullback preserves symplectic 2-form.)

Def: WF set

Thm $WF(x, \xi) \subseteq C \circ WF(V)$

Remark: C is a global diffeo - extends beyond caustics, when multipathing

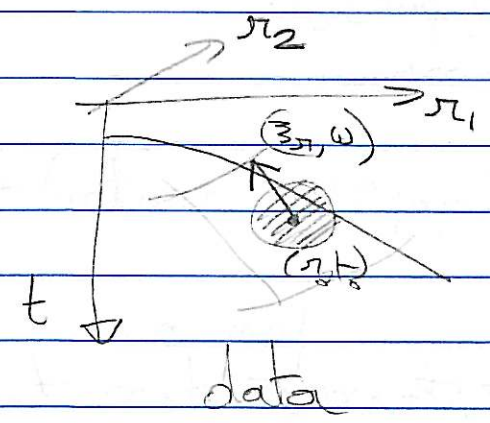
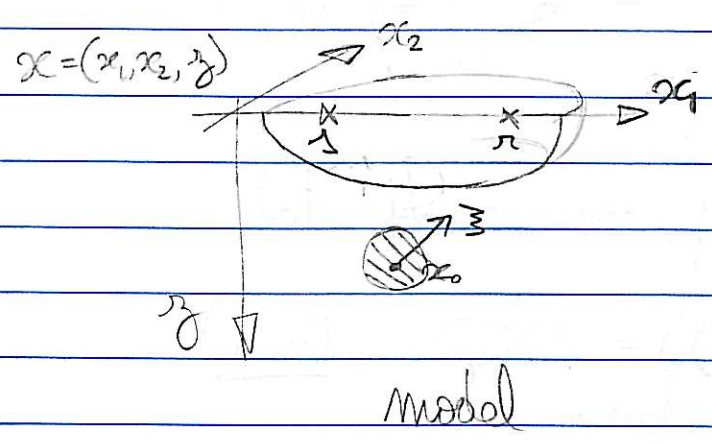
W/19/09 Kinematics.

Forward / demigration:

$$FV(x,t) = \int a(x,r) a(x,s) \delta''(t - \tau(x,r) - \tau(x,s)) V(x) dx$$

Migration:

$$F^* d(x) = \int a(x,r) a(x,s) \delta''(t - \tau(x,r) - \tau(x,s)) d(r,t) dr dt$$



Local oscillations = $(x_0, \vec{z}) \rightarrow \left(\begin{pmatrix} r_0 \\ t_0 \end{pmatrix}, \begin{pmatrix} \vec{z}_r \\ \omega \end{pmatrix} \right)$
 (map in phase-space)

trial. $V(x) = \chi(x-x_0) e^{i\vec{z} \cdot x}$, $\chi \in C^\infty$
 test $\Psi(r,t) = \chi(r-r_0) \chi(t-t_0) e^{i(\vec{z}_r \cdot r + \omega t)}$

then $\langle \Psi, FV \rangle_{(r,t)} = \iint \overline{\Psi(r,t)} FV(r,t) dr dt$
 $\hookrightarrow \Delta$ continuous r, t .

Rank: V complex = OK by linearity.

Assume $\frac{2\pi}{|\vec{z}|}, \frac{2\pi}{|\omega|} \ll \text{diam}(\text{supp}(\chi))$

$$\langle \psi, FV \rangle_{(n,t)} = \iint e^{-i(\frac{3}{2}\pi + \omega t)} e^{i3x} \chi(x-x_0) \chi(n-x_0) \chi(t-t_0) \\ \times [\text{amp}] \delta''(t - \tau(x,n) - \tau(x,s)) dx ds dt$$

Int by parts = $\frac{\partial^2}{\partial t^2} (e^{-i\omega t} \chi(t-t_0))$
 $= -\omega^2 e^{-i\omega t} \chi(t-t_0)$

$$\langle \psi, FV \rangle = \int e^{i(3x - \frac{3}{2}\pi - \omega(\tau(x,n) + \tau(x,s)))} [\chi] [\text{amp}] dx ds dt$$

Stat. phase. see p 4 last lecture. (2 pages here)

Concl. $C_F(x, \vec{z}) = \left(\begin{pmatrix} x \\ t \end{pmatrix}, \begin{pmatrix} \vec{z} \\ \omega \end{pmatrix} \right)$ is a map.
 - not defined at x st. $\nabla_x \tau(x,n) + \nabla_x \tau(x,s) = 0$
 - also,
 - ray from s may not reach x
 - $\xrightarrow{\quad} x \xleftarrow{\quad} s$
 (go off to ∞)

Called cononical reflection transformation
 (Surjectivity: perhaps not either)

Remark, Global object: generalizes to multivalued travel times

Remark, (Hw) For $\left(\begin{pmatrix} x \\ t \end{pmatrix}, \begin{pmatrix} \vec{z} \\ \omega \end{pmatrix} \right) \in \text{Ran } C_F$,
 (tough)

$$C_{F^*} \left(\begin{pmatrix} x \\ t \end{pmatrix}, \begin{pmatrix} \vec{z} \\ \omega \end{pmatrix} \right) = (x, \vec{z}) \text{ as above}$$

$$\rightarrow C_{F^*} = C_F^{-1}$$

\rightarrow resolution of ∞ in inv. prob, high freq.

Real definition of C_F : uses wavefront sets (microlocal analysis)

Def. (later) $WF(V) = \{ (x, \xi) \in \mathbb{R}^{2m} :$

" V is singular at x in the direction ξ " $\}$



Prop. Sing supp $V = \Pi WF(V) = \{ x = \exists \xi : (x, \xi) \in WF(V) \}$

Def. $C_F = WF(K(\sigma, t; x))$

\hookrightarrow kernel of F .

$= \{ (x, \xi; \begin{pmatrix} \sigma \\ t \end{pmatrix}, \begin{pmatrix} \xi \\ \omega \end{pmatrix}) : "K \text{ is singular at } (\begin{pmatrix} \sigma \\ t \end{pmatrix}; x) \text{ in the direction } (\begin{pmatrix} \xi \\ \omega \end{pmatrix}, \xi)" \}$

Def. Composition of WF. / relations

$A \subseteq X, B \subseteq X \times Y$

$B \circ A = \{ y \in Y : \exists x \in A : (x, y) \in B \}$

Thm. (Composition of distributions)

Assume C_F does not have elements of the kind $(x, \xi; 0, 0)$ or $(\infty, \begin{pmatrix} \xi \\ \omega \end{pmatrix})$

amount to $\partial_x^\alpha + \partial_x^\beta \neq 0$

Then $WF(FV) \subseteq \underbrace{WF(K)}_F \circ WF(V)$

(see Hörmander Vol 1)

Cor. $WF(F^*FV) \subseteq WF(V)$

(normal operator is microlocal)

Def. $WF(f)$

Fix x ; Def. the singular fiber $\Sigma_x(f)$ as follows:

$\nu \in (\Sigma_x(f))^c$ when there exist:

- $\phi \in C_0^\infty = \phi(x) \neq 0$,
- Γ a conical neighborhood of ν ,

such that

$$|\widehat{\phi f}(\xi)| \leq C_N (1+|\xi|)^{-N} \quad \forall \xi \in \Gamma$$

$$\forall N \in \mathbb{Z}_0^+$$

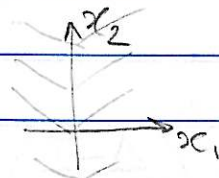
Then $WF(f) = \{(x, \xi) : \xi \in \Sigma_x(f)\}$.

Ex. $f \in C_0^\infty$; $WF(f) = \emptyset$.

$f(x) = \delta(x)$; $WF(f) = \{(0, \xi), \xi \in \mathbb{R}^n\}$.

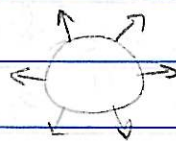
$f(x) = |x-a|$; $WF(f) = \{(a, \xi), \xi \in \mathbb{R}^n\}$

$f(x) = |x|$;
($x \in \mathbb{R}^2$)



$WF(f) = \left\{ \left(\begin{pmatrix} 0 \\ x_2 \end{pmatrix}, \begin{pmatrix} \xi_1 \\ 0 \end{pmatrix} \right), \right.$
 $\left. x_2, \xi_1 \in \mathbb{R} \right\}$

etc.



Next: Normal operator.

Rank: Canonical = symplectic = preserves areas

12/01/09 Microlocal analysis: formalize ideas such as

Map $(\alpha, \xi) \rightarrow \left(\begin{pmatrix} \alpha \\ \xi \end{pmatrix}, \begin{pmatrix} \alpha \\ \xi \end{pmatrix} \right)$
 (model) (data)

- + reflection, reflection
- + mapping of singularities
- + like the identity
- + directions of singularities
- + imaging vs. inversion.

obtained by stationary phase.

→ Fourier analysis of singularities
 Wavefront sets.

$k \mapsto \hat{u}$ Thm Let $u \in C^\infty(\mathbb{R}^n)$, and $\hat{u}(k) = \int e^{-ix \cdot k} u(x) dx$
 Then $|\hat{u}(k)| \leq C_N (1+|k|)^{-N}$ $N=1, 2, 3, \dots$

PF $\left[\begin{array}{l} k = \beta \hat{k}, \quad \beta = |k|, \quad \hat{k} = k/|k| \\ x \cdot k = \beta \phi(x) \quad \text{with} \quad \phi(x) = \hat{k} \cdot x \\ \nabla \phi = \hat{k} \neq 0 \\ \text{Stat. phase: } |\hat{u}(k)| \leq C_N \beta^{-N} \end{array} \right]$

$$\frac{(I - \Delta)}{1+|k|^2} e^{-ix \cdot k} = e^{-ix \cdot k}$$

$$\begin{aligned} \hat{u}(k) &= \int \left(\frac{I - \Delta}{1+|k|^2} \right)^N \left[e^{-ix \cdot k} \right] u(x) dx \\ &= (1+|k|^2)^{-N} \int e^{-ix \cdot k} (I - \Delta)^N u(x) dx \end{aligned}$$

$$|\hat{u}(k)| \leq (1+|k|^2)^{-N} \cdot \underbrace{\| (I - \Delta)^N u \|_{L^1}}_{\|u\|_{W^{2N}}} \quad \square$$

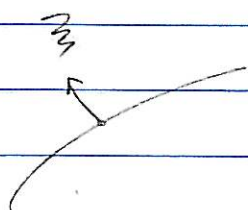
$$\|u\|_{W^{2N}} = C_N$$

Def. Let $u \in L^1(\mathbb{R}^m)$ or $\mathcal{D}'(\mathbb{R}^m)$

$$\text{Sing supp}(u) = \left\{ x \in \mathbb{R}^m : \nexists \text{ open neighborhood } V \text{ of } x \text{ s.t. } u \in C^\infty(V) \right\}$$

New $\text{WF}(u)$: set of couples $(x, \xi) \in \mathbb{R}^m \times \mathbb{R}^m \setminus \{0\}$

- $x \in \text{sing supp } u$.
- " u is singular in the direction ξ " (estimate $(1+|\xi|^2)^{-N}$ fails.)



Def. (conic neighborhood)

X is a conic neighborhood of $Y \subseteq \mathbb{R}^m$

- if
- X is open
 - $X \supseteq Y$
 - if $\xi \in X$ then $\lambda \xi \in X$ for all $\lambda > 0$.

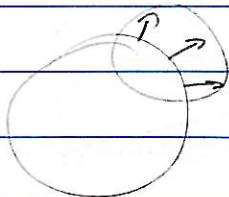
Def. (Singular cone).

$$\Gamma'(u) = \left\{ \eta \in \mathbb{R}^m \setminus \{0\} : \nexists \text{ conic neighborhood } W \text{ of } \eta \text{ s.t. for } \xi \in W, \right.$$

$$\left| \hat{u}(\xi) \right| \leq C_N (1+|\xi|^2)^{-N} \quad N=1,2,3,\dots \left. \right\}$$

= "directions in which u is nonsmooth, globally"

Lemma Let $\phi \in C_0^\infty(\mathbb{R}^m)$ and $u \in L^1(\mathbb{R}^m)$ or $\mathcal{D}'(\mathbb{R}^m)$.
Then $\Gamma'(\phi u) \subseteq \Gamma'(u)$.



Def (Singular fiber) For $x \in X \subseteq \mathbb{R}^n$,

$$\Gamma_x(u) = \bigcap_{\phi} \Gamma(\phi u); \quad \phi \in C_0^\infty(X), \phi(x) \neq 0.$$

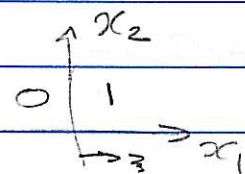
Alt. def, $\Gamma_x(u) = \lim_{j \rightarrow \infty} \Gamma(\phi_j u)$ $\phi_j \in C_0^\infty(X), \phi_j(x) \neq 0,$
 $\text{supp } \phi_j \rightarrow \{x\}.$

Def (Wavefront set)

$$WF(u) = \{ (x, \xi) \in \mathbb{R}^n \times \mathbb{R}^n \setminus \{0\} \text{ s.t. } \xi \in \Gamma_x(u) \}$$

Prop $\Pi_{\mathbb{R}^n} WF(u) = \text{sing supp}(u)$
 $\Pi_{\mathbb{R}^n \setminus \{0\}} WF(u) = \Gamma(u)$

Ex. $u = H(x_1) = \begin{cases} 1 & \text{if } x_1 > 0 \\ 0 & \text{if } x_1 \leq 0 \end{cases}$



$$WF(u) = \left\{ \left(\begin{pmatrix} 0 \\ x_2 \end{pmatrix}, \begin{pmatrix} \xi_1 \\ 0 \end{pmatrix} \right) : x_2 \in \mathbb{R}, \xi_1 \in \mathbb{R} \setminus \{0\} \right\}$$

Because: $\phi u \in C_0^\infty$ when $\text{supp } \phi \cap \{x_1 = 0\} = \emptyset$

$$\Rightarrow \text{sing supp} = \left\{ \begin{pmatrix} 0 \\ x_2 \end{pmatrix} : x_2 \in \mathbb{R} \right\}$$

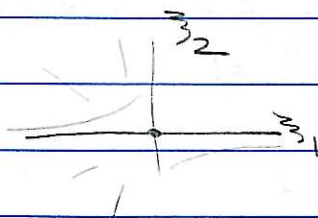
$$\hat{u}(\xi) = \int e^{-ix \cdot \xi} H(x_1) dx_1 dx_2$$

$$= \frac{1}{i\xi_1} \delta(\xi_2)$$

$$\Rightarrow \Gamma(u) = \left\{ \begin{pmatrix} \xi_1 \\ 0 \end{pmatrix} : \xi_1 \in \mathbb{R} \setminus \{0\} \right\}$$

$$\hat{\phi} u(\xi) = \hat{\phi} \star \hat{u}(\xi)$$

$$\Rightarrow \Gamma_x(u) \subseteq \Gamma(u)$$



Ex. $u = \delta(x) = \delta(x_1) \delta(x_2)$ $x \in \mathbb{R}^2$

$WF(u) = \left\{ \left(\begin{pmatrix} 0 \\ 0 \end{pmatrix}, \begin{pmatrix} \lambda_1 \\ \lambda_2 \end{pmatrix} \right) : \lambda_1, \lambda_2 \in \mathbb{R} \setminus \{0\} \right\}$

Because \cdot $\phi u = 0$ when $\phi(0) = 0$
 $\Rightarrow \text{sing supp}(u) = \{0\}$

$\cdot \hat{u}(\xi) = 1$
 $\Rightarrow \Gamma(u) = \mathbb{R}^2 \setminus \{0\}$

$\cdot \phi u(x) = \phi(0) \delta(x)$
 $\hat{\phi u}(\xi) = \phi(0)$
 $\Rightarrow \Gamma_x(u) = \mathbb{R}^2 \setminus \{0\}$ when $\phi(0) \neq 0$

Ex. $u = H(\psi(x))$ $\psi = \text{level set function, assume } \nabla \psi \neq 0$
 $\xi = \lambda \nabla \psi(x)$ $H = \text{Heaviside}$



$WF(u) = \left\{ (x, \lambda \nabla \psi(x)) : \psi(x) = 0, \lambda \neq 0 \right\}$

Break? (Pf. deform to $H(x_1)$.)

Concl. $(x_0, \xi_0) \in WF(u)$ if $(e^{ix \cdot \xi_0}, \phi_{x_0} u)$

is "large" when $\cdot \phi_{x_0}$ is loc. near x_0
 $\cdot \xi_0$ is very large
 $\text{rot}(1 + |\xi_0|^2)^{-1} \notin \mathbb{N}$



Remark. $FV(\eta, t) = \int K(x, \eta, t) V(x) dx$

check $\langle \phi_{(\frac{x_0}{t_0})} e^{i(\eta \beta_1 + \omega t)}, K \phi_{x_0} e^{i x \beta_1} \rangle$

$= \langle \phi_{x_0} \phi_{(\frac{x_0}{t_0})} e^{i(x \beta_1 + \eta \beta_1 + \omega t)}, K \rangle$

$= \int_{(x, \eta, t)} (\phi_{x_0} \phi_{(\frac{x_0}{t_0})} K) (\frac{\beta_1}{\omega}, \frac{\beta_1}{\omega}, \omega)$

Def. $C_F =$ "set of $(x, (\frac{\eta}{t}); \frac{\beta_1}{\omega} (\frac{\beta_1}{\omega}))$ s.t. the above is large"

$= WF(K)$

$=$ "set of $(\frac{\beta_1}{\omega}, \frac{\beta_1}{\omega})$ s.t.

K is singular at $(x, (\frac{\eta}{t}))$ in the direction $(\frac{\beta_1}{\omega}, \frac{\beta_1}{\omega})$ "

Remark. Not a map, but a relation

Def. Composition of relations

$A \subseteq X, B \subseteq X \times Y$

$B \circ A = \{ y \in Y : \exists x \in A : (x, y) \in B \}$

Thm. (Hörmander) Assume the elements of C_F obey $\beta_1 = 0 \Leftrightarrow (\frac{\beta_1}{\omega}) = 0$. Then

$WF(FV) \subseteq WF(K) \circ WF(V)$

(using kernel K)

\circ_F

↓
reflections

↓
canonical reflection int.

↓
reflectors

Rank. Simple geom optics \Rightarrow $C_{F^*} = C_F^{-1}$ on $\text{Ran } C_F$,
 when $F =$ Kirchhoff modeling
 $F^* =$ — migration
 Then

$$\begin{aligned} \text{WF}(F^*F V) &\subseteq \text{WF}(F^*) \circ \text{WF}(FV) \\ &\subseteq \text{WF}(F^*) \circ \text{WF}(F) \circ \text{WF}(V) \\ &\subseteq \text{WF}(V) \end{aligned}$$

\Rightarrow The normal operator F^*F does not move singularities. "like the identity".

Def. An operator T s.t. $\text{WF}(Tu) \subseteq \text{WF}(u)$ is called a microlocal operator

Rank. The normal operator needs not be invertible
 Inclusion but not identity of WF sets:
 Imaging but not (accurate) inversion
 Idealized ^{linearized} migration problem: given d ,
 find $\tilde{V} : \text{WF}(\tilde{V}) \subseteq \text{WF}(V)$.

Next: | Diff. op. are microlocal.
 Ψ DO are microlocal
 The normal op. is —
 Change of variables for WF sets. \rightarrow atang. bundle
 Can. rel. preserve areas \rightarrow symplectic

12/03/09

Higl - frequency imaging
! Simple geom. optics

$$WF(d) = WF(FV) \subseteq WF(F) \circ WF(V)$$

$$WF(F^*d) = WF(F^*FV) \subseteq WF(F^*) \circ WF(F) \circ WF(V)$$

$$\subseteq WF(V)$$

$V \xrightarrow{F} d$
 $WF(F) \downarrow$
 $WF(V) \xrightarrow{WF(F)} WF(d)$
 $WF(F^*) \leftarrow$

Q. $WF(F^*FV) \subseteq WF(V)$.
(conditions under which)

Can write kernel of F^*F as (change var.)

$$N(x,y) = \int e^{i(x-y) \cdot \xi} a(x, \xi) d\xi \approx \delta(x-y)$$

\hookrightarrow smooth ≈ 1
 (singular at $x=y$)

$$F^*F V(x) = \int N(x,y) V(y) dy$$

$$= \int \int e^{i(x-y) \cdot \xi} a(x, \xi) V(y) dy d\xi$$

$$= \int e^{ix \cdot \xi} a(x, \xi) \hat{V}(\xi) d\xi \approx V(x)$$

\rightarrow called a pseudodifferential operator.

Ex. $a(x, \xi) = (2\pi)^{-m} \sum_{|\alpha| \leq N} a_\alpha(x) (i\xi)^\alpha$ $\alpha = (\alpha_1, \dots, \alpha_m)$

(polynomial of $i\xi$)

$$(i\xi)^\alpha = (i\xi_1)^{\alpha_1} \dots (i\xi_m)^{\alpha_m}$$

$$|\alpha| = \alpha_1 + \dots + \alpha_m$$

$N =$ degree of the polynomial

if Hörmander applies and if $C_{F^*} = C_F^{-1}$

then
$$\int e^{i\alpha \cdot \xi} (2\pi)^{-m} \sum_{\alpha} a_{\alpha}(\alpha) (i\xi)^{\alpha} \widehat{V}(\xi) d\xi$$

$$= \sum_{\alpha} a_{\alpha}(\alpha) (2\pi)^{-m} \int e^{i\alpha \cdot \xi} (i\xi)^{\alpha} \widehat{V}(\xi) d\xi$$

$$= \sum_{\alpha} a_{\alpha}(\alpha) \nabla^{\alpha} V(x)$$

where $\nabla^{\alpha} = \left(\frac{\partial}{\partial x_1}\right)^{\alpha_1} \dots \left(\frac{\partial}{\partial x_m}\right)^{\alpha_m}$

→ differential operator.

Thm If $a(\alpha, \xi) = \sum_{\alpha} a_{\alpha}(\alpha) (i\xi)^{\alpha}$, $a_{\alpha}(\alpha) \in C_0^{\infty}(\mathbb{R}^m)$

and $Tu(x) = \int e^{i\alpha \cdot \xi} a(\alpha, \xi) \widehat{u}(\xi) d\xi$
then $WF(Tu) \subseteq WF(u)$, when $u \in C_0^{\infty}(\mathbb{R}^m)$

Rem $\Gamma(u) =$ "directions of slow decay of $\widehat{u}(\xi)$ "
 $\Gamma_{\infty}(u) = \lim_{j \rightarrow \infty} \Gamma(\phi_j u)$, wff $\phi_j \rightarrow \{x\}$

$$WF(u) = \{ (x, \xi) : \xi \in \Gamma_{\infty}(u) \}$$

PF of thm Check $|\xi| \rightarrow \infty$ in

$$T_j(\xi) = \int e^{-i\alpha \cdot \xi} \phi_j(x) Tu(x) dx$$

$$= \int e^{-i\alpha \cdot \xi} \phi_j(x) \sum_{\alpha} a_{\alpha}(\alpha) \nabla^{\alpha} u(x) dx$$

$$= \sum_{\alpha} \int u(x) (-1)^{|\alpha|} \nabla^{\alpha} \left(e^{-i\alpha \cdot \xi} \phi_j(x) a_{\alpha}(x) \right) dx$$

$$= \sum_{\alpha} \int u(x) \sum_{\beta \leq \alpha} \beta_{\alpha}(\alpha) e^{-i\alpha \cdot \xi} \dots dx$$

where $\beta_{\alpha j}(x) = \text{li. co. of } \phi_j \text{ and } a_\alpha$
and their derivatives
 $\in C^\infty(\mathbb{R}^n)$.

$$\text{supp } \beta_{\alpha j}(x) \subseteq \text{supp } \phi_j(x)$$
$$= \sum_{\alpha} \sum_{j} \int e^{-ix \cdot \xi} \left[\sum_{\alpha} \beta_{\alpha j}(x) \right] u(x) dx.$$

Fix $x \in \text{supp}(u)$, let $\xi \notin \Gamma_x(u)$.

$$\text{then } \exists j > 0 : \left| \int e^{-ix \cdot \xi} \left[\sum_{\alpha} \beta_{\alpha j}(x) \right] u(x) dx \right| \leq C_M (1+|\xi|^2)^{-M} \quad \forall M$$

$$\Rightarrow \left| \sum_{|\alpha| \leq N} \sum_{j} \int (\quad) \right|$$

$$\leq \sum_{|\alpha| \leq N} C_M \sum_{j} \xi^\alpha (1+|\xi|^2)^{-M} \quad \forall M$$

$$\leq C'_M (1+|\xi|^2)^{-M+|\alpha|} \quad \forall M$$

$$\leq C'_{M+1/2} (1+|\xi|^2)^{-M}$$

$$\Rightarrow \xi \notin \Gamma_x(Tu)$$

Concl : $(x, \xi) \notin \text{WF}(u) \Rightarrow (x, \xi) \notin \text{WF}(Tu)$

$$(x, \xi) \in \text{WF}(Tu) \Rightarrow (x, \xi) \in \text{WF}(u)$$

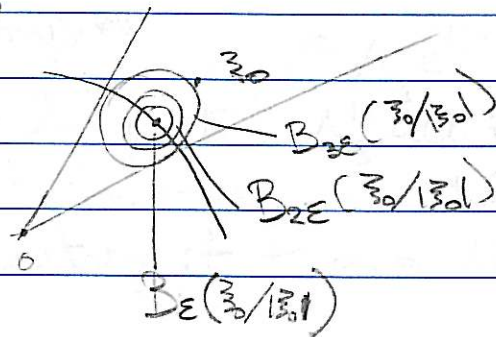
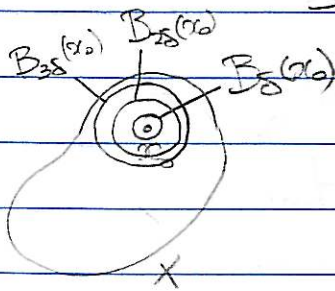
$$\text{WF}(Tu) \subseteq \text{WF}(u)$$

□.

Thm Let $Tu(x) = \int e^{ix \cdot \xi} a(x, \xi) \hat{u}(\xi) d\xi$
 with $|\nabla_{\xi}^{\alpha} \nabla_x^{\beta} a(x, \xi)| \leq C_{\alpha\beta} (1 + |\xi|^2)^{(m-k)/2}$
 and a compactly supported in x .
 Then $\forall u \in C_0^{\infty}(\mathbb{R}^m)$
 $WF(Tu) \subseteq WF(u)$

Ex. $T = F^*F$

Pf. Assume $(x_0, \xi_0) \notin WF(u)$, show $(x_0, \xi_0) \notin WF(Tu)$
 $\Rightarrow \exists X$ neigh. of x_0 ,
 $\exists W$ conic neigh. of ξ_0 ,



$\hat{\phi}u(\xi)$ decays fast
 for $\text{supp } \phi \subseteq X$
 $\xi \in W$

Let $\phi \in C_0^{\infty}(B_{\delta}(x_0))$, $\phi(x_0) \neq 0$.
 $\psi \in C_0^{\infty}(B_{3\delta}(x_0))$, $\psi(x) = 1$, $x \in B_{\delta}(x_0)$

Estimate $\widehat{\phi Tu}(\xi)$, $\xi \in B_{\epsilon}(\xi_0/|\xi_0|)$

Write $u = (1-\psi)u + \psi u$

Claims: 1) $\widehat{\phi T(1-\psi)u} \in C^{\infty}$

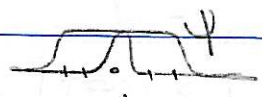
2) $\widehat{\phi T\psi u}(\xi)$ decays fast.

for $\xi \in B_{\delta}(\frac{\xi_0}{|\xi_0|})$

1) $\mathcal{F}T(1-\psi)u(x)$

$$= \phi(x) \int e^{ix\xi} a(x, \xi) \int e^{-iy\xi} (1-\psi(y)) u(y) dy d\xi$$

$$= \int d\xi \int dy a(x, \xi) \phi(x) (1-\psi(y)) e^{i(x-y)\xi} u(y)$$



$|x-y| > \delta$ for $x \in \text{supp } \phi$,
 $y \in \text{supp } (1-\psi)$

$$e^{i(x-y)\xi} = |x-y|^{-2} \Delta_{\xi} e^{i(x-y)\xi} \quad |x-y|^{-2} \text{ handles}$$

int by parts: $|\Delta_{\xi}^M a(x, \xi)| \leq (1+|\xi|^2)^{-M}$

→ int. converges with M large enough.

Some story if \mathcal{P} is a diff. op.

$\mathcal{P}\mathcal{F}T(1-\psi)u$ introduces powers of ξ
in the integrand

→ undone by $(1+|\xi|^2)^{-M}$

2). $\widehat{\mathcal{F}T\psi u}(\eta) = \iint a(x, \xi) \phi(x) \widehat{\psi u}(\xi) e^{ix \cdot (\xi - \eta)} dx d\xi$

$\eta = \tau\theta'$, $\theta' \in B_{\varepsilon}(\xi_0/|\xi_0|)$

$\xi = \tau\theta$

$$\int a(x, \tau\theta) \phi(x) \widehat{\psi u}(\tau\theta) e^{i\tau x \cdot (\theta - \theta')} dx$$

$\tau^m d\theta$

• Ques $|\theta - \theta'| > \varepsilon$:

$$\tau^2(\theta - \theta')^2 (-\Delta_x) e^{i\tau x \cdot (\theta - \theta')} = e^{i\tau x \cdot (\theta - \theta')}$$

Int. by parts: get τ^{-2M}
 Δ_x on a and ϕ is fine
 \rightarrow fast decay in τ

• Over $|\theta - \theta'| \leq \varepsilon$, $\theta \in B_{2\varepsilon}(\xi_0/|\xi_0|) \subseteq W$

Since $X \times W$ is disjoint from $V \setminus F(\mu)$,
 $|\widehat{\Psi \mu}(\tau \theta)| \leq C_N \tau^{-N}$
 \rightarrow fast decay in τ □

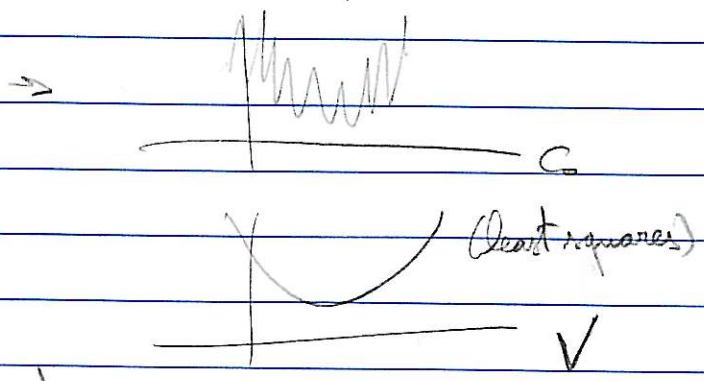
Open problems:

① Compute F^*F from $a(x, \xi)$
 and $(F^*F + S)^{-1} u(x) = \int e^{ix \cdot \xi} b(x, \xi) \widehat{u}(\xi) d\xi$

\rightarrow get b from a numerically, in compressed form

② Obtain c_0 , model velocity estimation.

$\min_{c_0, V} \|d - F[c_0]V\|_2$ prestack



\rightarrow new objectives

② Compressed sensing

$$f_i(x, t) \rightarrow F_j(x, t) = \sum_j a_{ij} f_i(x, t)$$

$$d(r, s, t) \rightarrow D_j(r, t) = \sum_j a_{ij} d(r, s, t)$$

Recover V from D_j - not more stack.
 Use sparsity ideas:

$$\min_R \sum_k |\langle V, \psi_k \rangle| \text{ s.t. predicted} = \text{actual}$$