

Tensor Product Decompositions for Modules Over Subregular W -algebras

Brian Li

Mission San Jose High School

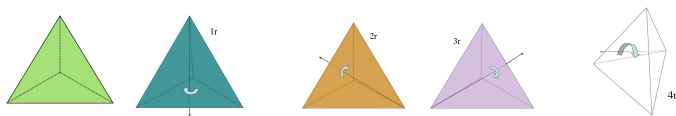
October 15, 2023

Concept of symmetry is very important in mathematics and physics (e.g. gauge theory).

- Mathematically formalized by groups

Instead of a formal definition, we give some examples of groups:

- 1 The symmetric group S_n , set of all permutations of n elements;
 - A bijection of a set onto itself, definition of "symmetry"
 - Symmetry group S_4 is an isometric permutation of vertices for tetrahedron:



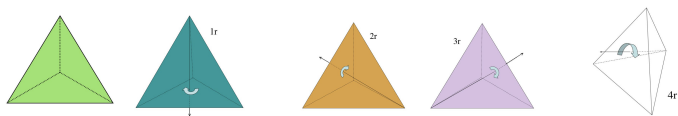
- 2 The group of invertible matrices GL_N over a field F with matrix multiplication.

Concept of symmetry is very important in mathematics and physics (e.g. gauge theory).

- Mathematically formalized by groups

Instead of a formal definition, we give some examples of groups:

- 1 The symmetric group S_n , set of all permutations of n elements;
 - A bijection of a set onto itself, definition of "symmetry"
 - Symmetry group S_4 is an isometric permutation of vertices for tetrahedron:



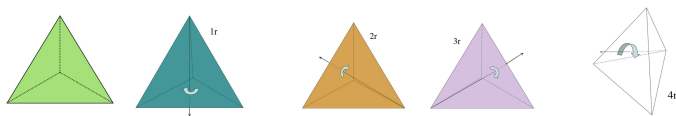
- 2 The group of invertible matrices GL_N over a field F with matrix multiplication.

Concept of symmetry is very important in mathematics and physics (e.g. gauge theory).

- Mathematically formalized by groups

Instead of a formal definition, we give some examples of groups:

- 1 The symmetric group S_n , set of all permutations of n elements;
 - A bijection of a set onto itself, definition of "symmetry"
 - Symmetry group S_4 is an isometric permutation of vertices for tetrahedron:



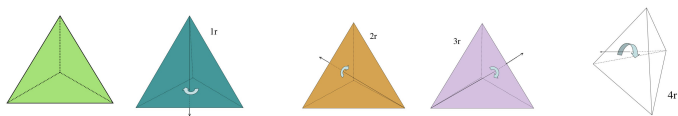
- 2 The group of invertible matrices GL_N over a field F with matrix multiplication.

Concept of symmetry is very important in mathematics and physics (e.g. gauge theory).

- Mathematically formalized by groups

Instead of a formal definition, we give some examples of groups:

- 1 The symmetric group S_n , set of all permutations of n elements;
 - A bijection of a set onto itself, definition of "symmetry"
 - Symmetry group S_4 is an isometric permutation of vertices for tetrahedron:

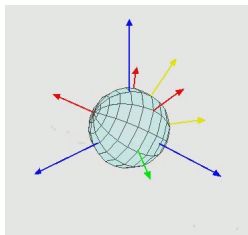


- 2 The group of invertible matrices GL_N over a field F with matrix multiplication.

When symmetries are continuous: **Lie groups**.

For example:

- Rotations of a sphere $SO(3)$



- Special linear groups $SL(2)$ given by

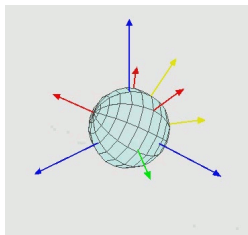
$$SL(2) = \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \mid ad - bc = 1 \right\}$$

Slogan 1: groups are difficult, infinitesimal transformations are easier — **Lie algebras**.

When symmetries are continuous: **Lie groups**.

For example:

- Rotations of a sphere $SO(3)$



- Special linear groups $SL(2)$ given by

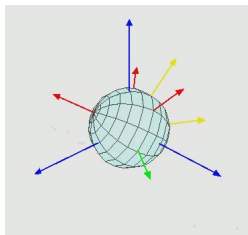
$$SL(2) = \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \mid ad - bc = 1 \right\}$$

Slogan 1: groups are difficult, infinitesimal transformations are easier — **Lie algebras**.

When symmetries are continuous: **Lie groups**.

For example:

- Rotations of a sphere $SO(3)$



- Special linear groups $SL(2)$ given by

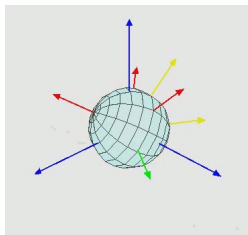
$$SL(2) = \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \mid ad - bc = 1 \right\}$$

Slogan 1: groups are difficult, infinitesimal transformations are easier — **Lie algebras**.

When symmetries are continuous: **Lie groups**.

For example:

- Rotations of a sphere $SO(3)$



- Special linear groups $SL(2)$ given by

$$SL(2) = \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \mid ad - bc = 1 \right\}$$

Slogan 1: groups are difficult, infinitesimal transformations are easier — **Lie algebras**.

Examples — derivations:

① $D_1 = \frac{d}{dx}$, infinitesimal version of *translations* by Taylor's formula:

$$f(x+t) = f(x) + t \cdot df/dx + O(t^2).$$

② $D_2 = x \frac{d}{dx}$: infinitesimal version of *dilations* $f(e^t x)$.

Observe both D_1, D_2 are derivations. However, they do not have an algebra structure:

- $D_1 \circ D_2$ is not a derivation
- But $D_1 \circ D_2 - D_2 \circ D_1 = \frac{d}{dx}$ is

We denote $D_1 \circ D_2 - D_2 \circ D_1$ by $[D_1, D_2]$.

Definition

A **Lie algebra** is a vector space \mathfrak{g} equipped with the skew-symmetric bilinear map $[-, -]$ satisfying the *Jacobi identity*

$$[a, [b, c]] + [b, [c, a]] + [c, [a, b]] = 0, \quad a, b, c \in \mathfrak{g}.$$

Examples:

- The set of derivations D
- \mathfrak{gl}_n : set of $n \times n$ matrices with commutator $[A, B] := A \cdot B - B \cdot A$;

Examples — derivations:

- 1 $D_1 = \frac{d}{dx}$, infinitesimal version of *translations* by Taylor's formula:

$$f(x + t) = f(x) + t \cdot df/dx + O(t^2).$$

- 2 $D_2 = x \frac{d}{dx}$: infinitesimal version of *dilations* $f(e^t x)$.

Observe both D_1, D_2 are derivations. However, they do not have an algebra structure:

- $D_1 \circ D_2$ is not a derivation
- But $D_1 \circ D_2 - D_2 \circ D_1 = \frac{d}{dx}$ is

We denote $D_1 \circ D_2 - D_2 \circ D_1$ by $[D_1, D_2]$.

Definition

A **Lie algebra** is a vector space \mathfrak{g} equipped with the skew-symmetric bilinear map $[-, -]$ satisfying the *Jacobi identity*

$$[a, [b, c]] + [b, [c, a]] + [c, [a, b]] = 0, \quad a, b, c \in \mathfrak{g}.$$

Examples:

- The set of derivations D
- \mathfrak{gl}_n : set of $n \times n$ matrices with commutator $[A, B] := A \cdot B - B \cdot A$;

Examples — derivations:

- ① $D_1 = \frac{d}{dx}$, infinitesimal version of *translations* by Taylor's formula:

$$f(x+t) = f(x) + t \cdot df/dx + O(t^2).$$

- ② $D_2 = x \frac{d}{dx}$: infinitesimal version of *dilations* $f(e^t x)$.

Observe both D_1, D_2 are derivations. However, they do not have an algebra structure:

- $D_1 \circ D_2$ is not a derivation
- But $D_1 \circ D_2 - D_2 \circ D_1 = \frac{d}{dx}$ is

We denote $D_1 \circ D_2 - D_2 \circ D_1$ by $[D_1, D_2]$.

Definition

A **Lie algebra** is a vector space \mathfrak{g} equipped with the skew-symmetric bilinear map $[-, -]$ satisfying the *Jacobi identity*

$$[a, [b, c]] + [b, [c, a]] + [c, [a, b]] = 0, \quad a, b, c \in \mathfrak{g}.$$

Examples:

- The set of derivations D
- \mathfrak{gl}_n : set of $n \times n$ matrices with commutator $[A, B] := A \cdot B - B \cdot A$;

Examples — derivations:

- 1 $D_1 = \frac{d}{dx}$, infinitesimal version of *translations* by Taylor's formula:

$$f(x + t) = f(x) + t \cdot df/dx + O(t^2).$$

- 2 $D_2 = x \frac{d}{dx}$: infinitesimal version of *dilations* $f(e^t x)$.

Observe both D_1, D_2 are derivations. However, they do not have an algebra structure:

- $D_1 \circ D_2$ is not a derivation
- But $D_1 \circ D_2 - D_2 \circ D_1 = \frac{d}{dx}$ is

We denote $D_1 \circ D_2 - D_2 \circ D_1$ by $[D_1, D_2]$.

Definition

A **Lie algebra** is a vector space \mathfrak{g} equipped with the skew-symmetric bilinear map $[-, -]$ satisfying the *Jacobi identity*

$$[a, [b, c]] + [b, [c, a]] + [c, [a, b]] = 0, \quad a, b, c \in \mathfrak{g}.$$

Examples:

- The set of derivations D
- \mathfrak{gl}_n : set of $n \times n$ matrices with commutator $[A, B] := A \cdot B - B \cdot A$;

Examples — derivations:

- ① $D_1 = \frac{d}{dx}$, infinitesimal version of *translations* by Taylor's formula:

$$f(x + t) = f(x) + t \cdot df/dx + O(t^2).$$

- ② $D_2 = x \frac{d}{dx}$: infinitesimal version of *dilations* $f(e^t x)$.

Observe both D_1, D_2 are derivations. However, they do not have an algebra structure:

- $D_1 \circ D_2$ is not a derivation
- But $D_1 \circ D_2 - D_2 \circ D_1 = \frac{d}{dx}$ is

We denote $D_1 \circ D_2 - D_2 \circ D_1$ by $[D_1, D_2]$.

Definition

A **Lie algebra** is a vector space \mathfrak{g} equipped with the skew-symmetric bilinear map $[-, -]$ satisfying the *Jacobi identity*

$$[a, [b, c]] + [b, [c, a]] + [c, [a, b]] = 0, \quad a, b, c \in \mathfrak{g}.$$

Examples:

- The set of derivations D
- \mathfrak{gl}_n : set of $n \times n$ matrices with commutator $[A, B] := A \cdot B - B \cdot A$;

Slogan 2: groups are difficult, linear algebra is also difficult, but well-studied. Hence: study groups via linear algebra — *representation theory*. E.g., representations correspond to particles in physics.

Definition

A *representation* of \mathfrak{g} is a vector space \mathbb{C}^n with a map of Lie algebras $\mathfrak{g} \rightarrow \mathfrak{gl}_n$.

This map represents each element in \mathfrak{g} as a matrix.

Examples:

- Tautologically, \mathbb{C}^n for \mathfrak{gl}_n

Slogan 2: groups are difficult, linear algebra is also difficult, but well-studied. Hence: study groups via linear algebra — *representation theory*. E.g., representations correspond to particles in physics.

Definition

A *representation* of \mathfrak{g} is a vector space \mathbb{C}^n with a map of Lie algebras $\mathfrak{g} \rightarrow \mathfrak{gl}_n$.

This map represents each element in \mathfrak{g} as a matrix.

Examples:

- Tautologically, \mathbb{C}^n for \mathfrak{gl}_n

Slogan 2: groups are difficult, linear algebra is also difficult, but well-studied. Hence: study groups via linear algebra — *representation theory*. E.g., representations correspond to particles in physics.

Definition

A *representation* of \mathfrak{g} is a vector space \mathbb{C}^n with a map of Lie algebras $\mathfrak{g} \rightarrow \mathfrak{gl}_n$.

This map represents each element in \mathfrak{g} as a matrix.

Examples:

- Tautologically, \mathbb{C}^n for \mathfrak{gl}_n

Representations of \mathfrak{sl}_2

Main object for today: \mathfrak{sl}_2 .

$$\mathfrak{sl}_2 = \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \mid a + d = 0 \right\}$$

Over complex numbers: same as the Lie algebra of $SO(3)$.

For representations, take

$$E \mapsto \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}, \quad F \mapsto \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix}, \quad H \mapsto \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}.$$

More abstractly: spanned by E, F, H with relations

$$[H, E] = 2E, \quad [H, F] = -2F, \quad [E, F] = H.$$

How to study representations? Basic building block — irreducibles:

Definition

A representation V is *irreducible* if it does not contain a non-trivial subrepresentation.

Representations of \mathfrak{sl}_2

Main object for today: \mathfrak{sl}_2 .

$$\mathfrak{sl}_2 = \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \mid a + d = 0 \right\}$$

Over complex numbers: same as the Lie algebra of $SO(3)$.

For representations, take

$$E \mapsto \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}, \quad F \mapsto \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix}, \quad H \mapsto \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}.$$

More abstractly: spanned by E, F, H with relations

$$[H, E] = 2E, \quad [H, F] = -2F, \quad [E, F] = H.$$

How to study representations? Basic building block — irreducibles:

Definition

A representation V is *irreducible* if it does not contain a non-trivial subrepresentation.

Main object for today: \mathfrak{sl}_2 .

$$\mathfrak{sl}_2 = \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \mid a + d = 0 \right\}$$

Over complex numbers: same as the Lie algebra of $SO(3)$.

For representations, take

$$E \mapsto \begin{bmatrix} 0 & 1 \\ 0 & 0 \end{bmatrix}, \quad F \mapsto \begin{bmatrix} 0 & 0 \\ 1 & 0 \end{bmatrix}, \quad H \mapsto \begin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}.$$

More abstractly: spanned by E, F, H with relations

$$[H, E] = 2E, \quad [H, F] = -2F, \quad [E, F] = H.$$

How to study representations? Basic building block — irreducibles:

Definition

A representation V is **irreducible** if it does not contain a non-trivial subrepresentation.

Irreducible representations

For \mathfrak{sl}_2 — complete classification:

Theorem

Irreducible finite-dimensional representations V_n of \mathfrak{sl}_2 are classified by a natural number n and are of the form $V_n := \text{span}(v, Fv, F^2v, \dots, F^nv)$, where the vector v satisfies $Ev = 0$ (highest-weight vector).

Natural operation on representations: tensor product (for instance, corresponds to combined system of particles)

- Decomposition of tensor products are actually completely determined by highest weight vectors

Theorem (Clebsch-Gordan)

For irreducible representations V_n, V_m , we have $V_n \otimes V_m \cong \bigoplus_{k=0}^{\min(n,m)} V_{n+m-2k}$.

For instance, let $\mathbb{C}^2 = \text{span}(v_1, v_2)$ where $v_1 = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$, $v_2 = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$. Then, highest-weight vectors of $\mathbb{C}^2 \otimes \mathbb{C}^2$ are

$$(v_1 \otimes v_1, v_1 \otimes v_2 - v_2 \otimes v_1) \in V_2 \oplus V_0.$$

Irreducible representations

For \mathfrak{sl}_2 — complete classification:

Theorem

Irreducible finite-dimensional representations V_n of \mathfrak{sl}_2 are classified by a natural number n and are of the form $V_n := \text{span}(v, Fv, F^2v, \dots, F^nv)$, where the vector v satisfies $Ev = 0$ (highest-weight vector).

Natural operation on representations: tensor product (for instance, corresponds to combined system of particles)

- Decomposition of tensor products are actually completely determined by highest weight vectors

Theorem (Clebsch-Gordan)

For irreducible representations V_n, V_m , we have $V_n \otimes V_m \cong \bigoplus_{k=0}^{\min(n,m)} V_{n+m-2k}$.

For instance, let $\mathbb{C}^2 = \text{span}(v_1, v_2)$ where $v_1 = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$, $v_2 = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$. Then, highest-weight vectors of $\mathbb{C}^2 \otimes \mathbb{C}^2$ are

$$(v_1 \otimes v_1, v_1 \otimes v_2 - v_2 \otimes v_1) \in V_2 \oplus V_0.$$

Irreducible representations

For \mathfrak{sl}_2 — complete classification:

Theorem

Irreducible finite-dimensional representations V_n of \mathfrak{sl}_2 are classified by a natural number n and are of the form $V_n := \text{span}(v, Fv, F^2v, \dots, F^nv)$, where the vector v satisfies $Ev = 0$ (highest-weight vector).

Natural operation on representations: tensor product (for instance, corresponds to combined system of particles)

- Decomposition of tensor products are actually completely determined by highest weight vectors

Theorem (Clebsch-Gordan)

For irreducible representations V_n, V_m , we have $V_n \otimes V_m \cong \bigoplus_{k=0}^{\min(n,m)} V_{n+m-2k}$.

For instance, let $\mathbb{C}^2 = \text{span}(v_1, v_2)$ where $v_1 = \begin{bmatrix} 1 \\ 0 \end{bmatrix}$, $v_2 = \begin{bmatrix} 0 \\ 1 \end{bmatrix}$. Then, highest-weight vectors of $\mathbb{C}^2 \otimes \mathbb{C}^2$ are

$$(v_1 \otimes v_1, v_1 \otimes v_2 - v_2 \otimes v_1) \in V_2 \oplus V_0.$$

To decompose $V_n \otimes V_m$, enough to find highest-weight vectors in this tensor product.

- What if $E v = v$? Called **Whittaker vectors**, generate *Whittaker modules*. Naturally arise in physics (Toda system).
- Decomposition of Whittaker modules: likewise, completely classified by Whittaker vectors $\text{Whit}(\mathcal{W})$.

Theorem (Kalmykov, 2021)

For any Whittaker module \mathcal{W} and a finite-dimensional representation V of \mathfrak{sl}_2 , we have $\text{Whit}(\mathcal{W} \otimes V) \cong \text{Whit}(\mathcal{W}) \otimes V$ canonically.

To decompose $V_n \otimes V_m$, enough to find highest-weight vectors in this tensor product.

- What if $E v = v$? Called **Whittaker vectors**, generate *Whittaker modules*. Naturally arise in physics (Toda system).
- Decomposition of Whittaker modules: likewise, completely classified by Whittaker vectors $\text{Whit}(\mathcal{W})$.

Theorem (Kalmykov, 2021)

For any Whittaker module \mathcal{W} and a finite-dimensional representation V of \mathfrak{sl}_2 , we have $\text{Whit}(\mathcal{W} \otimes V) \cong \text{Whit}(\mathcal{W}) \otimes V$ canonically.

Application: non-standard quantization of SL_2 . Two ways to compute Whittaker vectors in $W \otimes U \otimes V$:

$$\text{Whit}(W \otimes U \otimes V) \cong \text{Whit}(W \otimes U) \otimes V \cong (\text{Whit}(W) \otimes U) \otimes V,$$

$$\text{Whit}(W \otimes U \otimes V) \cong \text{Whit}(W) \otimes (U \otimes V).$$

Differ by the action on $U \otimes V$ of

$$J = \sum_k J_k^{(1)} \otimes J_k^{(2)} = \sum_{k \geq 0} \frac{(-1)^k}{2^k k!} F^k \otimes \prod_{i=0}^{k-1} (H - 2i).$$

Deforms multiplication on functions on SL_2 :

$$f * g := \sum_i J_i^{(1)}(f) \cdot J_i^{(2)}(g), \quad f, g \in \text{Fun}(SL_2).$$

Application: non-standard quantization of SL_2 . Two ways to compute Whittaker vectors in $W \otimes U \otimes V$:

$$\text{Whit}(W \otimes U \otimes V) \cong \text{Whit}(W \otimes U) \otimes V \cong (\text{Whit}(W) \otimes U) \otimes V,$$

$$\text{Whit}(W \otimes U \otimes V) \cong \text{Whit}(W) \otimes (U \otimes V).$$

Differ by the action on $U \otimes V$ of

$$J = \sum_k J_k^{(1)} \otimes J_k^{(2)} = \sum_{k \geq 0} \frac{(-1)^k}{2^k k!} F^k \otimes \prod_{i=0}^{k-1} (H - 2i).$$

Deforms multiplication on functions on SL_2 :

$$f * g := \sum_i J_i^{(1)}(f) \cdot J_i^{(2)}(g), \quad f, g \in \text{Fun}(SL_2).$$

Application: non-standard quantization of SL_2 . Two ways to compute Whittaker vectors in $W \otimes U \otimes V$:

$$\text{Whit}(W \otimes U \otimes V) \cong \text{Whit}(W \otimes U) \otimes V \cong (\text{Whit}(W) \otimes U) \otimes V,$$

$$\text{Whit}(W \otimes U \otimes V) \cong \text{Whit}(W) \otimes (U \otimes V).$$

Differ by the action on $U \otimes V$ of

$$J = \sum_k J_k^{(1)} \otimes J_k^{(2)} = \sum_{k \geq 0} \frac{(-1)^k}{2^k k!} F^k \otimes \prod_{i=0}^{k-1} (H - 2i).$$

Deforms multiplication on functions on SL_2 :

$$f * g := \sum_i J_i^{(1)}(f) \cdot J_i^{(2)}(g), \quad f, g \in \text{Fun}(SL_2).$$

Application: non-standard quantization of SL_2 . Two ways to compute Whittaker vectors in $W \otimes U \otimes V$:

$$\text{Whit}(W \otimes U \otimes V) \cong \text{Whit}(W \otimes U) \otimes V \cong (\text{Whit}(W) \otimes U) \otimes V,$$

$$\text{Whit}(W \otimes U \otimes V) \cong \text{Whit}(W) \otimes (U \otimes V).$$

Differ by the action on $U \otimes V$ of

$$J = \sum_k J_k^{(1)} \otimes J_k^{(2)} = \sum_{k \geq 0} \frac{(-1)^k}{2^k k!} F^k \otimes \prod_{i=0}^{k-1} (H - 2i).$$

Deforms multiplication on functions on SL_2 :

$$f * g := \sum_i J_i^{(1)}(f) \cdot J_i^{(2)}(g), \quad f, g \in \text{Fun}(SL_2).$$

Generalization: ***W*-algebras** (Whittaker Modules for \mathfrak{gl}_n).

Our research: tensor product decomposition for subregular *W*-algebras.

Theorem (Kalmykov-L., 2023)

For any subregular Whittaker module \mathcal{W} and the vector representation V of \mathfrak{gl}_n , there is an explicit identification

$$\text{Whit}(\mathcal{W} \otimes V) \cong \text{Whit}(\mathcal{W}) \otimes V.$$

In particular, allows to construct canonically Whittaker vectors in $\mathcal{W} \otimes U$ for *any* finite-dimensional representation U of \mathfrak{gl}_n .

Likewise, gives non-standard quantization of the group GL_N .

I would like to kindly thank:

- My mentor, Dr. Artem Kalmykov, for guiding me through the tough mathematical readings and being patient with me throughout the entire research process
- MIT PRIMES organizers, in particular Prof. Pavel Etingof, Dr. Slava Gerovitch, and Dr. Tanya Khovanova, for providing this wonderful opportunity for me to do math research
- My parents for always being so supportive.

- T. Arakawa. “Introduction to W -algebras and their representation theory”. *Perspectives in Lie theory*. Vol. 19. Springer INdAM Ser. Springer, Cham, 2017, pp. 179–250.
- P. Etingof and O. Schiffmann. “Lectures on the dynamical Yang-Baxter equations”. *Quantum Groups and Lie Theory (Durham, 1999)*, London Math. Soc. Lecture Note Ser 290 (2001), pp. 89– 129
- P. Etingof and O. Schiffmann. *Lectures on quantum groups*. Second. *Lectures in Mathematical Physics*. International Press, Somerville, MA, 2002, pp. xii+242.
- W. Fulton and J. Harris. *Representation theory*. Vol. 129. *Graduate Texts in Mathematics*. A first course, *Readings in Mathematics*. Springer-Verlag, New York, 1991, pp. xvi+551. URL: <https://doi.org/10.1007/978-1-4612-0979-9>.
- S. M. Goodwin. “Translation for finite W -algebras”. *Represent. Theory* 15 (2011), pp. 307–346. URL: <https://doi.org/10.1090/S1088-4165-2011-00388-5>.

- J. E. Humphreys. Introduction to Lie algebras and representation theory. Vol. 9. Graduate Texts in Mathematics. Second printing, revised. Springer-Verlag, New York-Berlin, 1978, pp. xii+171.
- A. Kalmykov. Geometric and categorical approaches to dynamical representation theory. eng. Zürich, 2021.
- B. Kostant. “On Whittaker vectors and representation theory”. Invent. Math. 48.2 (1978), pp. 101–184. URL: <https://doi.org/10.1007/BF01390249>.
- I. Losev. “Finite W -algebras”. Proceedings of the International Congress of Mathematicians. Volume III. Hindustan Book Agency, New Delhi, 2010, pp. 1281–1307.