# Lecture 22: Lévy Distributions

Scribe: Neville E. Sanjana

Department of Brain and Cognitive Sciences, MIT

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#### Central Limit Theorem (from previously) 1

Let  $\{x_i\}$  be IID random variables with finite mean  $\mu$  and finite variance  $\sigma^2$ . Then the RV

$$S_n = \frac{1}{\sigma\sqrt{n}} \left( \sum_{j=1}^n x_j - n\mu \right) \tag{1}$$

converges to a standard N(0,1) Gaussian density.

#### $\mathbf{2}$ Introduction

This lectures covers similar material to Lectures 12 and 13 in the 2003 lecture notes. Also, see Hughes  $\S 4.3$ .

Motivating question: What are the limiting distributions for a sum of IID variables? (And, in particular, those who violate the assumptions of the CLT.)

### Lévy stability laws $\mathbf{3}$

Definitions:  $\{x_i\}$  RVs and  $X_N = \sum_{i=1}^{N} x_i$ 

Characteristic function:

$$\hat{p}_n(k) = \left[\hat{p}(k)\right]^n \tag{2}$$

$$\hat{p}_n(k) = [\hat{p}(k)]^n$$

$$\hat{p}_{n \times m}(k) = [\hat{p}_m(k)]^n = [\hat{p}(k)]^{n \times m}$$
(2)
(3)

Introduce new variables  $Z_N = \frac{X_N}{a_N}$ . The PDF of the  $Z_N$ 's is given by  $F_N(x) = p_N \frac{x}{a_N} \frac{1}{a_N}$ .

$$\hat{F}_{N}\left(a_{N}k\right) = \hat{p}_{N}\left(k\right) \tag{4}$$

Substituting in,

$$\hat{F}_{n \times m} \left( a_{nm} k \right) = \left[ \hat{F}_m \left( a_m k \right) \right]^n \tag{5}$$

<sup>&</sup>lt;sup>1</sup>B. Hughes, Random Walks and Random Environments, Vol. I, Sec. 2.1 (Oxford, 1995).

As  $n \to \infty$ , we want  $\hat{F}_n(k)$  to tend to some fixed distribution  $\hat{F}(k)$ . Say that  $\lim_{m \to \infty} \frac{a_{nm}}{a_m} = c_n$ ,

$$\hat{F}(kc_n) = \left[\hat{F}(k)\right]^n \tag{6}$$

Suppose I want to find

$$\hat{F}(k\mu(\lambda)) = \left[\hat{F}(k)\right]^{\lambda} \tag{7}$$

Introduce  $\psi(k) = \log \hat{F}(k)$ 

$$\psi\left(k\mu\left(\lambda\right)\right) = \lambda\psi\left(k\right) \tag{8}$$

Differentiate:

$$k\mu'(\lambda)\psi'(k\mu(\lambda)) = \psi(\lambda)$$
 (9)

$$\mu \left( \lambda = 1 \right) = 1 \tag{10}$$

$$k\mu'(1)\psi'(k) = \psi(k) \tag{11}$$

$$k\mu'(1)\psi'(k) = \psi(k)$$

$$\frac{d\psi}{dk} = \frac{\psi}{\mu'(1)k}$$
(11)

This has a solution:

$$\psi(k) = \begin{cases} v_1 |k|^{\alpha} & k > 0 \\ v_2 |k|^{\alpha} & k < 0 \end{cases}$$

$$(13)$$

$$\hat{F}(k) = \begin{cases} \exp(v_1 |k|^{\alpha}) & k > 0 \\ \exp(v_2 |k|^{\alpha}) & k < 0 \end{cases}$$
(14)

$$v_1, v_2 \in \mathcal{C} \tag{15}$$

Thus we have this condition:

$$\hat{F}(-k) = \hat{F}(k)^* \tag{16}$$

$$\Longrightarrow \hat{F}(k) = \exp\left(\left(c_1 + ic_2 sgn\left(k\right)\right) |k|^{\alpha}\right) \tag{17}$$

$$c_1, c_2 \in \Re \tag{18}$$

We can also write:

$$\hat{F}(k) = \exp\left[-a|k|^{\alpha} \left(1 - i\beta \tan \frac{\alpha \pi}{2} sgn(k)\right)\right]$$
(19)

In this equation,  $\beta$  relates to the skewness of the distribution.

### Lévy distribution 4

For above  $\hat{F}$  when  $\beta = 0$  (ie. symmetric limiting distribution), you get the Lévy distribution. It is defined in terms of its Fourier transform.

$$\hat{L}_{\alpha}(a,k) = \hat{F}(k) = \exp(-a|k|^{\alpha})$$
(20)

Two special cases of the the Lévy distribution are  $\alpha = 1, 2$ .

$$\alpha = 2$$
  $\hat{F}(k) = \exp(-ak^2)$  like Gaussian but more general (21)

$$\alpha = 1$$
  $\hat{F}(k) = \exp(-a|k|)$  Cauchy (22)

(25)

## 4.1 Large x expansions of the Lévy distribution

$$L_{\alpha}(a,x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \exp ikx - a |k|^{\alpha} dk$$
 (23)

$$= \frac{1}{\pi} \int_{0}^{\infty} \cos kx \exp -ak^{\alpha} dk \tag{24}$$

Integrate by parts:

 $=\frac{a\alpha}{\pi}\int\limits_{0}^{\infty}\sin\frac{k|x|}{|x|}\exp-ak^{\alpha}k^{\alpha-1}\mathrm{d}k \quad =\frac{a\alpha}{\pi\left|x\right|^{1+\alpha}}\int\limits_{0}^{\infty}\xi^{\alpha-1}\sin\xi\exp\frac{-a\xi^{\alpha}}{\left|x\right|^{\alpha}}\mathrm{d}\xi\left(26\right)$ 

Where we have made the following substitutions:

$$\xi = k |x|, d\xi = dk |x| \tag{27}$$

Taking the limit  $x \to \infty$ :

$$\sim_{x \to \infty} \frac{a\alpha}{\pi |x|^{1+\alpha}} \int_{0}^{\infty} \xi^{\alpha-1} \sin \xi d\xi$$
 (28)

$$\sim \frac{a\alpha\Gamma(\alpha)\sin\frac{\pi\alpha}{2}}{\pi|x|^{1+\alpha}} \quad \text{power-law tail}$$
 (29)

# 4.2 Small x expansion

$$L_{\alpha}(a,x) = \frac{1}{\pi} \int_{0}^{\infty} \cos kx \exp{-ak^{\alpha}} dk$$
 (30)

Taylor expand cosine: 
$$= \frac{1}{\pi} \int_{0}^{\infty} \sum_{n=0}^{\infty} (-1)^n \frac{(kx)^{2n}}{(2n)!} \exp{-ak^{\alpha} dk}$$
 (31)

Substitute:  $w = ak^{\alpha}$   $dw = \alpha ak^{\alpha-1}dk$  (32)

$$= \frac{1}{\pi} \sum_{n=0}^{\infty} \int_{0}^{\infty} (-1)^n \frac{(x)^{2n}}{(2n)!} \frac{w^{2n/\alpha}}{a} \exp -w \frac{\mathrm{d}w}{\alpha a \frac{w}{a} \frac{\alpha-1}{\alpha}}$$
(33)

We're only interested in evaluating this equation as  $x \to 0$ :

$$= \frac{1}{\pi} \sum_{n=0}^{\infty} \frac{(-1)^n}{(2n)!} \frac{(x)^{2n}}{a^{\frac{2n+1}{\alpha}}} \Gamma\left(\frac{2n+1}{\alpha}\right)$$
 (34)

$$\sim \frac{1}{\pi} \frac{\Gamma(\frac{1}{\alpha})}{\alpha a^{\alpha/\alpha}}$$
 increases as  $\alpha \to 0$  (35)

# 5 Analogy with Central Limit Theorem?

### 5.1 Gnedenko-Doblin Theorem

 $\frac{1}{a_N} p_N \frac{x}{a_N} \to L_{\alpha,\beta}(a,x)$  iff the CDF P(x) satisfies

1. 
$$\lim_{x \to \infty} \frac{p(-x)}{a - p(x)} = \frac{1 - \beta}{1 + \beta}$$

This is akin to the "amount" of probability on one side of the tail.

$$2. \lim_{x \to \infty} \frac{1 - p(x) + p(-x)}{1 - p(rx) + p(-rx)} = r^{\alpha} \quad \forall r$$

We can think of this as saying how fast it is decaying.

## 5.2 Theorem

The distribution p(x) is in the basin of attraction of  $L_{\alpha,\beta}(a,x)$  if

$$p(x) \sim \frac{A_{\pm}}{|x|^{1+\alpha}} \quad x \to \pm \infty$$
 (36)

$$\beta = \frac{A_{+} - A_{-}}{A_{+} + A_{-}} \tag{37}$$

If p(x) has power-law tails, then the sum of RVs goes to a Lévy distribution.

## 5.3 Scaling and superdiffusivity

Suppose that  $\{x_i\}$  follow a  $L_{\alpha}(a,x)$  distribution, then  $\hat{p}(k) = \exp{-a|k|^{\alpha}}$  and then the characteristic function at  $x_N = \sum_{n=1}^N x_n$  is given by

$$\hat{p}_N(k) = [\exp -a \, |k|^{\alpha}]^N = \exp -a \, |k|^{\alpha} \, N$$
 (38)

$$p_N(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \exp ikx - aN |k|^{\alpha} dk$$
 (39)

Now we substitute:

$$k = VN^{-1/\alpha}, dk = -dVN^{-1/\alpha}$$

$$\tag{40}$$

$$p_N(k) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \exp iV\left(\frac{x}{N^{1/\alpha}}\right) - a|V|^{\alpha} \frac{\mathrm{d}V}{N^{1/\alpha}} = \frac{1}{N^{1/\alpha}} L_{\alpha}\left(a, \frac{x}{N^{1/\alpha}}\right)$$
(41)

Thus, the width scales like  $N^{1/\alpha}$ .

If  $\alpha > 2$ , we get scaling that exceeds "square-root" scaling, ie. a superdiffusive process.