

Lecture 22, May 13, 2021 Green's functions

Green's function IVP for $T_t = T_{xx}$, $-\infty < x < \infty$

$G_t = G_{xx}$ with I.D. $G = \delta(x)$ †

① If $T(x)$ a solution then $T(\alpha x, \alpha^2 t) \times$
is a sol for any const.

② $\alpha \delta(\alpha x) = \delta(x)$ $\int \alpha \delta(\alpha x) f(x) dx = \int \delta(z) f(\frac{z}{\alpha}) dz = f(0)$

† Extra condition G is bounded! Why?

$\alpha G(\alpha x, \alpha^2 t)$ solves same problem
Uniqueness of sol \Rightarrow $\alpha G(\alpha x, \alpha^2 t) = G(x, t)$

Take $\alpha = \frac{1}{\sqrt{t_0}}$ $\frac{1}{\sqrt{t_0}} G(\frac{x}{\sqrt{t_0}}, \frac{t}{t_0}) = G(x, t)$

\therefore at $t=t_0$ $\frac{1}{\sqrt{t_0}} G(\frac{x}{\sqrt{t_0}}, 1) = G(x, t_0)$

But to obtain $G(x, t) = \frac{1}{\sqrt{t}} \underbrace{G(\frac{x}{\sqrt{t}}, 1)}_{g(x/\sqrt{t})}$

$G(x, t) = \frac{1}{\sqrt{t}} g(\frac{x}{\sqrt{t}})$ $\frac{x}{\sqrt{t}} = \xi$

Substitute into
equ.

$$\frac{d^2 g}{d\xi^2} + \frac{1}{2} \xi \frac{dg}{d\xi} + \frac{1}{2} g = 0$$

$\underbrace{\hspace{10em}}_{\frac{d}{d\xi}(\frac{1}{2} \xi g)}$

$$\frac{dg}{d\xi} + \frac{1}{2} \xi g = a$$

Same argument as above for symmetry $x \rightarrow -x$ shows $\left. \begin{array}{l} G \text{ even in } x \\ \Rightarrow \\ g \text{ even in } \xi \end{array} \right\}$

$$\Rightarrow a = 0$$

$$g = C e^{-\frac{1}{4} \xi^2}$$

$$C = ?$$

$$G = \frac{C}{\sqrt{t}} e^{-\frac{x^2}{4t}}$$

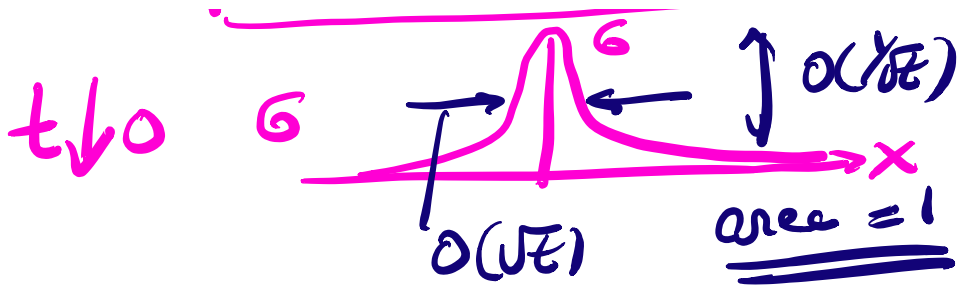
$$\int_{-\infty}^{\infty} G dx = C \int_{-\infty}^{\infty} e^{-\frac{1}{4} \xi^2} d\xi = 1$$

$\xi = \frac{x}{\sqrt{t}} \quad \underbrace{\hspace{10em}}_{2\sqrt{\pi}}$

$$C = \frac{1}{2\sqrt{\pi}}$$

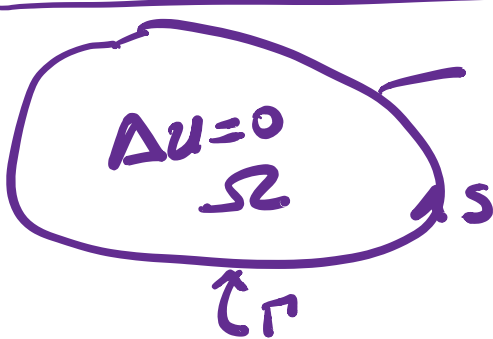
$$G = \frac{1}{2\sqrt{\pi t}} e^{-x^2/4t}$$

Green's
Function
Heat Eqn.



$$Ax = y \quad \longrightarrow \quad x = Gy \quad G = A^{-1}$$

$$\sum_m a_{nm} x_m = y_n \quad \underline{x_n = \sum_m G_{nm} y_m}$$



$u = f(s)$ on Γ

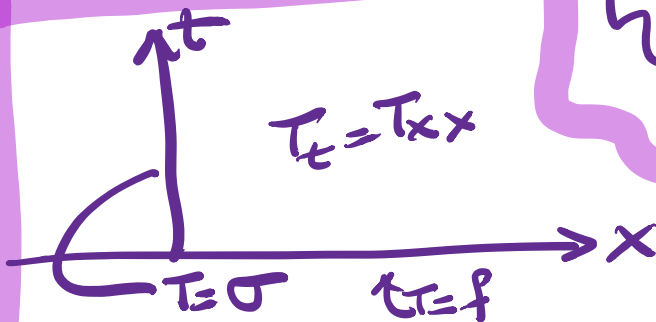
Solve

$$\Delta G = 0 \text{ in } \Omega$$

$$G = \delta(s - s_0) \text{ on } \Gamma$$

$$G = G(x, s)$$

$$u = \int_{\Gamma} G(x, s) f(s) ds$$



$$T_t = T_{xx} \quad x \geq 0 \quad t \geq 0$$

$$T = \sigma \quad \text{on } x = 0$$

$$T = f \quad \text{on } t = 0$$

Soln rule

$$T = T_1 + T_2$$

T_1 satisfies $(T_1)_t = (T_1)_{xx} \quad x > 0, t > 0$

$$\begin{aligned} \rightarrow T_1 &= 0 \text{ on } x=0 \quad \leftarrow \\ &= f \text{ on } t=0 \end{aligned}$$

T_2 satisfies $(T_2)_t \neq (T_2)_{xx} \quad x > 0, t > 0$

$$\begin{aligned} T_2 &= \sigma \text{ on } x=0 \\ T_2 &= 0 \text{ on } t=0 \end{aligned}$$

Solve for T_1

$-T(-x, t)$ is a solution

T_1 problem is equivalent to

$$\begin{cases} T_t = T_{xx} & -\infty < x < \infty, t > 0 \\ T = \tilde{f}(x) & \text{for } t=0 \end{cases} \quad (x > 0)$$

where \tilde{f} is odd and $\tilde{f}(x) = -\tilde{f}(-x)$

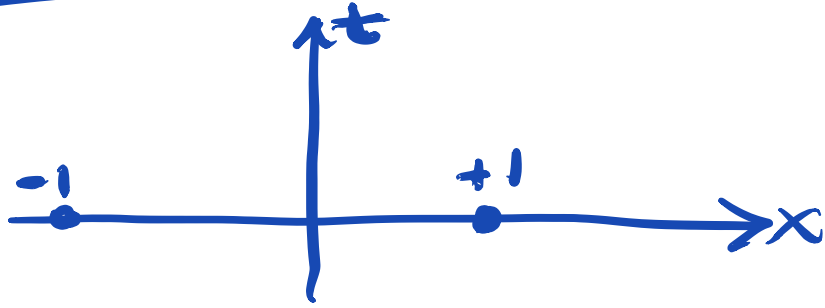
Green's function solve with

$$\text{i.i.d. } T = \int \delta(x-y) - \delta(x+y)$$

$$G(x, y) = \frac{1}{2\sqrt{\pi t}} \left[e^{-\frac{(x-y)^2}{4t}} - e^{-\frac{(x+y)^2}{4t}} \right]$$

Green's function for $\left\{ \begin{array}{l} T_t = T_{xx} \quad x > 0, t > 0 \\ T = f \quad \text{at } t = 0 \end{array} \right.$

$$T = \int_0^{\infty} G(x, y) f(y) dy$$



Same idea: $T_t = T_{xx} \quad x > 0, t > 0$
 $T = f \quad t = 0$
insulator on $x=0 \rightarrow T_x = 0 \quad \text{on } x = 0$

$$G = \frac{1}{2\sqrt{\pi t}} \left[e^{-\frac{(x-y)^2}{4t}} + e^{-\frac{(x+y)^2}{4t}} \right]$$

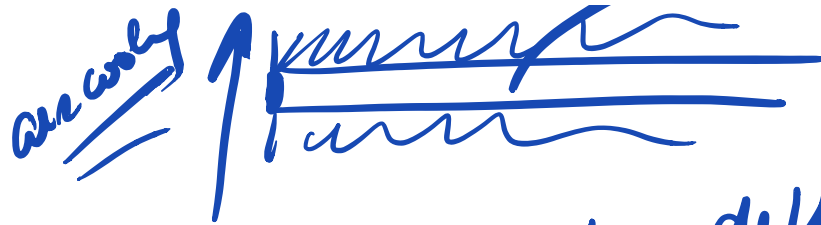
Third prob

Prescribe $T + \alpha T_x$ on $x=0$

\Rightarrow
Robin B.C

$$T + \alpha T_x = 0 \quad \text{on } \underline{\underline{x=0}}$$

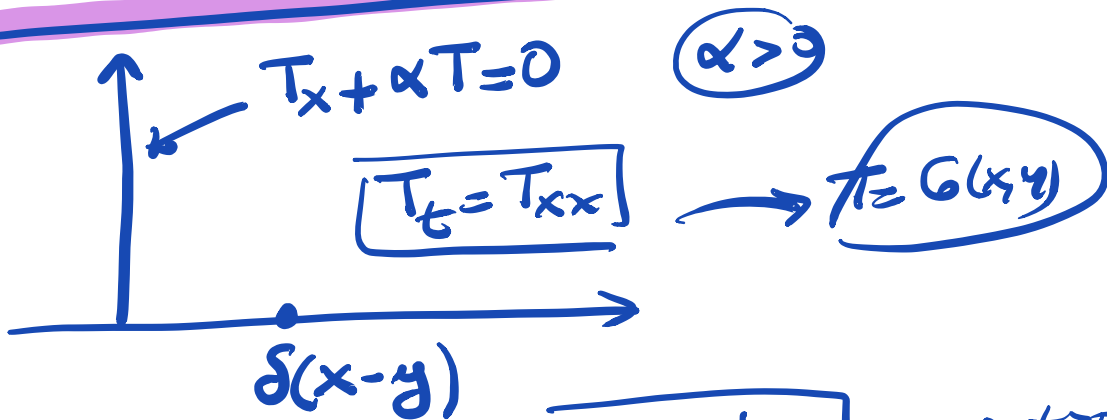
$$T_t = T_{xx}$$



heat flux \propto Temperature difference

$$-\kappa T_x = C(T - T_0) \quad \underline{C > 0}$$

$$\boxed{\kappa T_x + C T = C T_0(t)} \quad \text{Robin b.c.}$$



$$\boxed{\phi = T_x + \alpha T}$$

$$\boxed{\phi_t = \phi_{xx}} \quad \underline{x > 0, t > 0}$$

$$\phi = 0 \quad \text{on } \underline{x = 0}$$

and

$$\phi = \delta'(x-y) + \alpha \delta(x-y)$$

Extend to $-\infty < x < \infty$ and solve with

$$\boxed{\begin{aligned} \phi &= \delta'(x-y) + \alpha \delta(x-y) \\ &\quad - \delta'(x+y) + \alpha \delta(x+y) \end{aligned}}$$

\uparrow Green's function ϕ .

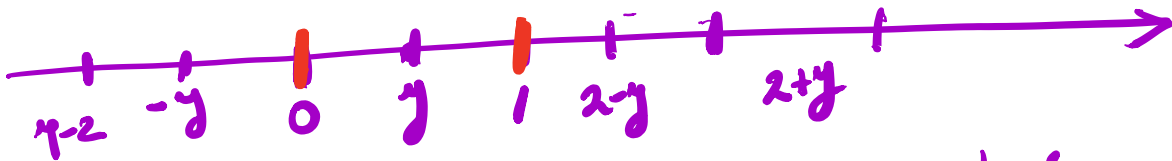
$$T = e^{-\alpha x} \int e^{\alpha y} \phi(y) dy$$

Finite interval

$$T_t = T_{xx} \quad 0 < x < 1$$

$$T_x = 0 \quad \text{on } x=0, 1$$

$$T = \delta(x-y) \quad \text{for } t=0 \text{ some } 0 < y < 1$$



Del'tas at $x = y + 2n$ | n units
 $-y + 2n$

$$G = \frac{1}{2\sqrt{\pi t}} \sum_n e^{-\frac{(x-y_n)^2}{4t}}$$

$$G = G(x, y, t)$$

$$T_t = T_{xx} \quad 0 < x < 1$$

$$T = f(x) \quad t=0$$

$$T = 0 \quad \text{at } x=0, 1$$

Can write
see two
ways

$$T = \int_0^1 G(x, y) f(y) dy$$

Good at t is small

What about to solve

$$f(x) = \sum_1^{\infty} f_n \sin(nx)$$

$$T = \sum_1^{\infty} f_n \sin(nx) e^{-n^2 t}$$

Spectral
method

Eigenvalues and eigenvectors for
self adjoint matrices

$$\langle A\vec{u}, \vec{v} \rangle = \langle \vec{u}, A\vec{v} \rangle$$

$A = A^* \Rightarrow A$ has orthonormal
basis of eigenvectors

Eigenvalues are real $A\vec{u} = \lambda\vec{u}$

$$\langle A\vec{u}, \vec{u} \rangle = \langle \vec{u}, A\vec{u} \rangle = \lambda \frac{\langle \vec{u}, \vec{u} \rangle}{\|\vec{u}\|^2}$$
$$\lambda \frac{\langle \vec{u}, \vec{u} \rangle}{\|\vec{u}\|^2} = \lambda \frac{\langle \vec{u}, \vec{u} \rangle}{\|\vec{u}\|^2} \quad \lambda = \bar{\lambda}$$

Take A has one eigenvector

$$A\vec{u}_1 = \lambda_1 \vec{u}_1, \quad \|\vec{u}_1\| = 1$$

Let H be set of all vectors that
are $\perp \vec{u}_1$ Hyperplane

Take $\vec{v} \in H$ then $A\vec{v}$ in H

$$\langle A\vec{v}, \vec{u}_1 \rangle = \langle \vec{v}, A\vec{u}_1 \rangle = \lambda_1 \langle \vec{v}, \vec{u}_1 \rangle = 0$$

\therefore Can think of A as operator in H

Also self adjoint in H

Then use induction!

For matrices this works also for

normal matrices

$$\boxed{AA^+ = A^+A}$$

Open \rightarrow her extension to ∞ dim.

If A is not normal
big unknown for ∞ operators

Example self adjoint etc

Heat

$$T_t = \Delta u$$

\uparrow self-adjoint for
reasonable B.C.



$$\langle u, \Delta v \rangle = \int_{\Omega} \frac{u \cdot \Delta v}{L} dV = \int_{\Omega} \text{div}(u \nabla v) - (\nabla u \cdot \nabla v)$$

$$= \int_{\partial \Omega} (u \cdot \nabla v) \hat{n} dA - \int_{\Omega} (\nabla u \cdot \nabla v) dV$$

$$\langle \Delta u, v \rangle = \int_{\partial \Omega} (\nabla u \cdot \nabla v) \hat{n} dA - \int_{\Omega} (\nabla u \cdot \nabla v) dV$$

if this vanishes
then self-adjoint

$$u_t = u_x + u_{xxx} \quad L = \partial_x + \partial_x^3$$

skew-adjoint on $-\infty < x < \infty$

$$\langle Lu, v \rangle = - \langle u, Lv \rangle$$

$$L^+ = -L$$

iL is self-adjoint

\therefore eigenvalues L purely $im.$

$$u_t = \frac{u_x + u_{xxx}}{Lu} \quad \underline{x > 0} \quad \text{plus some B.C. at } x = \infty$$

$$\begin{aligned}
 \langle Lu, v \rangle &= \int_0^\infty (u_x + u_{xxx})v \, dx = \\
 &= (uv) \Big|_0^\infty + (v u_{xx}) \Big|_0^\infty - \int_0^\infty (uv_x + u_{xx}v) \, dx \\
 &= \underbrace{(uv + v u_{xx} - v_x u_x + v_{xx} u)}_{\langle u, Lv \rangle} \Big|_0^\infty - \int_0^\infty u (v_x + v_{xxx}) \, dx
 \end{aligned}$$